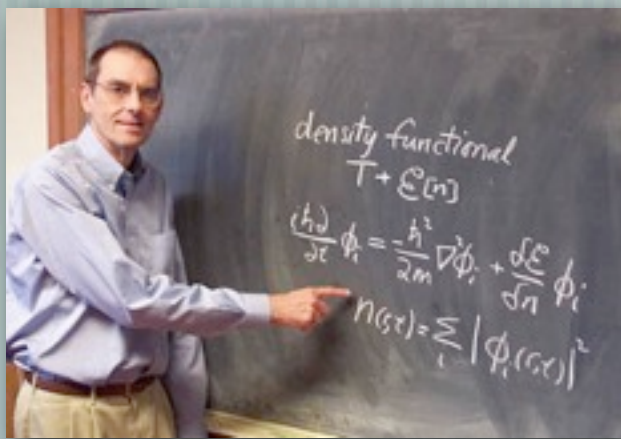


Photons from Relativistic Heavy Ion Collisions: Progress and Puzzles

OUTLINE

- Sources & EM emissivity: Rates
- Modelling the evolving system:
 - 3D hydro
 - 3D viscous hydro
 - Fluctuating initial states
- How are the photon yields dependent on the dynamics?
- Is it the same for dileptons?
- Status of our understanding of the data



BertschFest



Charles Gale
McGill University

find a bertsch, g and t photon

Brief format

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Sort by:

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latest first :: desc. :: - or rank by - :: 25 results :: single list ::

HEP

6 records found

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1. Hard photons in proton-nucleus collisions.

K. Nakayama (Julich, Forschungszentrum), G.F. Bertsch (Michigan State U., NSCL). Dec 1989. 6 pp.

Published in **Phys.Rev. C40 (1989) 2520-2526**

DOI: [10.1103/PhysRevC.40.2520](https://doi.org/10.1103/PhysRevC.40.2520)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[Journal Server: Phys. Rev. C Server](#)

[Detailed record](#) - [Cited by 4 records](#)

2. Potential model calculation of energetic photon production from heavy ion collisions.

K. Nakayama (Julich, Forschungszentrum & Georgia U.), G.F. Bertsch (Michigan State U., NSCL). Aug 1989. 6 pp.

Published in **Phys.Rev. C40 (1989) 685-691**

DOI: [10.1103/PhysRevC.40.685](https://doi.org/10.1103/PhysRevC.40.685)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[Journal Server: Phys. Rev. C Server](#)

[Detailed record](#) - [Cited by 6 records](#)

3. Production Of Energetic Photons Via Excitation Of The Delta Resonance In Heavy Ion Collisions.

W. Bauer, G.F. Bertsch (Michigan State U.). 1989.

Published in **Phys.Lett. B229 (1989) 16-19**

DOI: [10.1016/0370-2693\(89\)90146-9](https://doi.org/10.1016/0370-2693(89)90146-9)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[Journal Server: KEK scanned document](#) ; [ADS Abstract Service](#); [Science Direct](#)

[Detailed record](#) - [Cited by 3 records](#)

4. Ultradipole photons in low-energy heavy-ion reactions.

K. Nakayama, G.F. Bertsch (Michigan State U.). Nov 1987. 4 pp.

Published in **Phys.Rev. C36 (1987) 1848-1852**

DOI: [10.1103/PhysRevC.36.1848](https://doi.org/10.1103/PhysRevC.36.1848)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[Journal Server: Phys. Rev. C Server](#)

[Detailed record](#) - [Cited by 9 records](#)

5. High energy photon production in nuclear collisions.

K. Nakayama, G. Bertsch (Michigan State U., NSCL). Dec 1986. 10 pp.

Published in **Phys.Rev. C34 (1986) 2190-2200**

DOI: [10.1103/PhysRevC.34.2190](https://doi.org/10.1103/PhysRevC.34.2190)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[Journal Server: Phys. Rev. C Server](#)

[Detailed record](#) - [Cited by 29 records](#)

6. Energetic photons from intermediate energy proton- and heavy-ion-induced reactions.

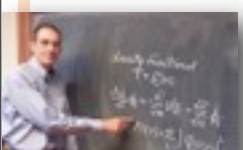
W. Bauer, G.F. Bertsch, Wolfgang Cassing, Ulrich Mosel (Michigan State U., NSCL & Giessen U.). Dec 1986. 6 pp.

Published in **Phys.Rev. C34 (1986) 2127-2133**

DOI: [10.1103/PhysRevC.34.2127](https://doi.org/10.1103/PhysRevC.34.2127)

[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
[Journal Server: Phys. Rev. C Server](#)

[Detailed record](#) - [Cited by 85 records](#)



Sources of photons in a relativistic nuclear collision:

Hard direct photons. pQCD with shadowing
Non-thermal

Fragmentation photons. pQCD with shadowing
Non-thermal

Thermal photons
Thermal

Jet-plasma photons
Thermal

Jet in-medium bremsstrahlung
Thermal



INFO CARRIED BY THE RADIATION

$$dR = -\frac{g^{\mu\nu}}{2\omega} \frac{d^3k}{(2\pi)^3} \frac{1}{Z} \sum_i e^{-\beta K_i} \sum_f (2\pi)^4 \delta(p_i - p_f - k) \\ \times \langle j | J_\mu | i \rangle \langle i | J_\nu | j \rangle$$

Thermal ensemble average of the current-current correlator

Emission rates:

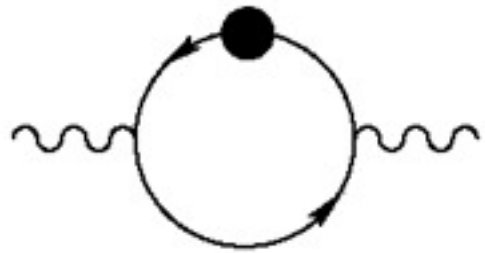
$$\omega \frac{d^3 R}{d^3 k} = -\frac{g^{\mu\nu}}{(2\pi)^3} \text{Im} \Pi_{\mu\nu}^R(\omega, k) \frac{1}{e^{\beta\omega} - 1} \quad (\text{photons})$$

$$E_+ E_- \frac{d^6 R}{d^3 p_+ d^3 p_-} = \frac{2e^2}{(2\pi)^6} \frac{1}{k^4} L^{\mu\nu} \text{Im} \Pi_{\mu\nu}^R(\omega, k) \frac{1}{e^{\beta\omega} - 1} \quad (\text{dileptons})$$

McLerran, Toimela (85), Weldon (90), Gale, Kapusta (91)



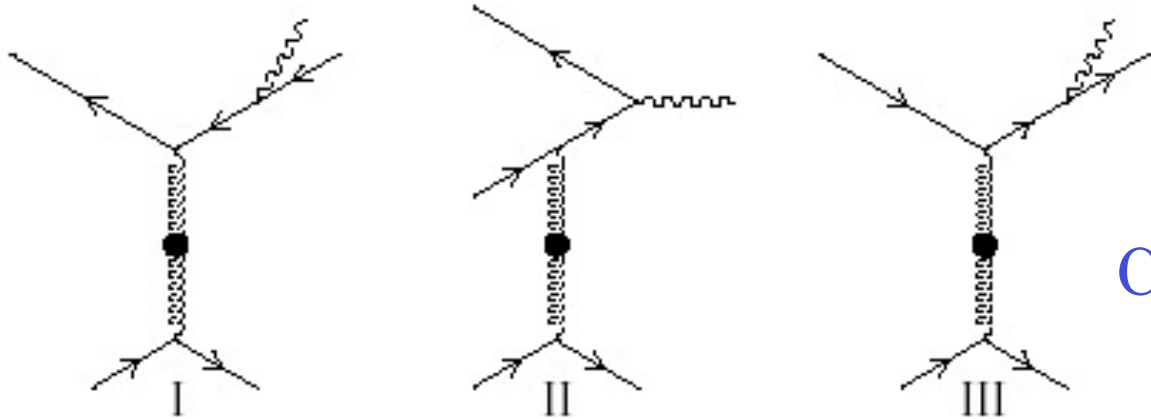
Thermal Photons from hot QCD: HTL program (Klimov (1981), Weldon (1982), Braaten & Pisarski (1990); Frenkel & Taylor (1990))



$$\text{Im } \Pi_{R\mu}^{\mu} \sim \ln \left(\frac{\varpi T}{(m_{th} (\sim gT))^2} \right)$$

Kapusta, Lichard, Seibert (1991)
Baier, Nakkagawa, Niegawa, Redlich (1992)

Going to two loops: Aurenche, Kobes, Gelis, Petitgirard (1996)
Aurenche, Gelis, Kobes, Zaraket (1998)



Co-linear singularities:

$$\alpha_s^2 \left(\frac{T^2}{m_{th}^2} \right) \sim \alpha_s$$

2001: Results complete at $O(\alpha_s)$

Arnold, Moore, and Yaffe JHEP **12**, 009 (2001); JHEP **11**, 057 (2001)
Incorporate LPM; Inclusive treatment of collinear enhancement,
photon and gluon emission

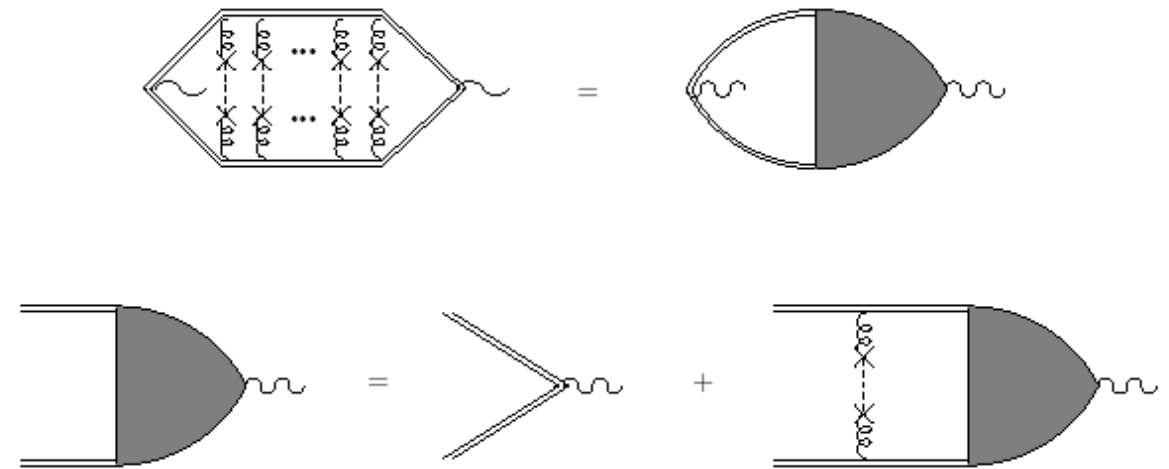


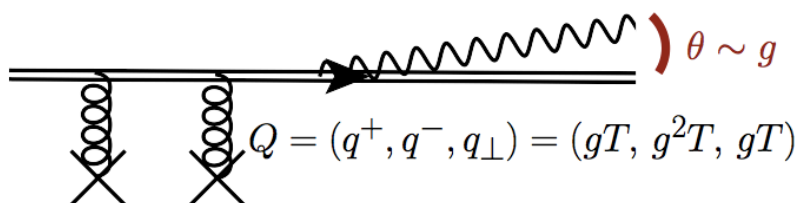
Are these rates final?

- Approach is LO, but

$$\alpha_s \sim 0.2 - 0.3$$

- Integral equation can be written in terms of a single-gluon scattering kernel:



NLO: 

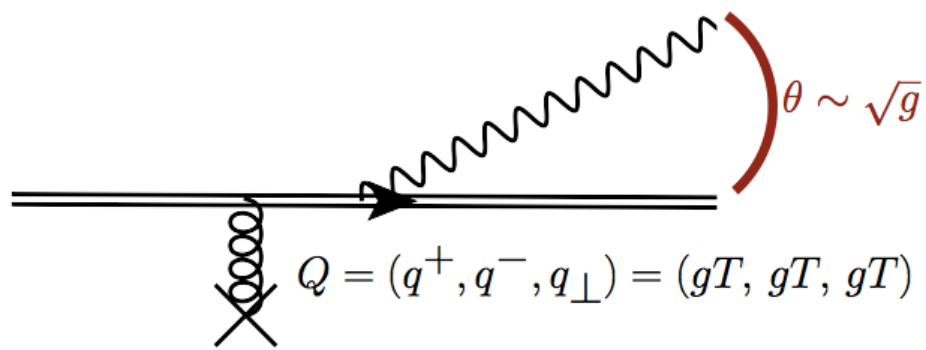
$$C(q_T)_{\text{LO}} = \frac{Tg^2 m_D}{q_T (q_T + m_D)} \Rightarrow \text{NLO}$$

Aurenche, Géelis, Zaraket (2002)

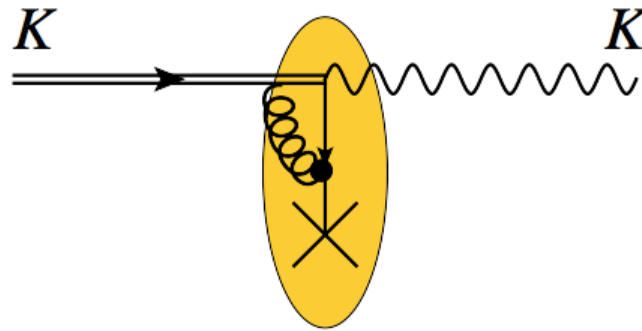
Simon Caron-Huot PRD (2010)



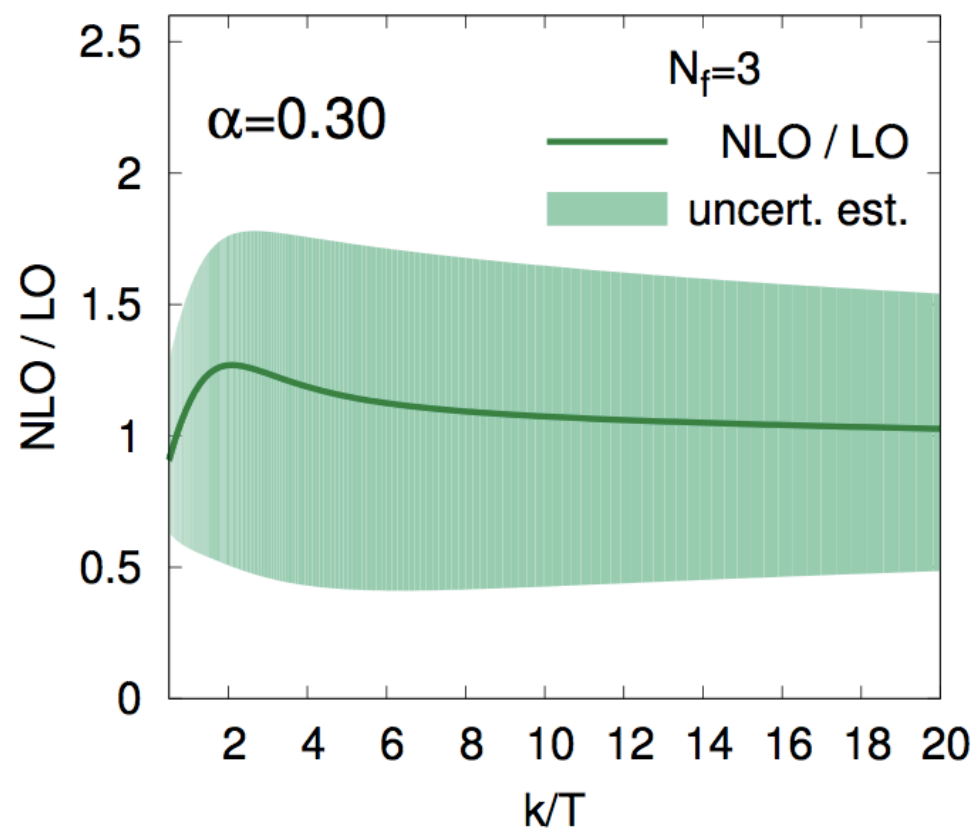
Larger angle emission



Conversion photons



$$R_{\text{NLO}} \sim g^3 \ln(1/g) + g^3$$

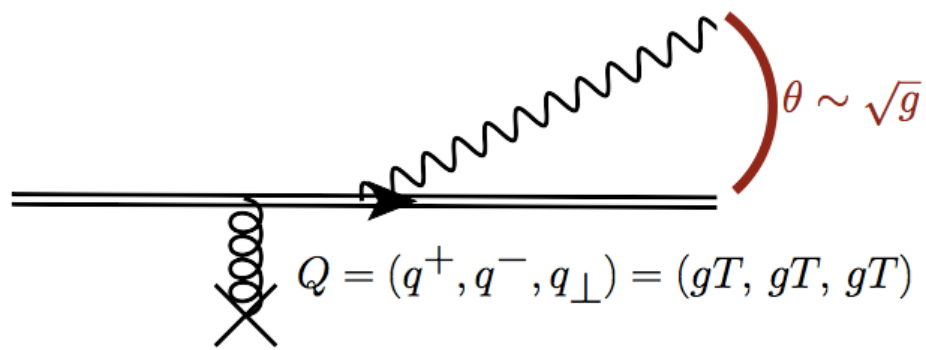


- Net correction to emission rate not numerically important in region up to $k/T \sim 10$
- Techniques developed here will have other applications in FTFT

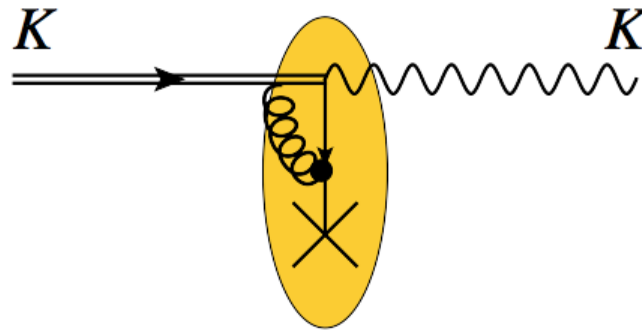
J.Ghiglieri, J.Hong, A.Kurkela, E.Lu, G.D.Moore, D.Teaney (2012)



Larger angle emission

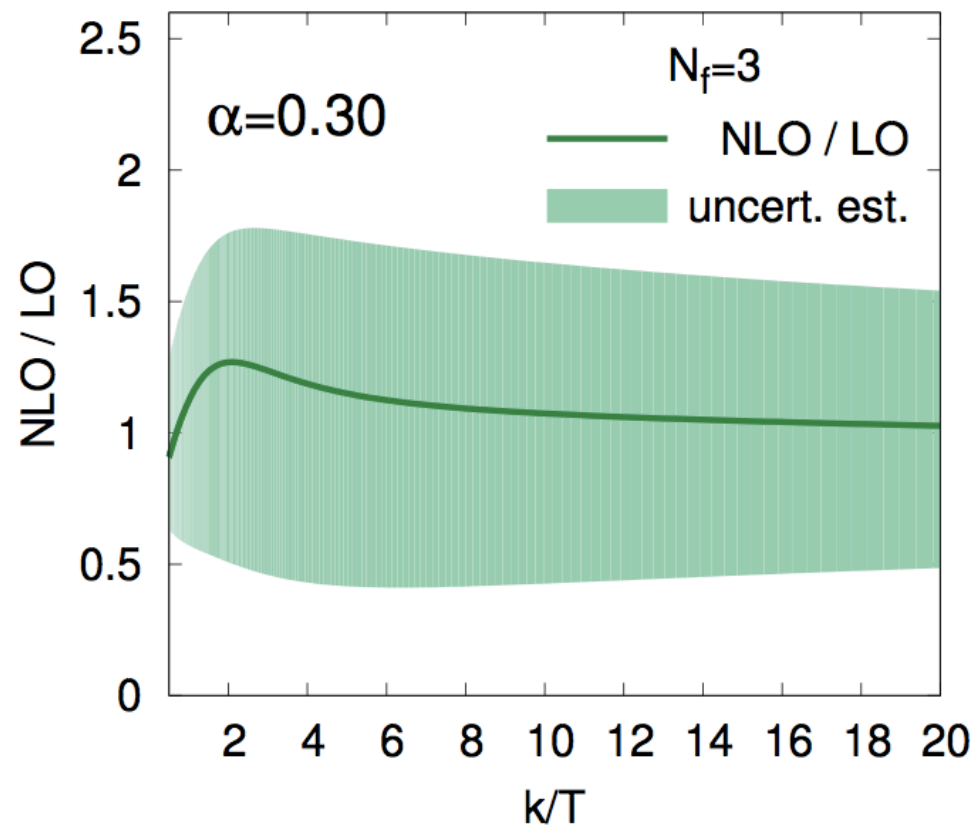


Conversion photons



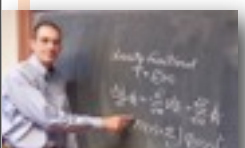
$$R_{\text{NLO}} \sim g^3 \ln(1/g) + g^3$$

PROGRESS



- Net correction to emission rate not numerically important in region up to $k/T \sim 10$
- Techniques developed here will have other applications in FTFT

J.Ghiglieri, J.Hong, A.Kurkela, E.Lu, G.D.Moore, D.Teaney (2012)



ELECTROMAGNETIC RADIATION FROM HADRONS

Chiral, Massive Yang-Mills:

O. Kaymakcalan, S. Rajeev, J. Schechter, PRD 30, 594 (1984)

$$\begin{aligned} \mathcal{L} = & \frac{1}{8} F_\pi^2 \text{Tr} D_\mu U D^\mu U^\dagger + \frac{1}{8} F_\pi^2 \text{Tr} M (U + U^\dagger) \\ & - \frac{1}{2} \text{Tr} (F_{\mu\nu}^L F^{L\mu\nu} + F_{\mu\nu}^R F^{R\mu\nu}) + m_0^2 \text{Tr} (A_\mu^L A^{L\mu} + A_\mu^R A^{R\mu}) \\ & + \text{non-minimal terms} \end{aligned}$$

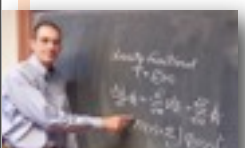
Parameters and form factors are constrained by hadronic phenomenology:

- Masses & strong decay widths
- Electromagnetic decay widths
- Other hadronic observables:

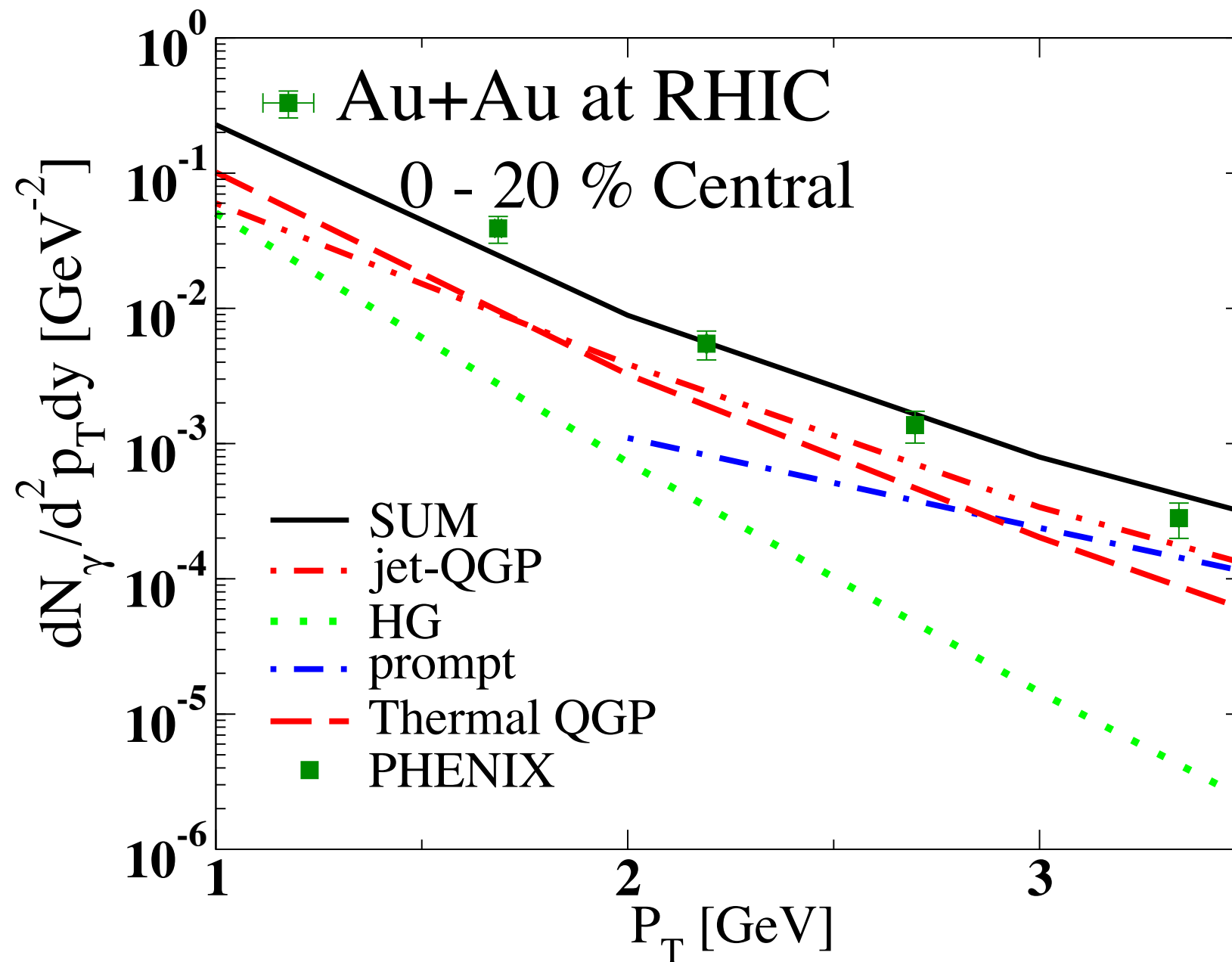
• *e.g.* $a_1 \rightarrow \pi \rho$ D/S (See also, Lichard and Vojik, Nucl. Phys. (2010); Lichard and Juran, PRD (2008))

EM emissivities computed: Turbide, Rapp, Gale, PRC (2004);
Turbide, McGill PhD (2006)

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APPLYING THIS TO INTERPRET PHOTONS MEASURED @ RHIC: RATES ARE INTEGRATED USING RELATIVISTIC HYDRODYNAMIC MODELING

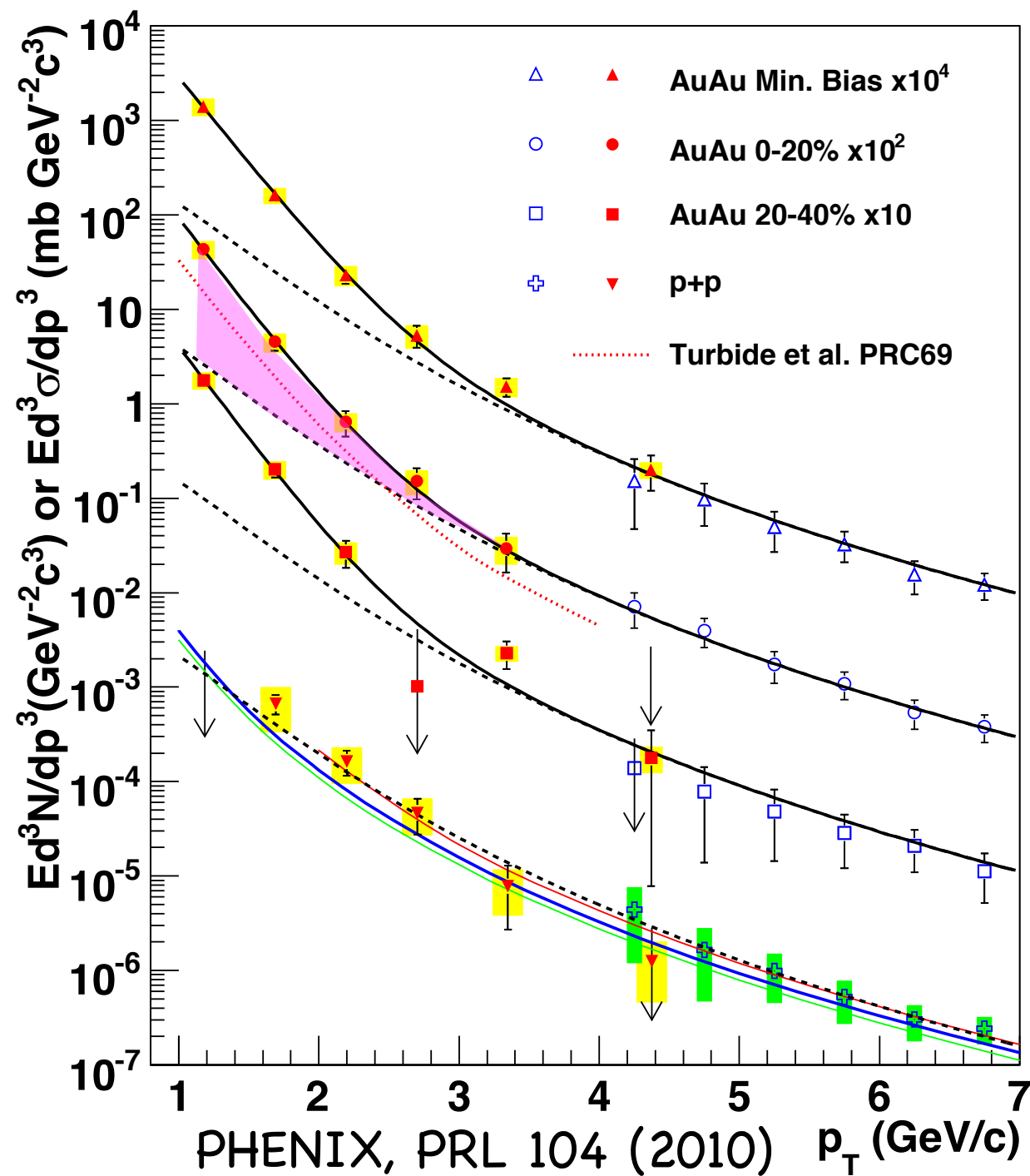


- At low p_T , spectrum dominated by thermal components (HG, QGP)
- At high p_T , spectrum dominated by pQCD
- Window for jet-QPG contributions at mid- p_T

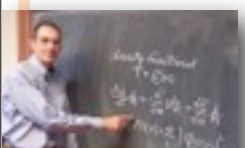
Turbide, Gale, Frodermann, Heinz, PRC (2008);
Higher p_T : G. Qin et al., PRC (2009)



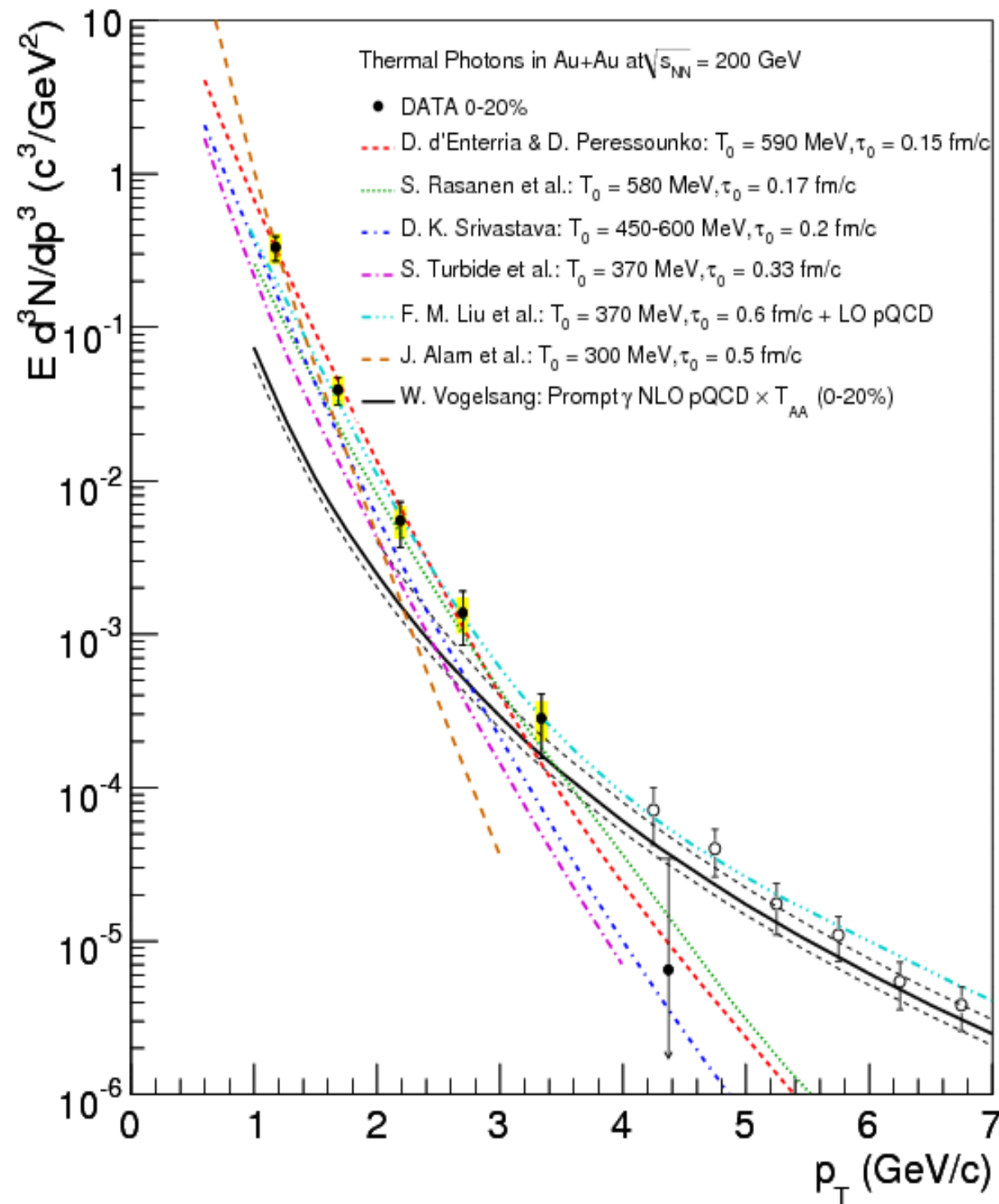
ONE OF THE USES OF PHOTONS: CHARACTERIZING THE HOT MATTER CREATED AT RHIC



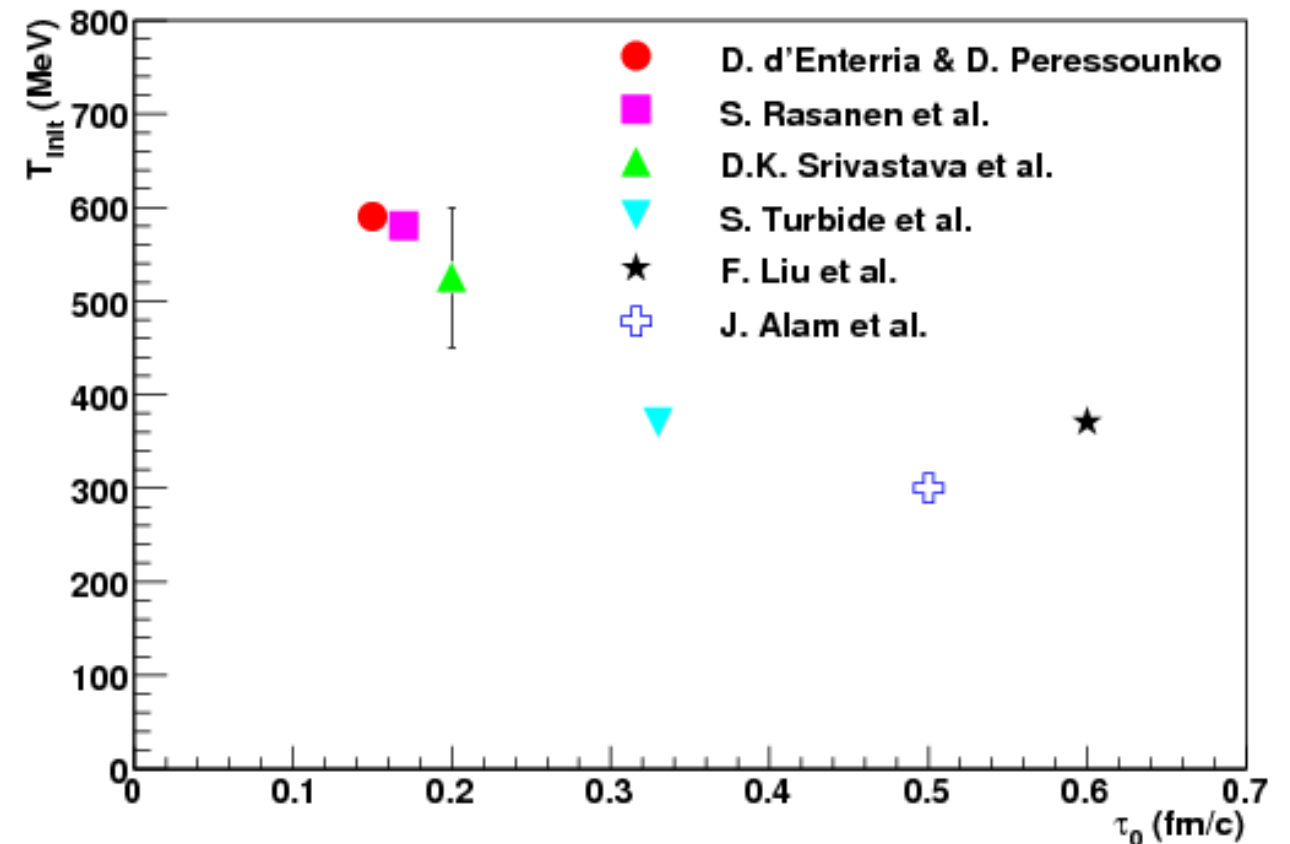
$$T_{\text{excess}} = 221 \pm 19 \pm 19 \text{ MeV}$$



ONE OF THE USES OF PHOTONS: CHARACTERIZING THE HOT MATTER CREATED AT RHIC



$$T_{\text{excess}} = 221 \pm 19 \pm 19 \text{ MeV}$$



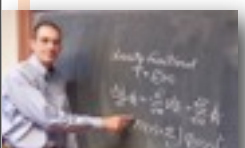
$$T_{\text{ini}} = 300 \text{ to } 600 \text{ MeV}$$

$$\tau_0 = 0.15 \text{ to } 0.5 \text{ fm/c}$$

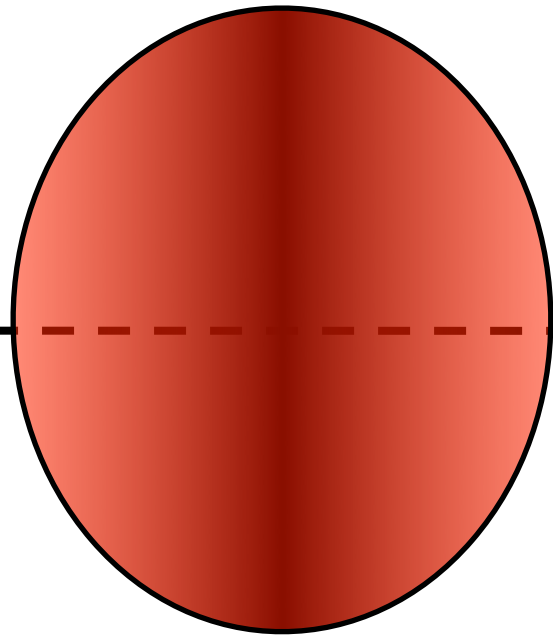
D'Enterria & Peressounko, Eur. Phys. J. (2006)

Knowing rates alone is not enough to guarantee predictive power or even characterization ability

Charles Gale



BEYOND SIMPLE SPECTRA: FLOW AND CORRELATIONS



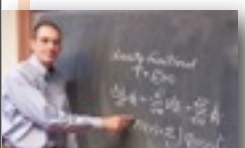
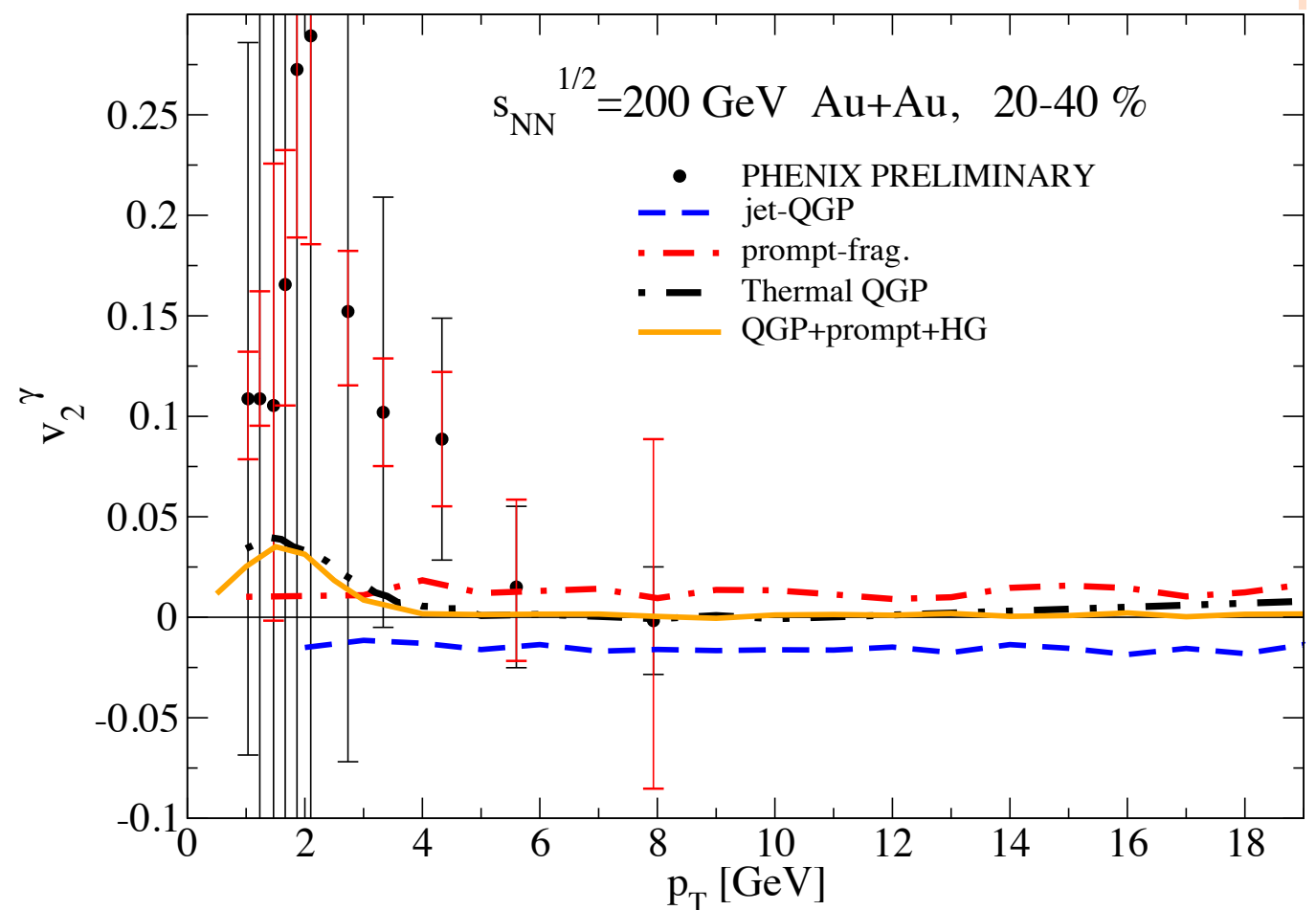
$$\frac{dN}{p_T dp_T d\phi} = \frac{dN}{2\pi p_T dp_T} \left[1 + \sum_n 2v_n \cos(n\phi) \right]$$

- Soft photons will go with the flow
- Jet-plasma photons: a negative v_2
- Details will matter: flow, $T(t)$. . .

Turbide, Gale, Fries PRL (2006)

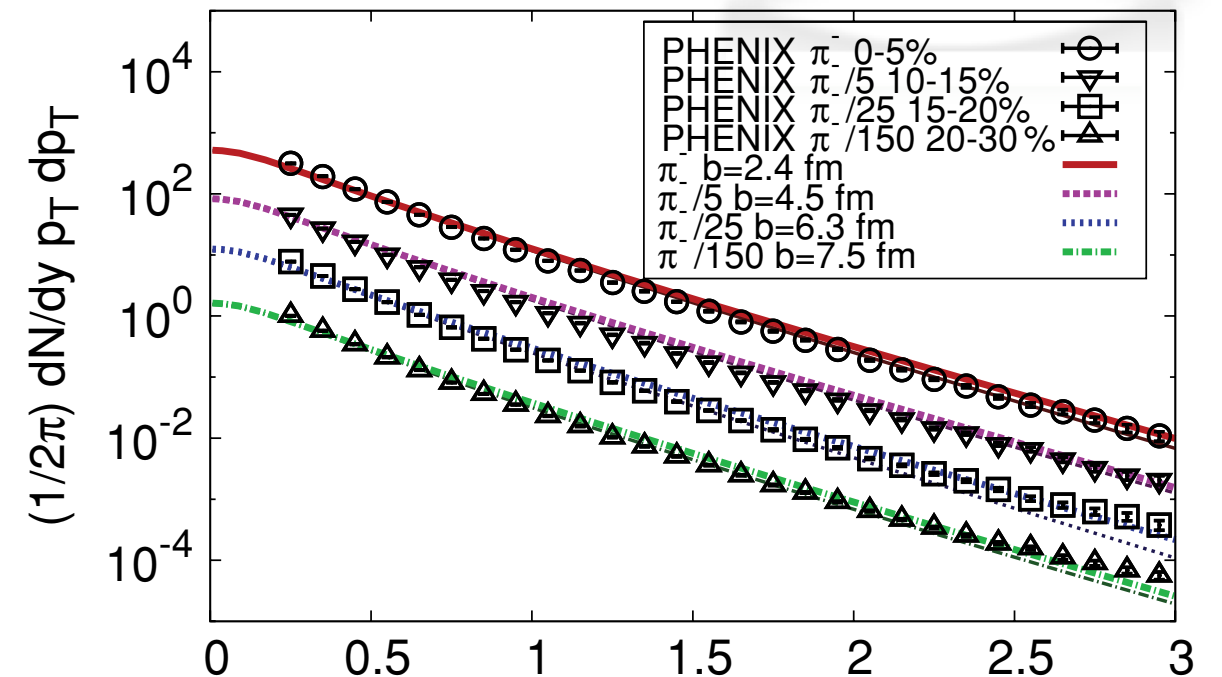
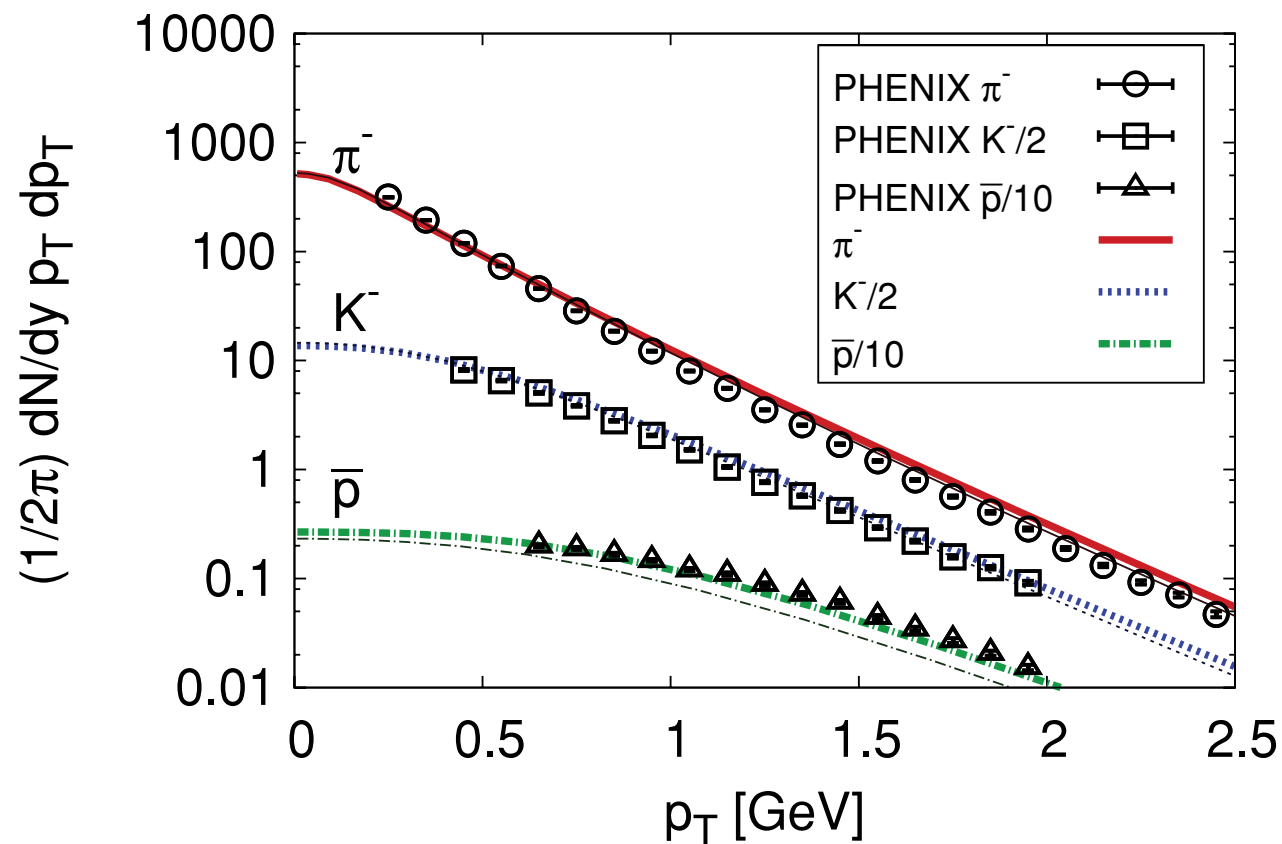
Low p_T : Chatterjee *et al.*, PRL (2006)

All p_T : Turbide *et al.*, PRC (2008)



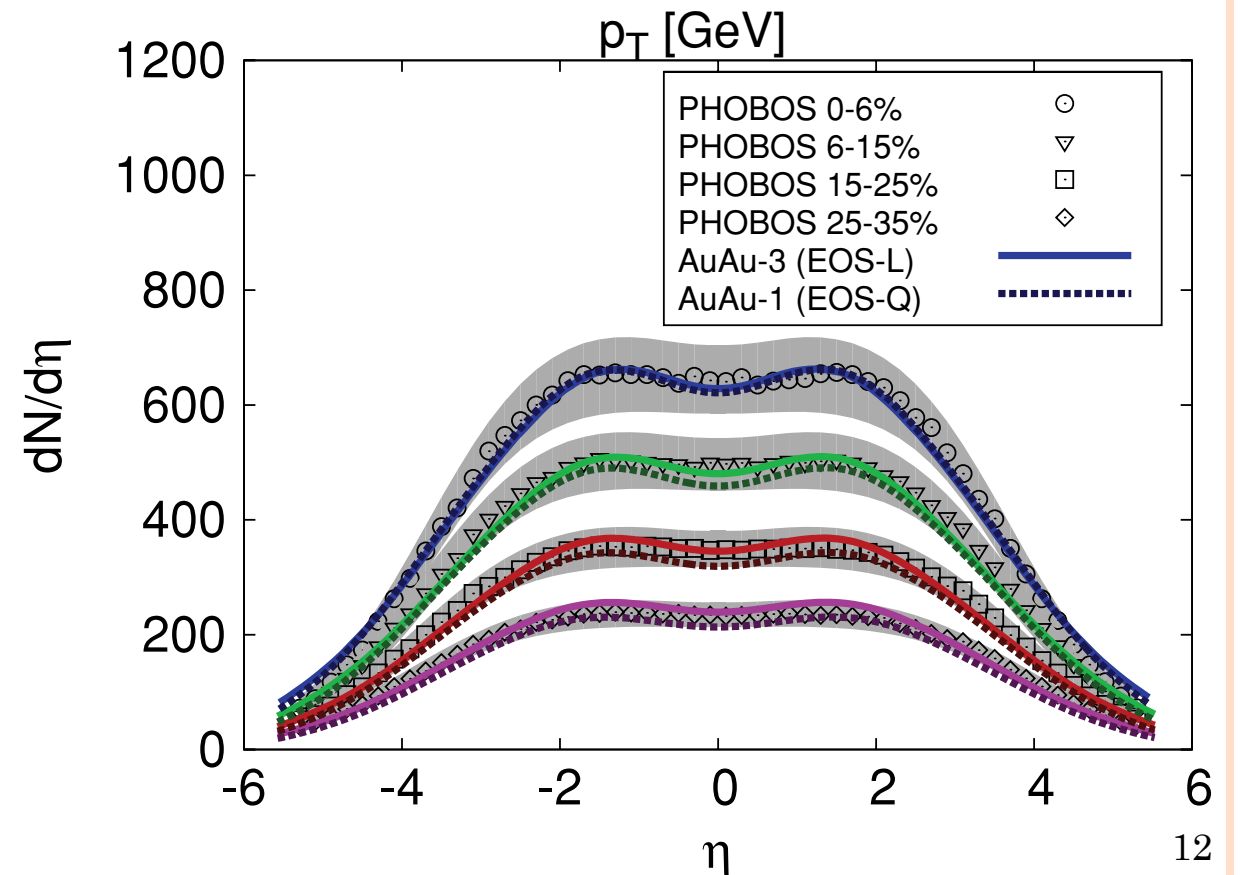
PROGRESS IN CHARACTERIZATION TOOL: 3D VISCOUS RELATIVISTIC HYDRODYNAMICS

MUSIC:



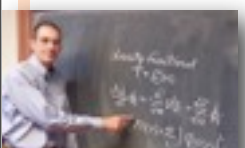
MUSIC: 3D relativistic hydro

- Ideal: Schenke, Jeon, and Gale, PRC (2010)
- **FIC and Viscous**: Schenke, Jeon, Gale, PRL (2011)



Viscosity effects on EM observables?

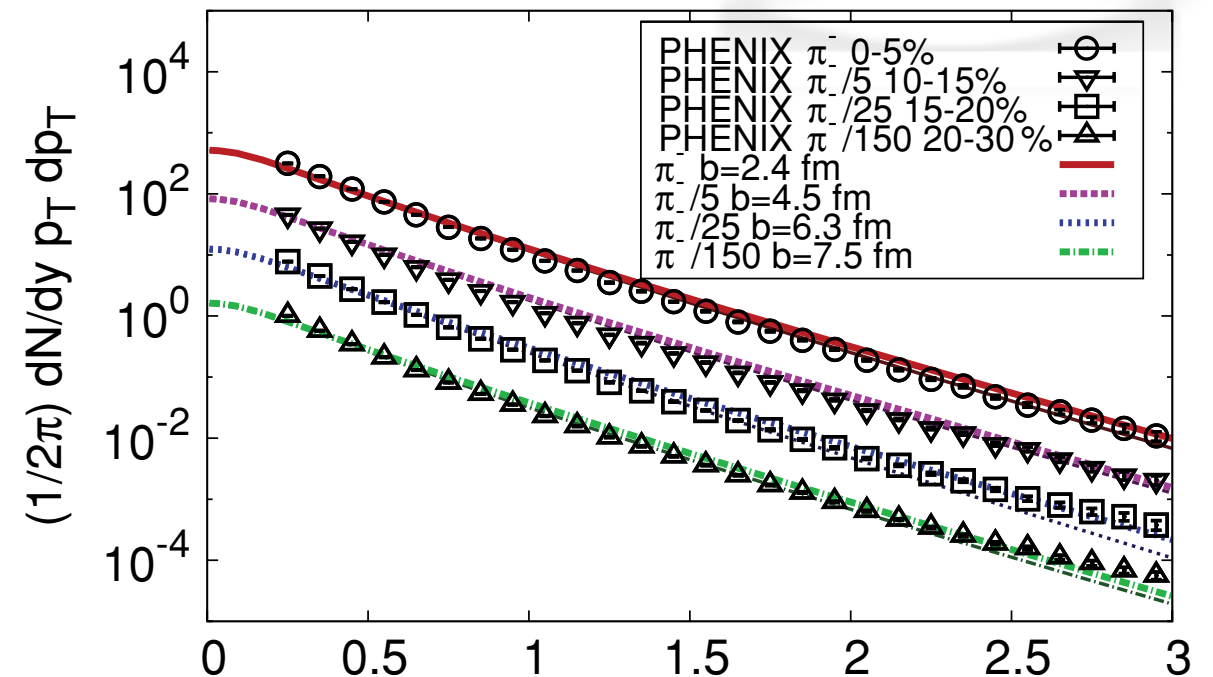
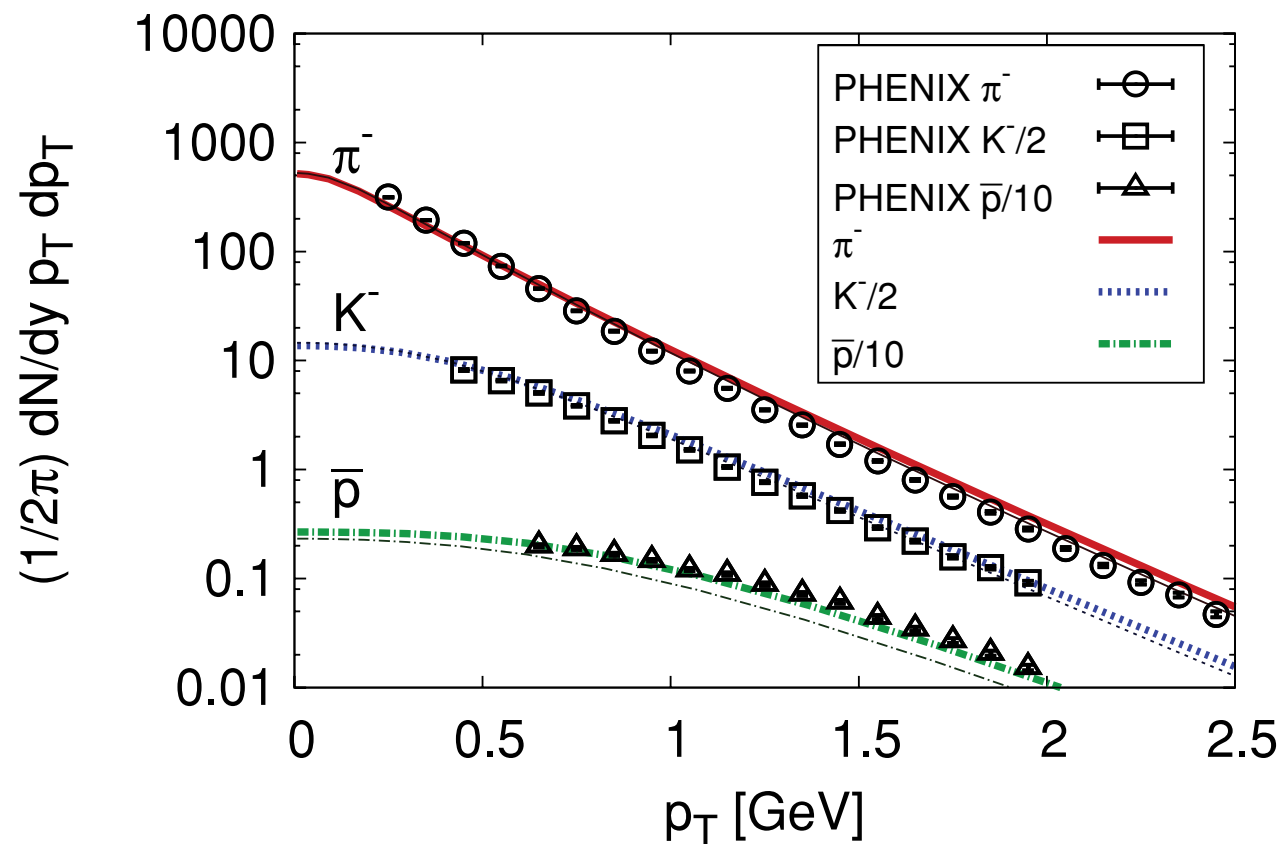
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PROGRESS

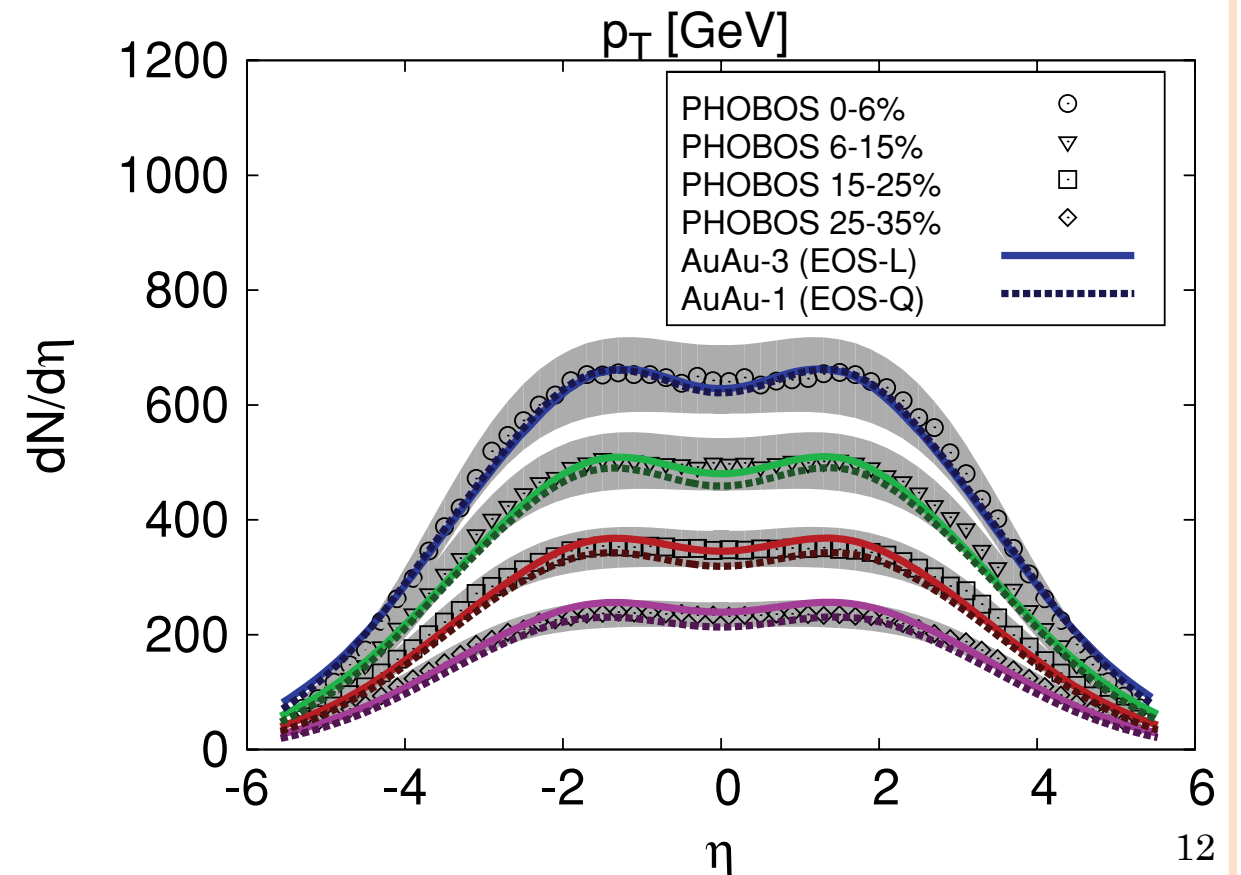
PROGRESS IN CHARACTERIZATION TOOL: 3D VISCOUS RELATIVISTIC HYDRODYNAMICS

MUSIC:



MUSIC: 3D relativistic hydro

- Ideal: Schenke, Jeon, and Gale, PRC (2010)
- **FIC and Viscous:** Schenke, Jeon, Gale, PRL (2011)



Viscosity effects on EM observables?

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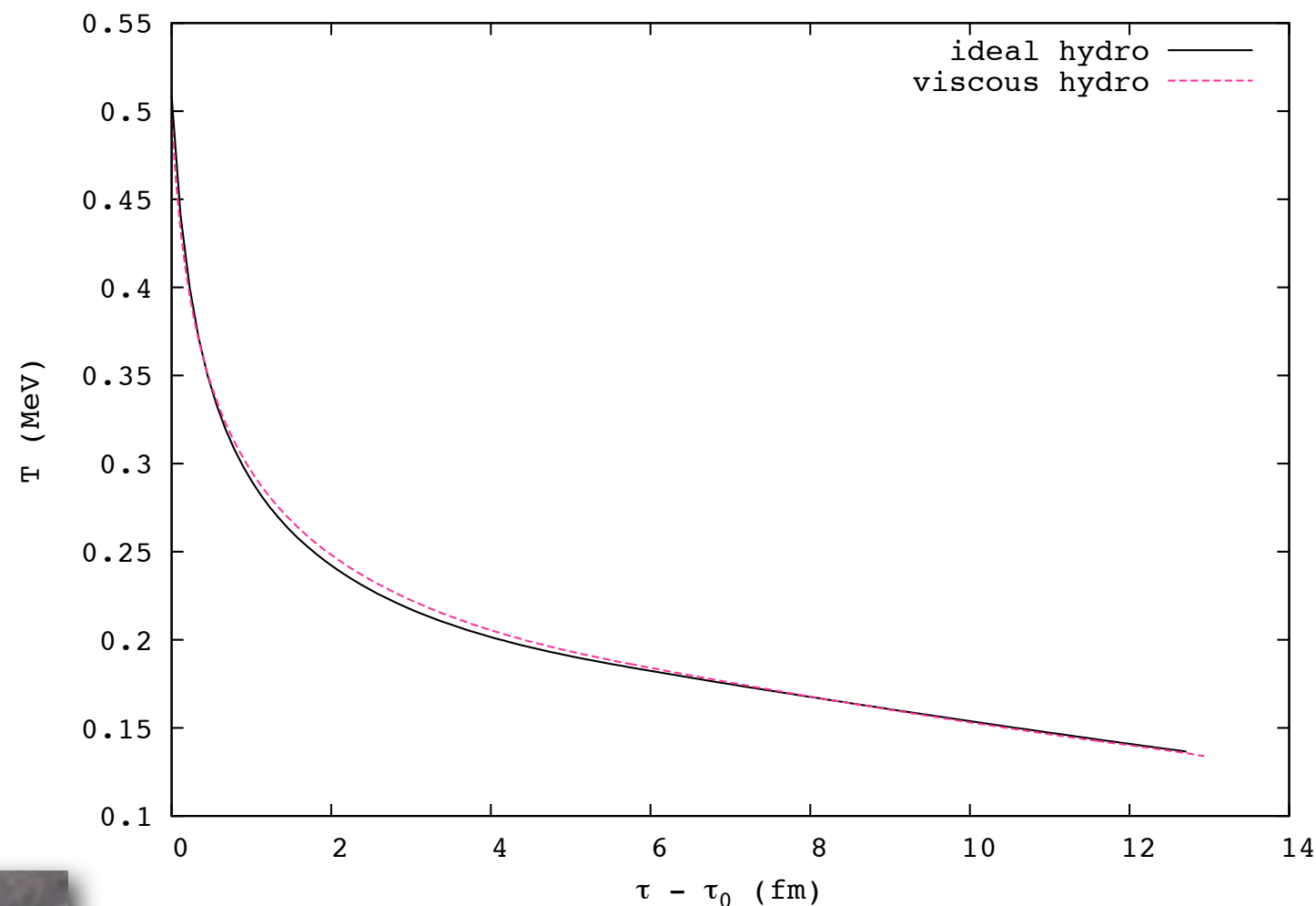
THE EFFECTS OF SHEAR VISCOSITY ON BULK DYNAMICS

$$T_{\text{ideal}}^{\mu\nu} = (\varepsilon + P)u^\mu u^\nu - Pg^{\mu\nu}$$

$$T^{\mu\nu} = T_{\text{ideal}}^{\mu\nu} + \pi^{\mu\nu}$$

Israël & Stewart, Ann. Phys. (1979), Baier et al., JHEP (2008), Luzum and Romatschke, PRC (2008)

$$\partial_\mu (su^\mu) \propto \eta$$



- Viscous evolution starts with a lower T
- T drop is slower than ideal case



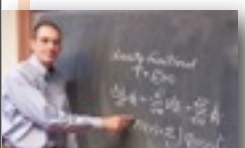
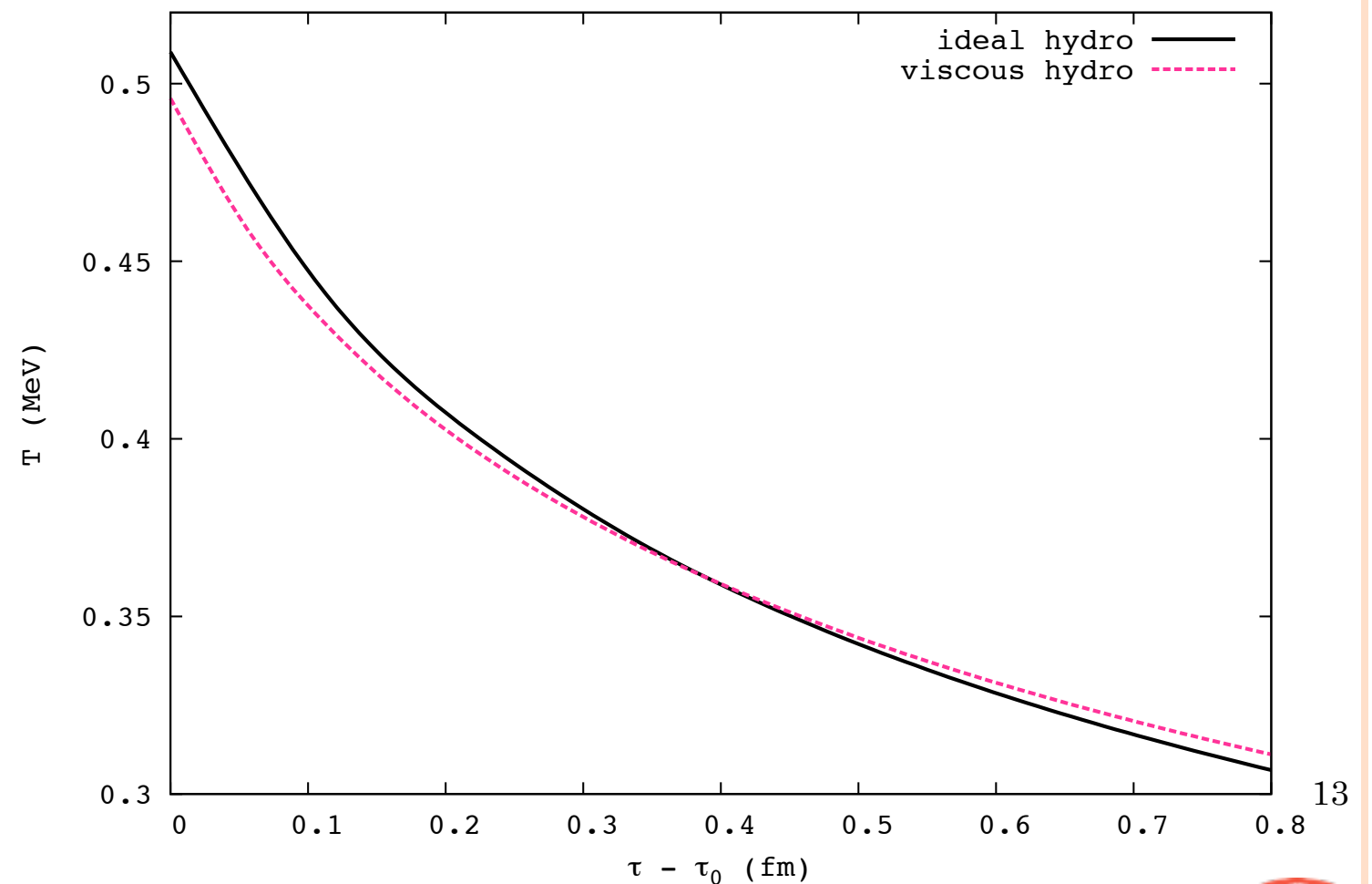
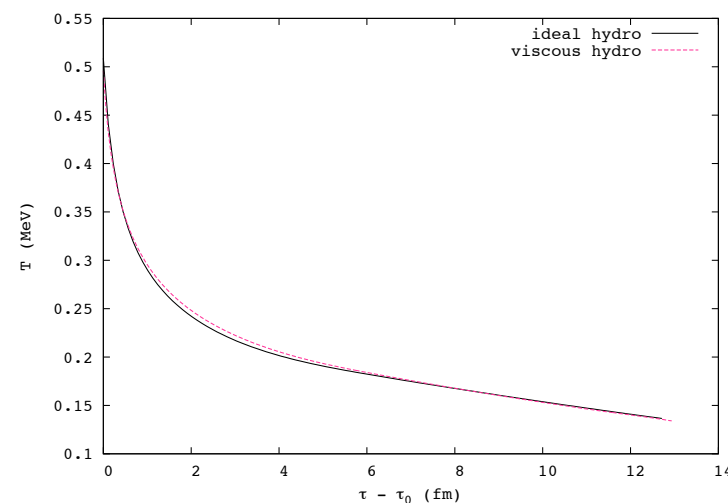
THE EFFECTS OF SHEAR VISCOSITY ON BULK DYNAMICS

$$T_{\text{ideal}}^{\mu\nu} = (\varepsilon + P)u^\mu u^\nu - Pg^{\mu\nu}$$

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Israël & Stewart, Ann. Phys. (1979), Baier et al., JHEP (2008), Luzum and Romatschke, PRC (2008)

$$\partial_\mu (su^\mu) \propto \eta$$



THE EFFECTS OF SHEAR VISCOSITY ON THE PHOTON DISTRIBUTION

In-medium **hadrons**:

$$f_0(u^\mu p_\mu) = \frac{1}{(2\pi)^3} \frac{1}{\exp[(u^\mu p_\mu - \mu)/T] \pm 1}$$

$$f \rightarrow f_0 + \delta f, \quad \delta f = f_0(1 \pm (2\pi)^3 f_0) p^\alpha p^\beta \pi_{\alpha\beta} \frac{1}{2(\varepsilon + P)T^2}$$

$$q_0 \frac{d^3 R}{d^3 q} = \int \frac{d^3 p_1}{2(2\pi)^3 E_1} \frac{d^3 p_2}{2(2\pi)^3 E_2} \frac{d^3 p_3}{2(2\pi)^3 E_3} (2\pi)^4 |M|^2 \delta^4(\dots) \frac{f(E_1)f(E_2)[1 \pm f(E_3)]}{2(2\pi)^3}$$

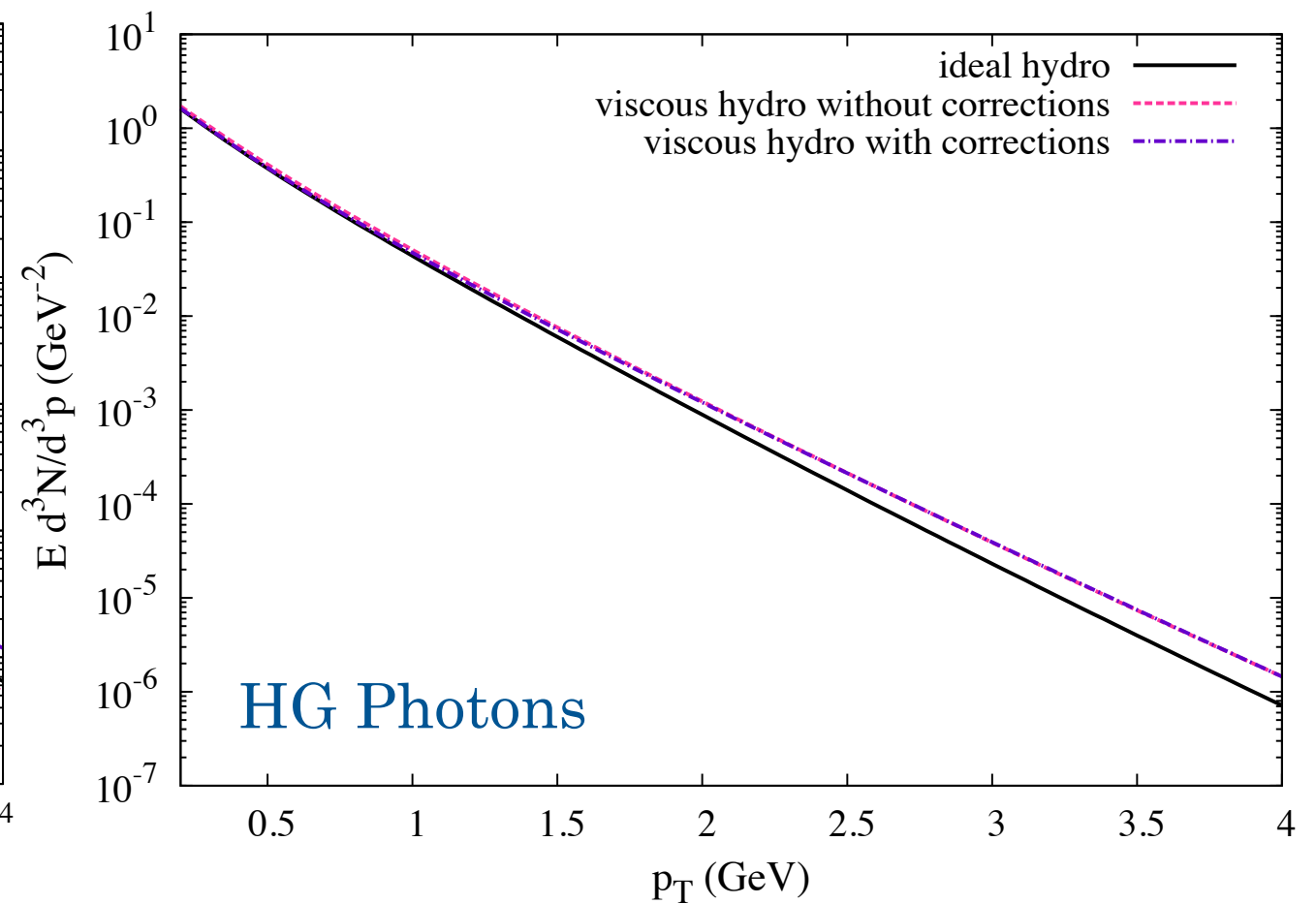
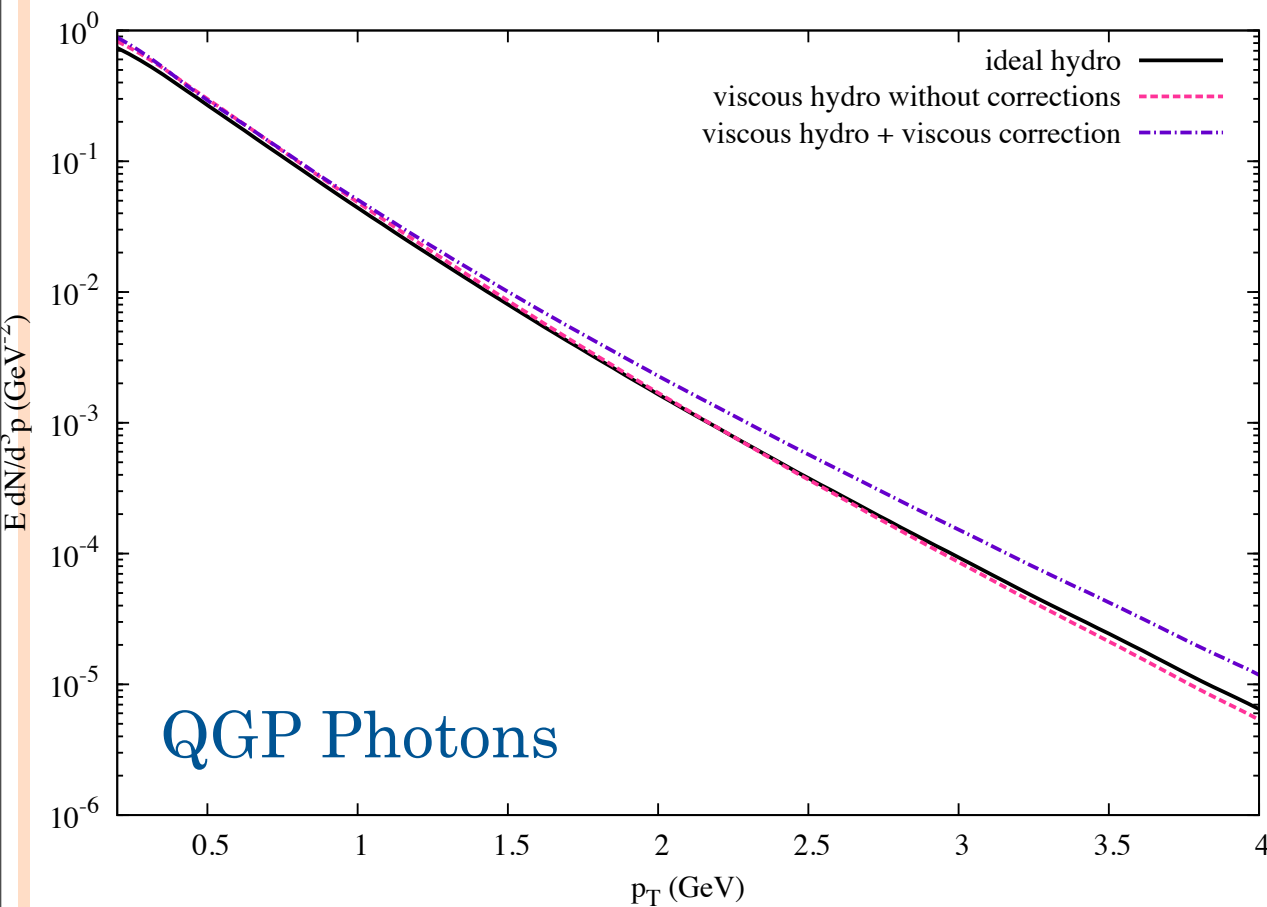
One considers all the reaction and radiative decay channels of external state combinations of:

$\{\pi, K, \rho, K^*, a_1\}$ With hadronic form factors

+ QGP Photons



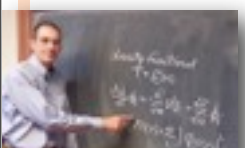
THE EFFECTS OF SHEAR VISCOSITY ON THE PHOTON DISTRIBUTION



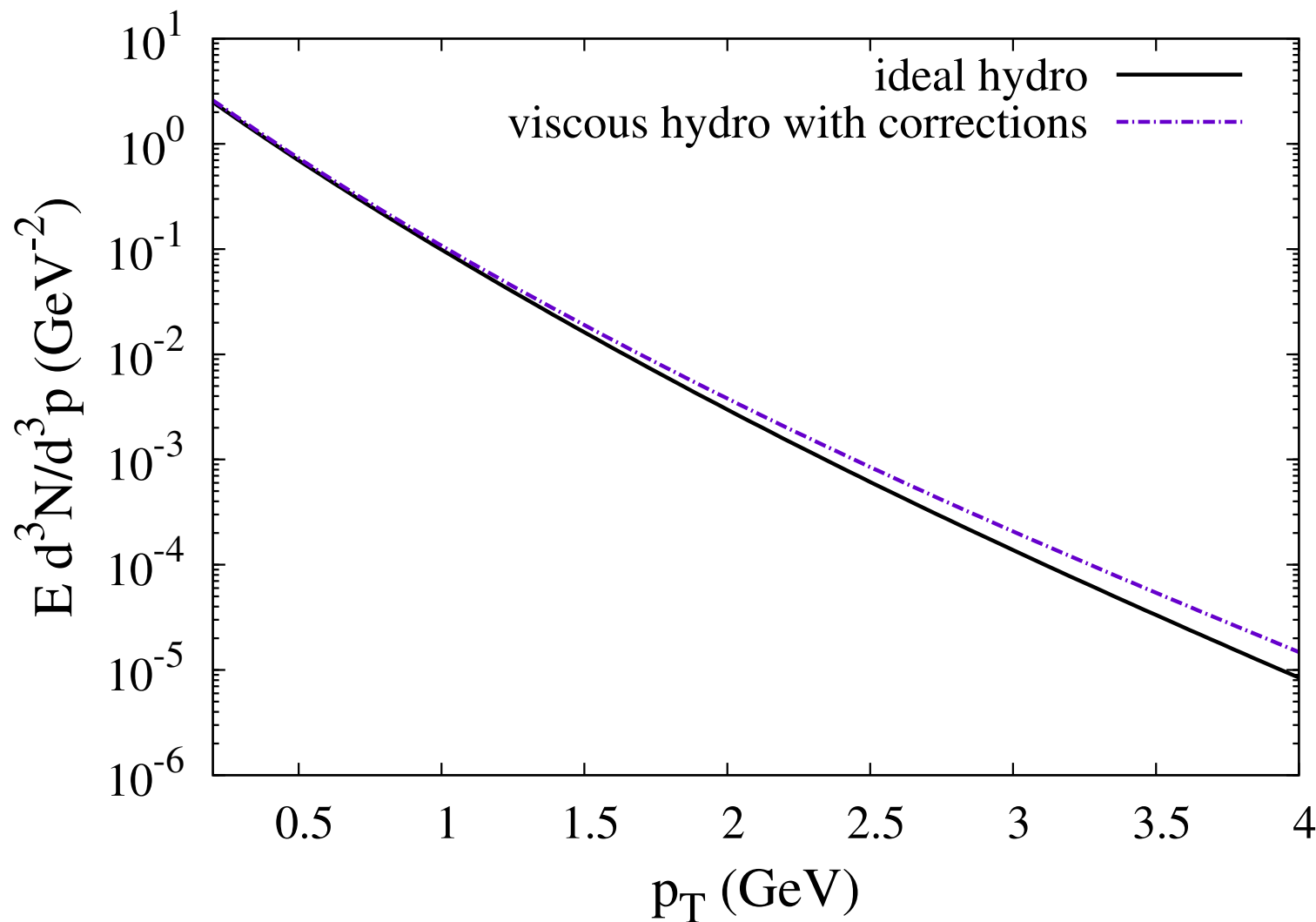
K. Dusling NPA (2010)
Chaudhuri & Sinha, PRC (2011)

Viscous effects harden
the photon spectrum

M. Dion et al., PRC (2011)



THE NET THERMAL PHOTON YIELD

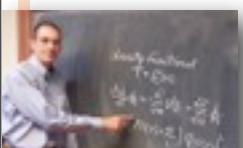
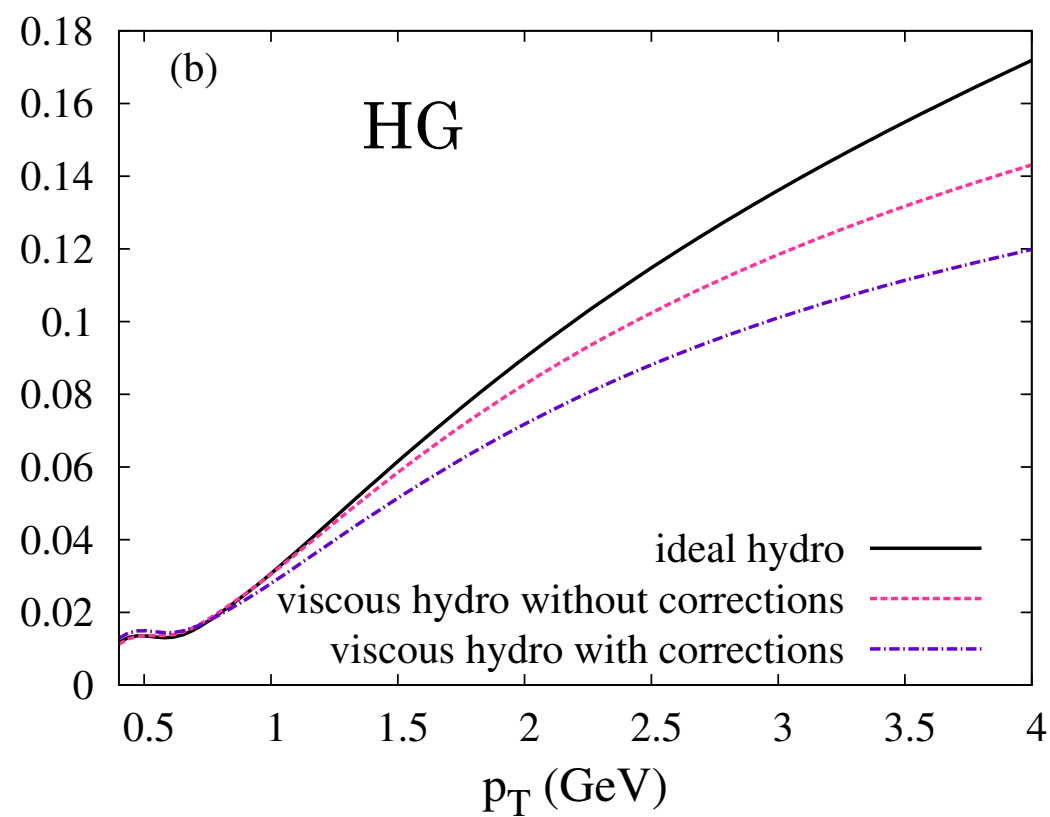
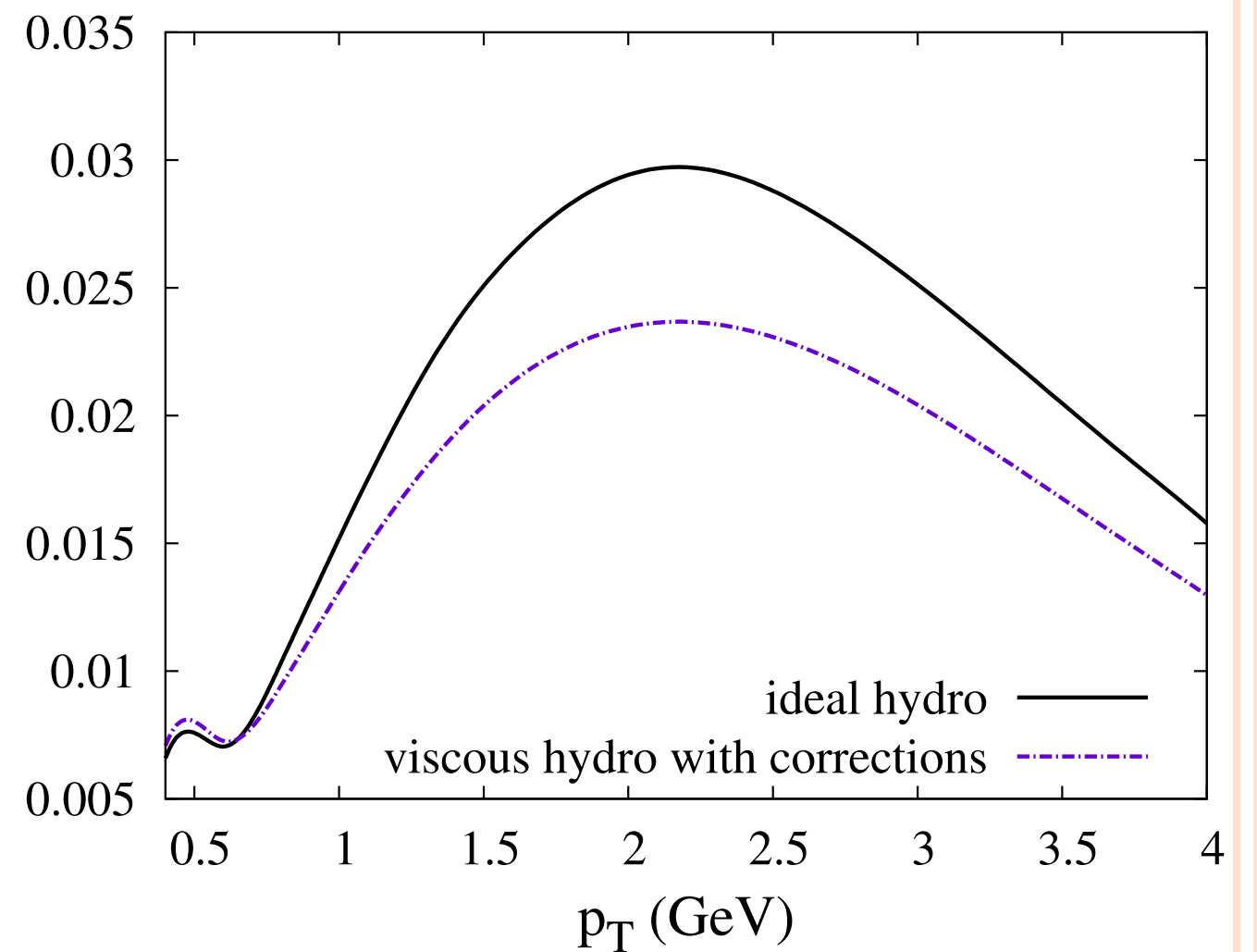
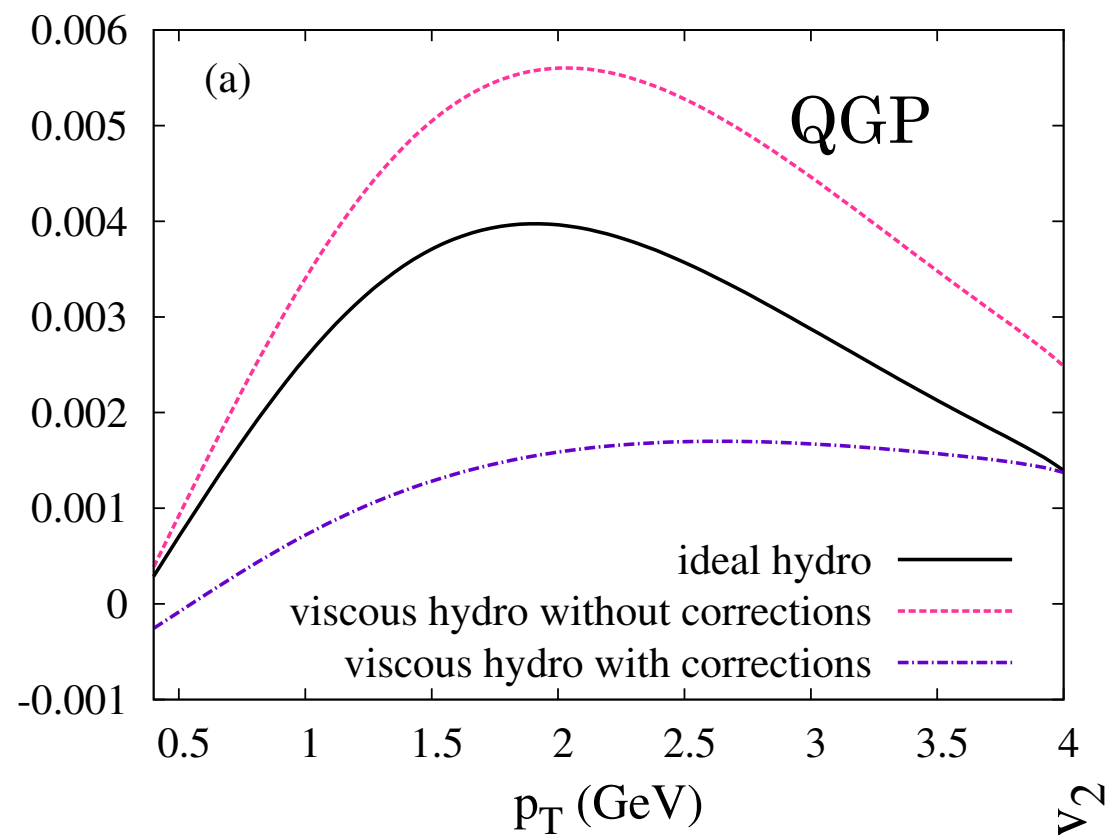


PROGRESS

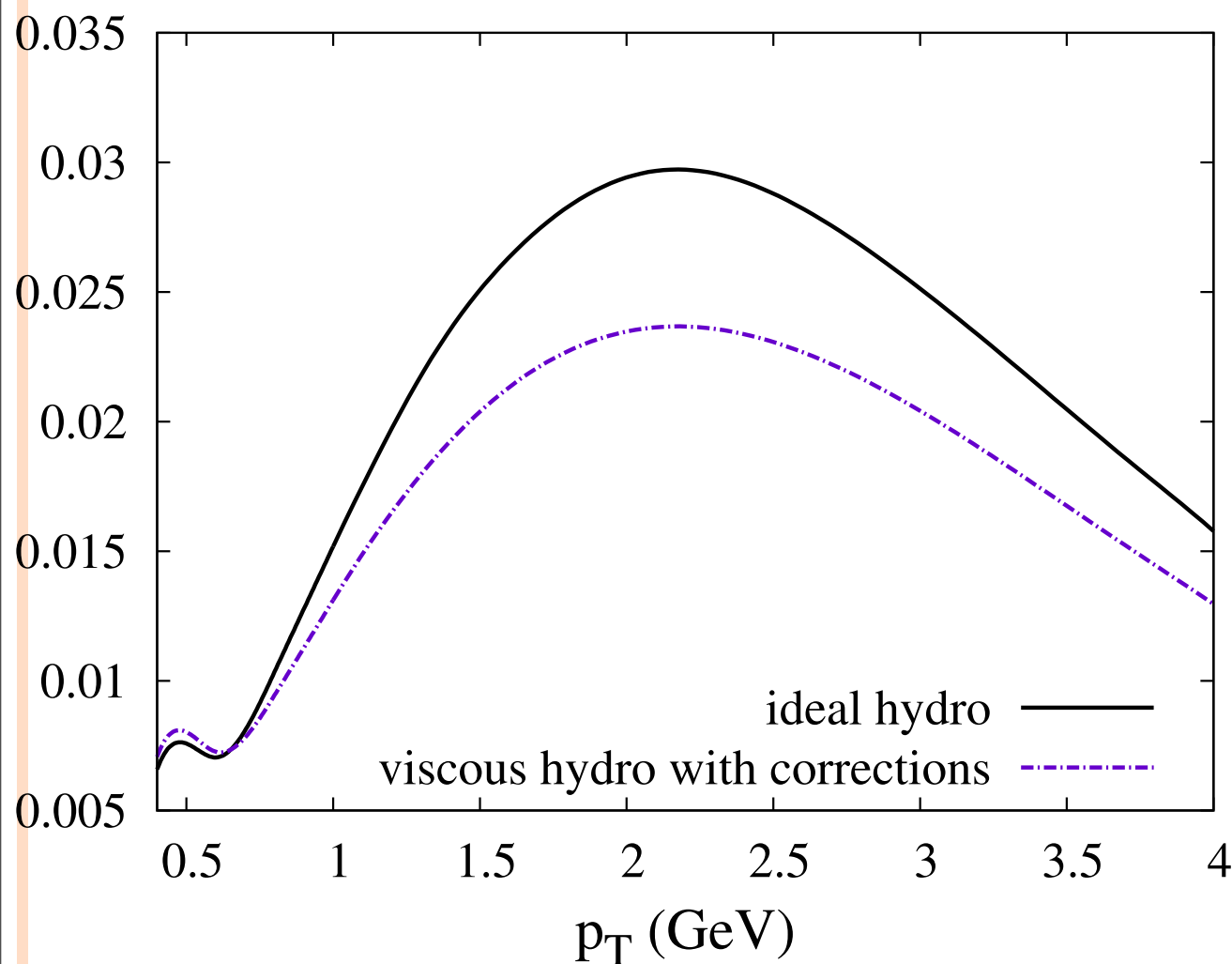
- Viscous corrections make the spectrum harder, $\approx 100\%$ at $p_T = 4 \text{ GeV}$.
- Increase in the slope of $\approx 15\%$ at $p_T = 2 \text{ GeV}$.
- Extracting the viscosity from the photon spectra will be challenging
- Once pQCD photons are included: a few % effect from viscosity
- More work is still needed to properly include all photon sources in a consistent way



SHEAR VISCOSITY AND PHOTON V_2



SHEAR VISCOSITY AND PHOTON v_2

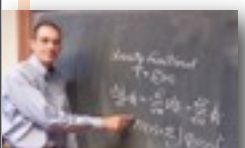


- The net elliptic flow is a *weighted average*. A larger QGP yield will yield a smaller v_2 . Same story - *mutatis mutandis* - for the HG
- The turnover at $p_T \approx 2$ GeV is QGP-driven (*)
- The net effect of viscous corrections makes the photon elliptic flow smaller, as it does for hadrons

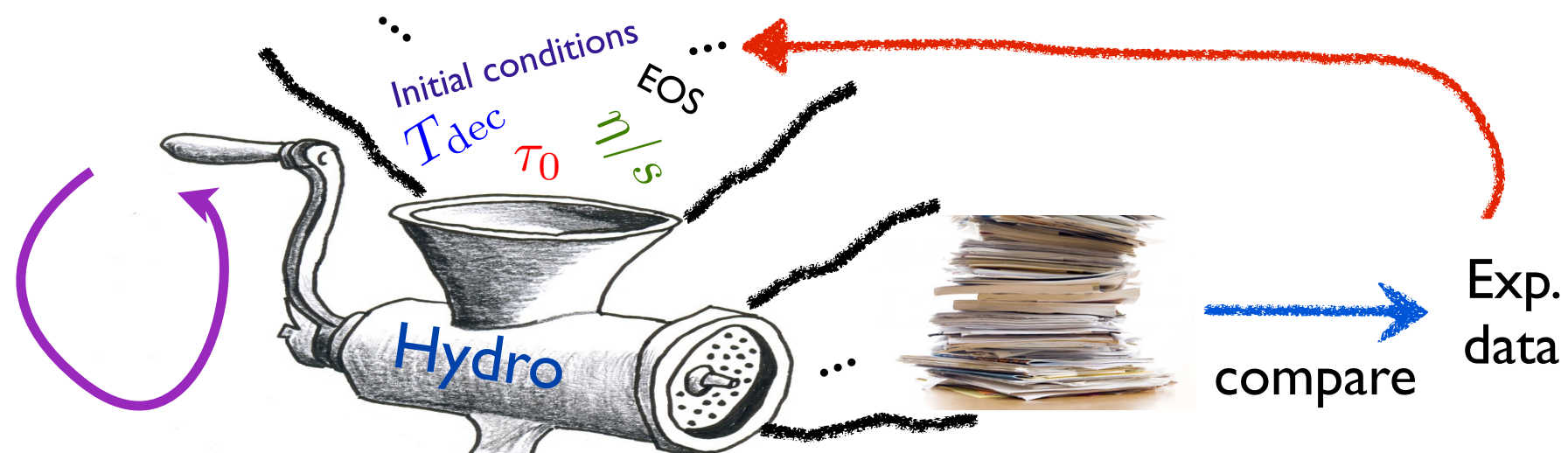
PROGRESS

M. Dion *et al.*, PRC 2011

Charles Gale



LOOKING UNDER THE HOOD: NON-EQUILIBRIUM EFFECTS



Chun Shen (OSU)

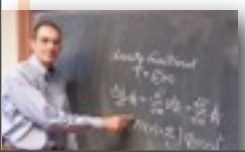
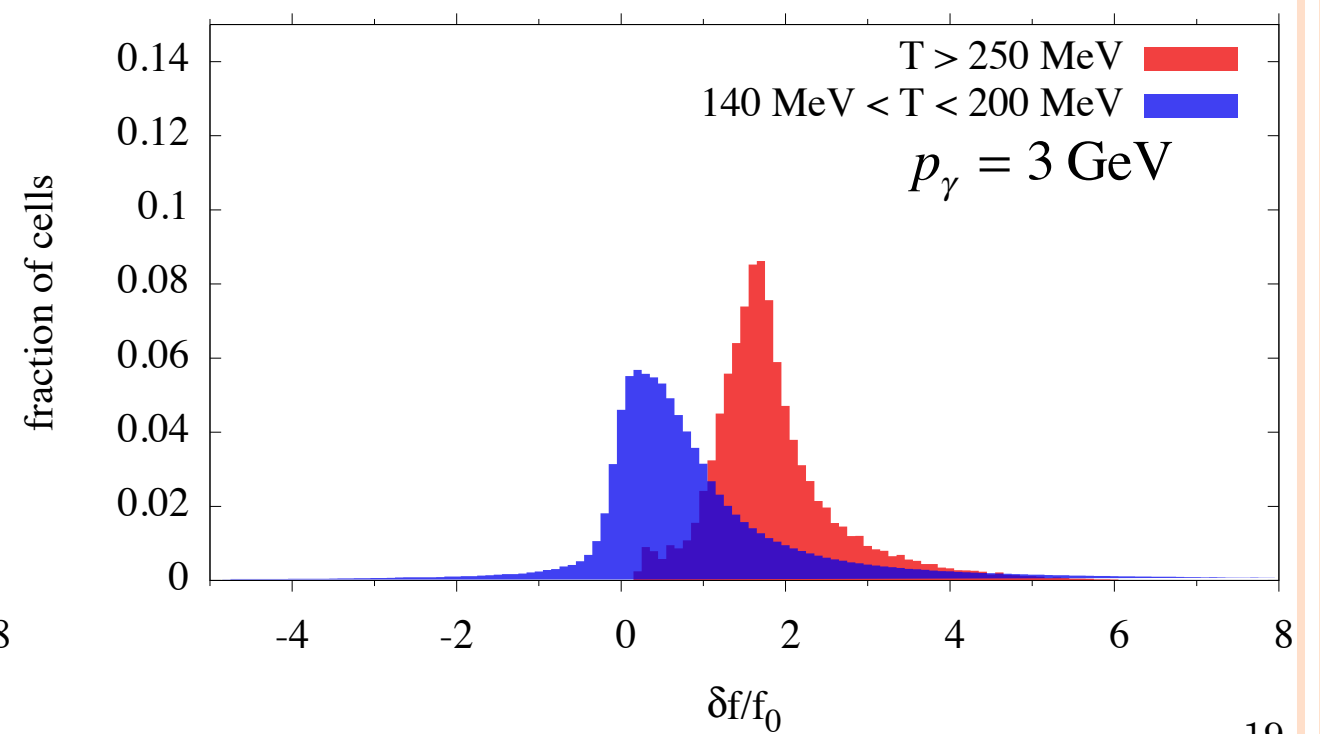
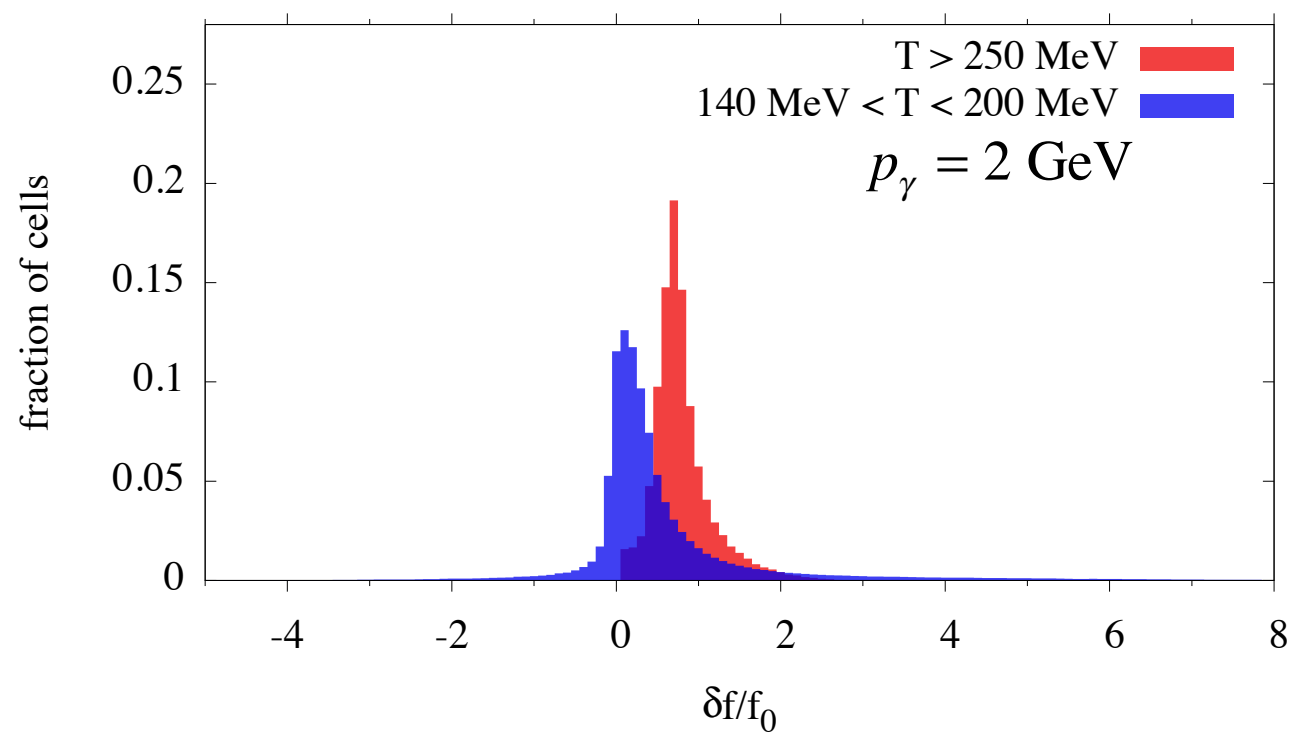
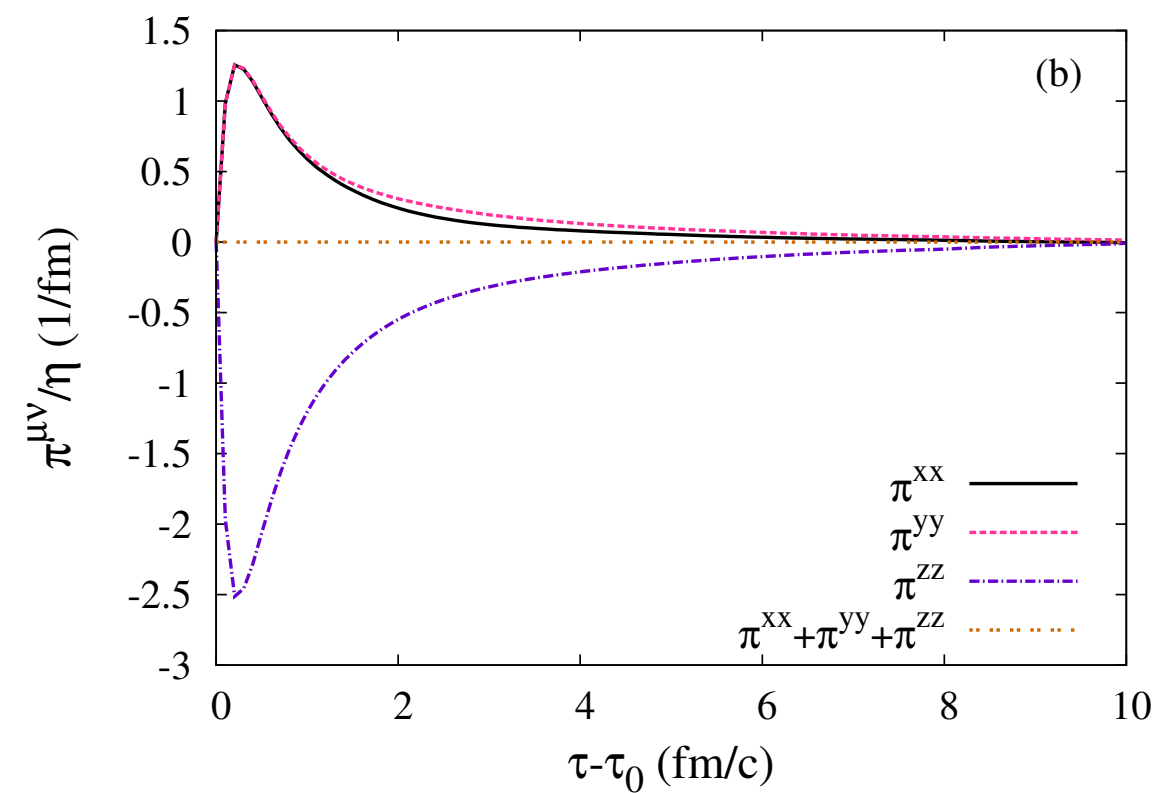
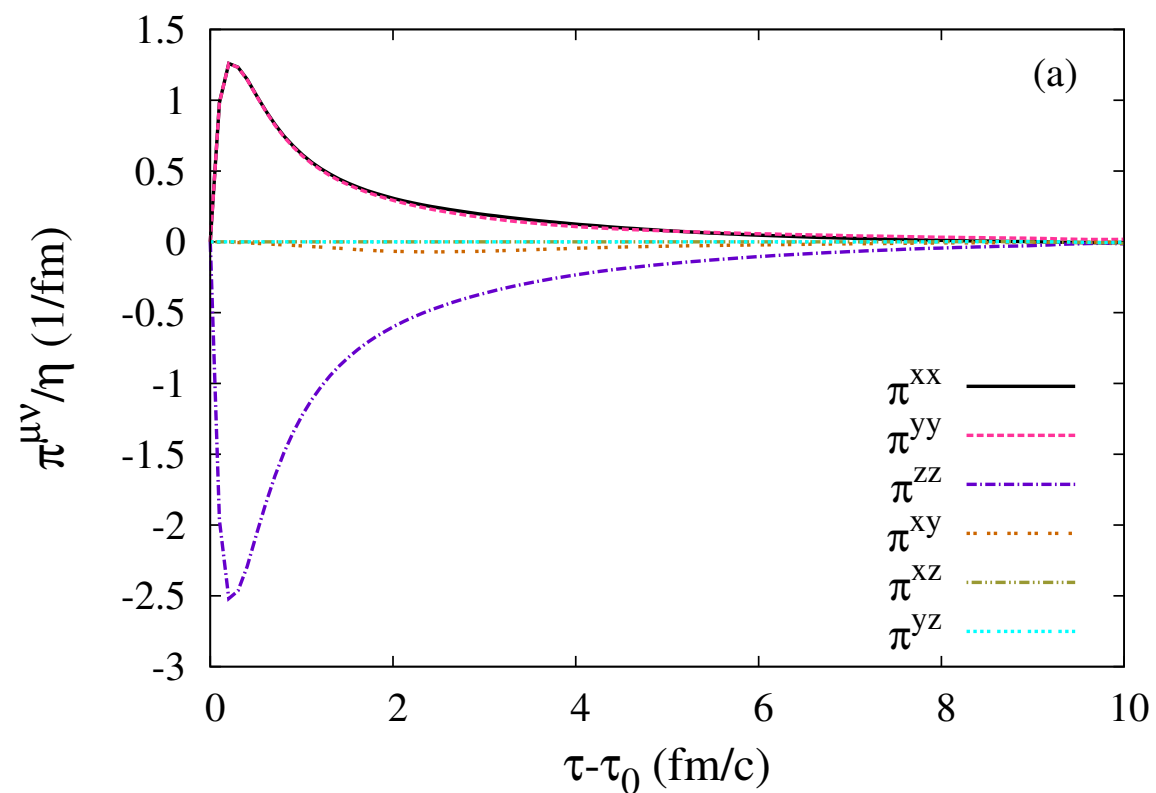
18



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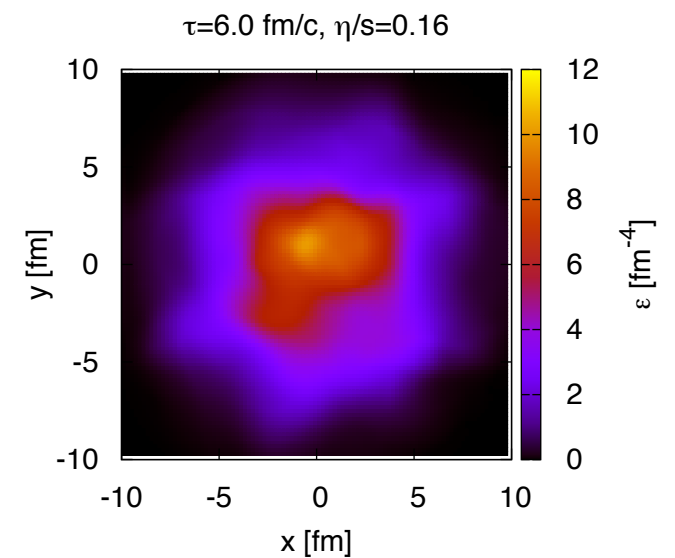
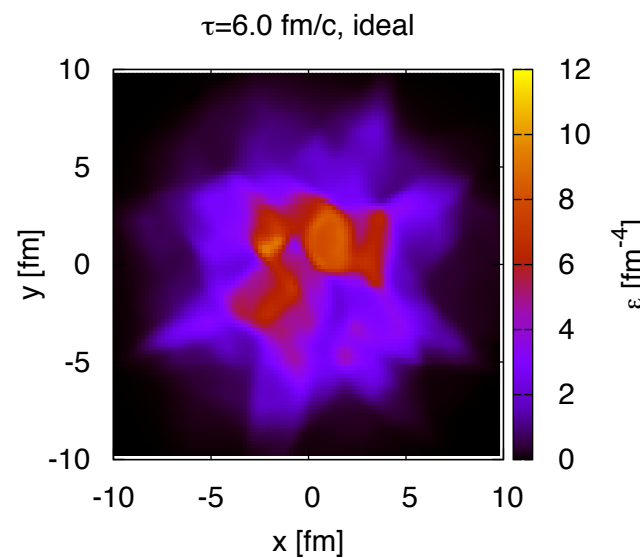
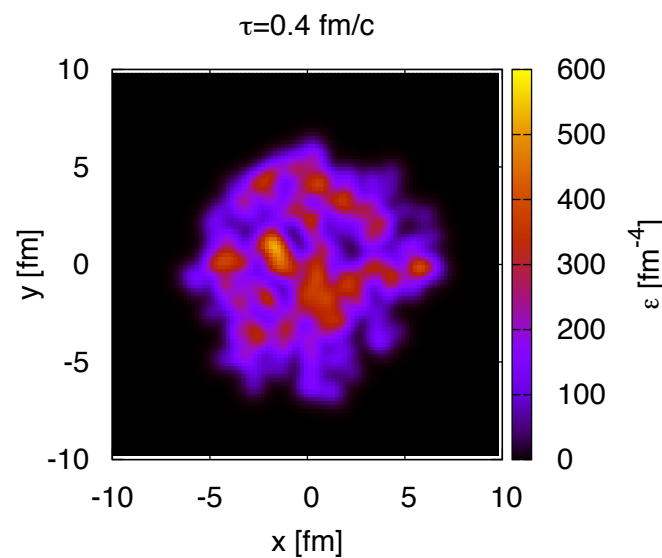
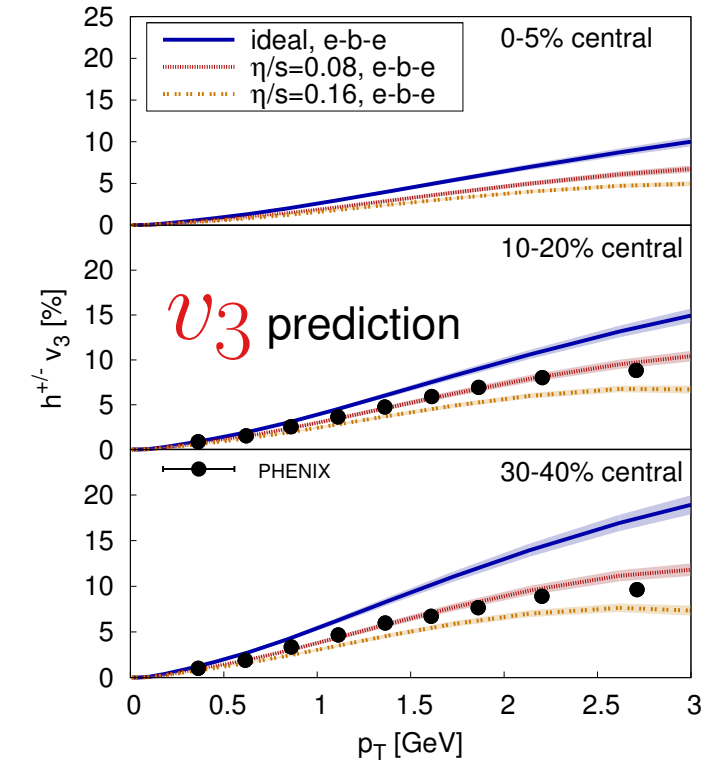
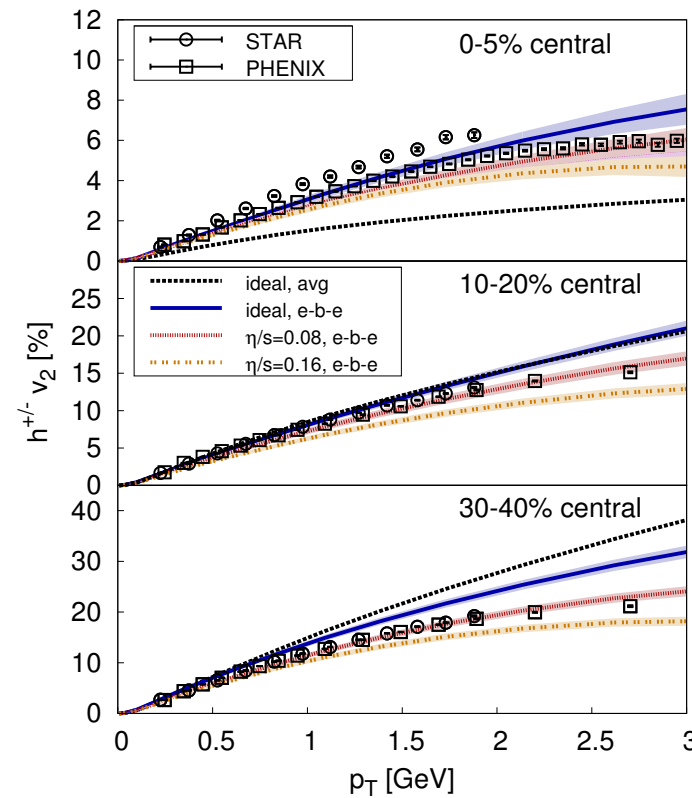
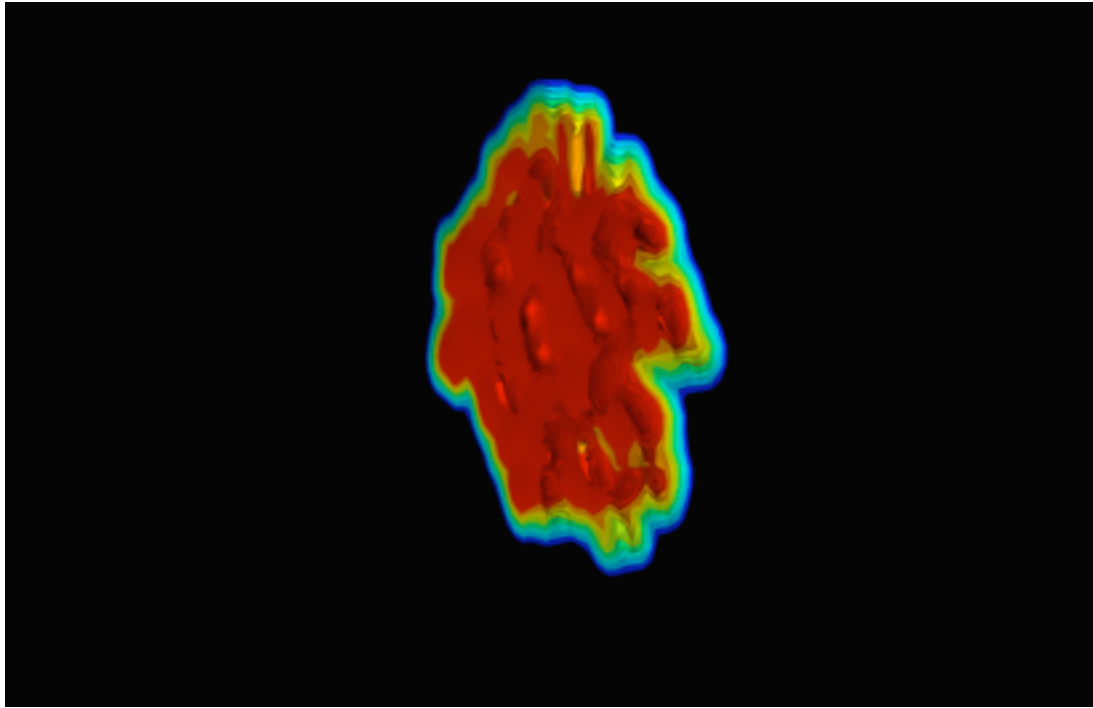


LOOKING UNDER THE HOOD: NON-EQUILIBRIUM EFFECTS



INITIAL STATE FLUCTUATIONS: A PARADIGM SHIFT IN HEAVY ION ANALYSES

Lumpy
MUSIC



Schenke, Jeon, Gale, PRL (2011)

20

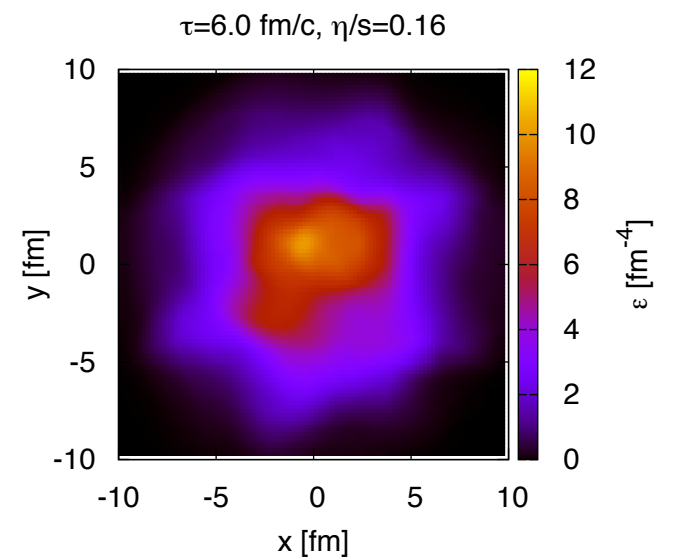
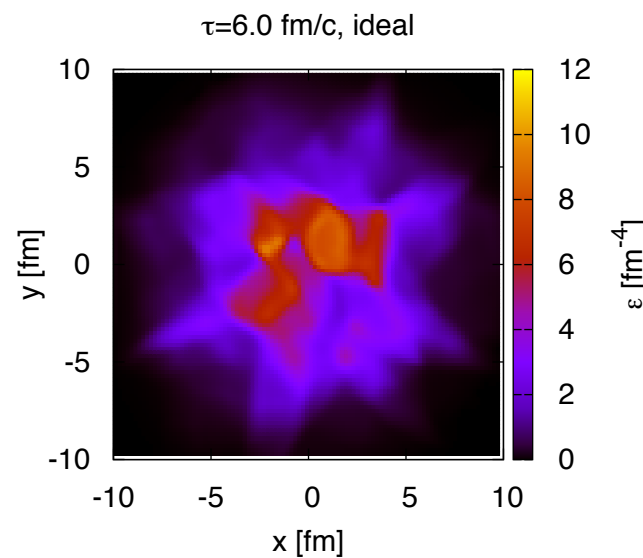
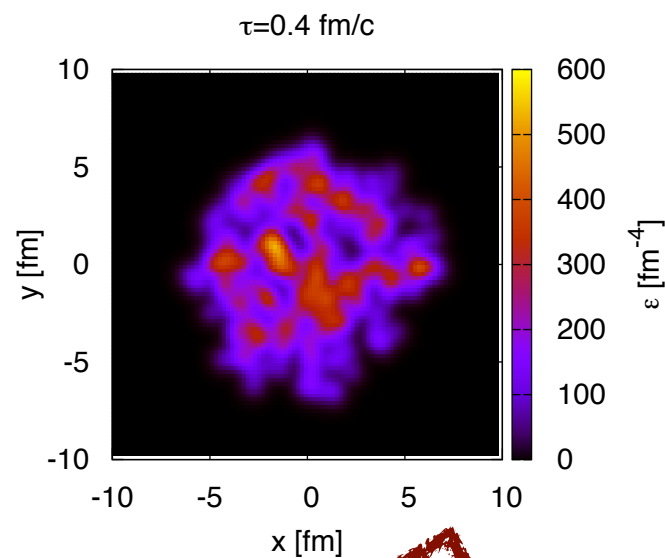
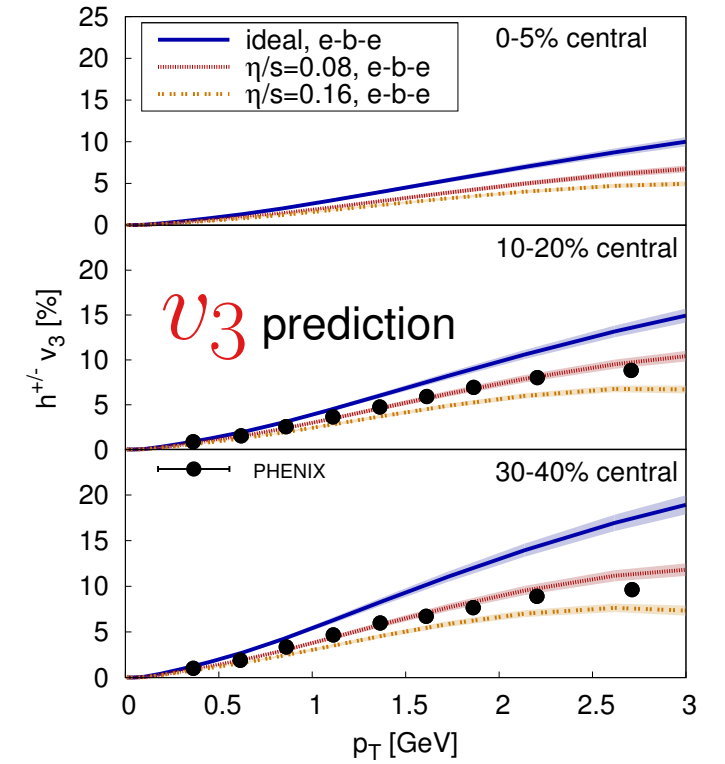
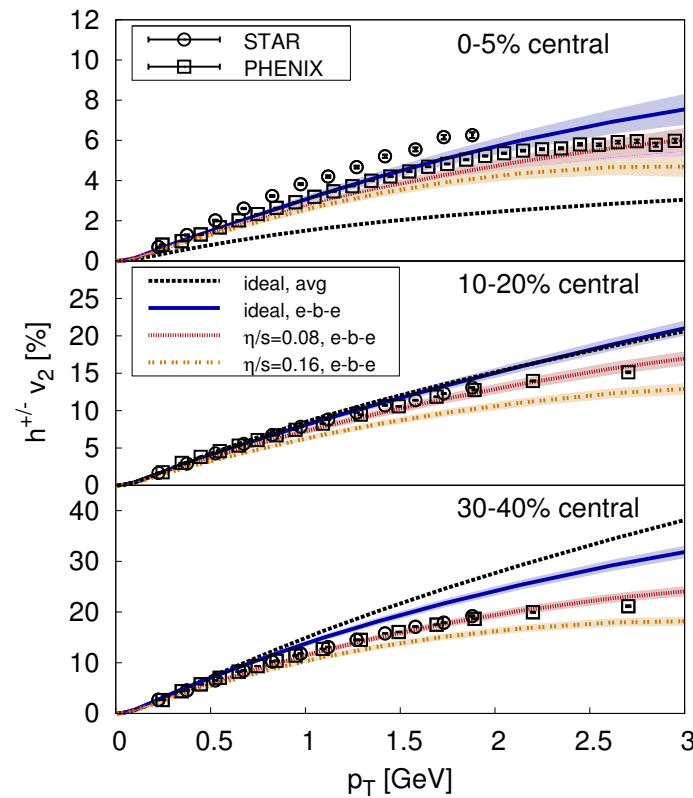
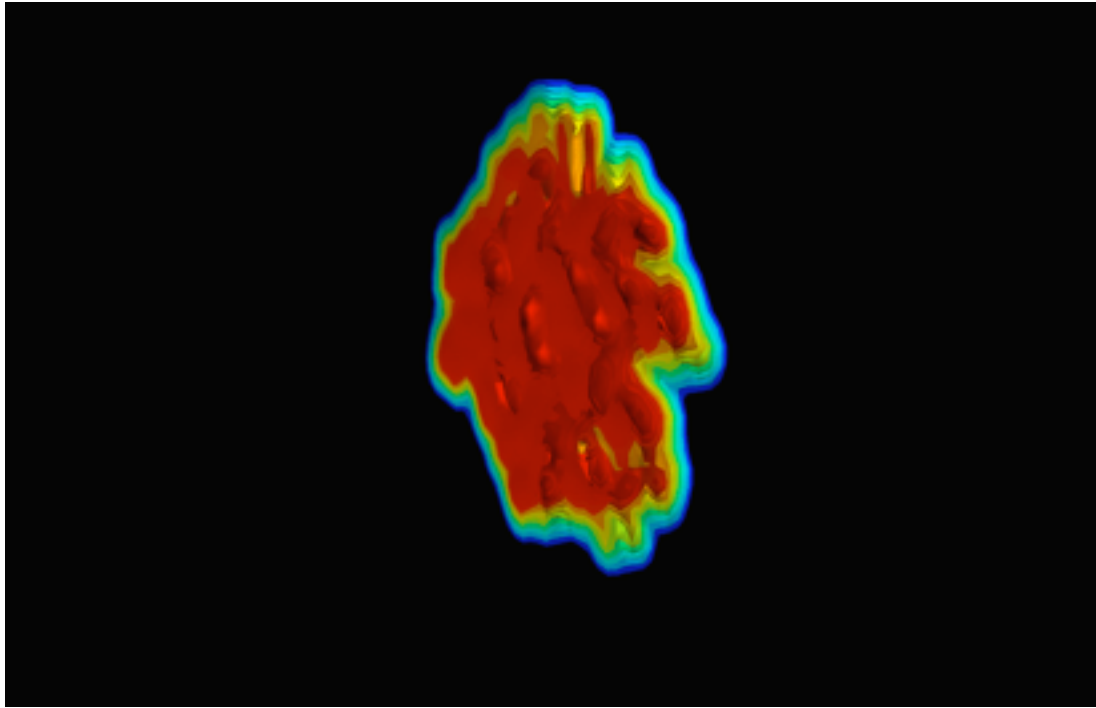


Charles Gale



INITIAL STATE FLUCTUATIONS: A PARADIGM SHIFT IN HEAVY ION ANALYSES

Lumpy
MUSIC

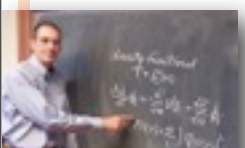


PROGRESS

Schenke, Jeon, Gale, PRL (2011)

20

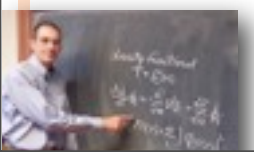
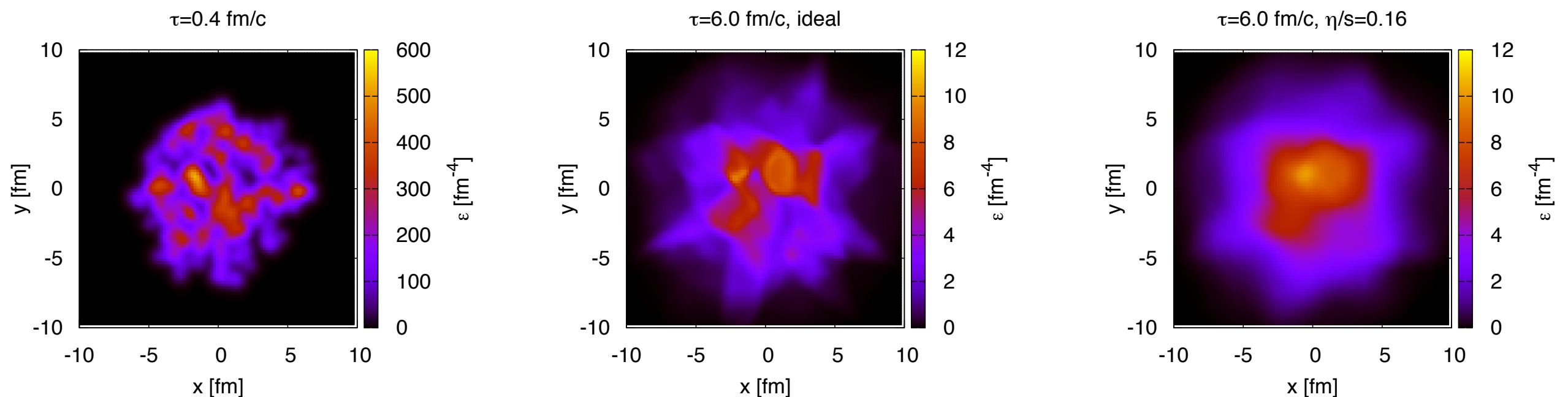
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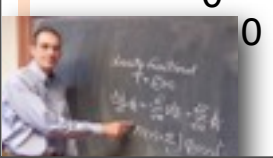
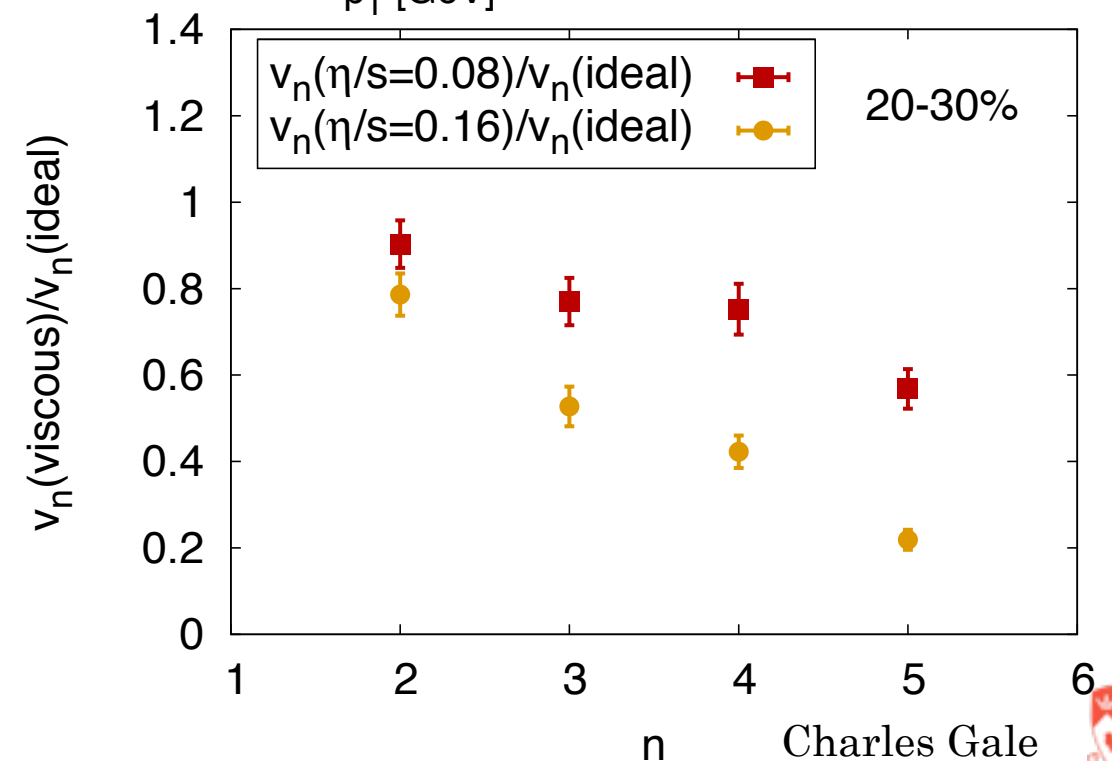
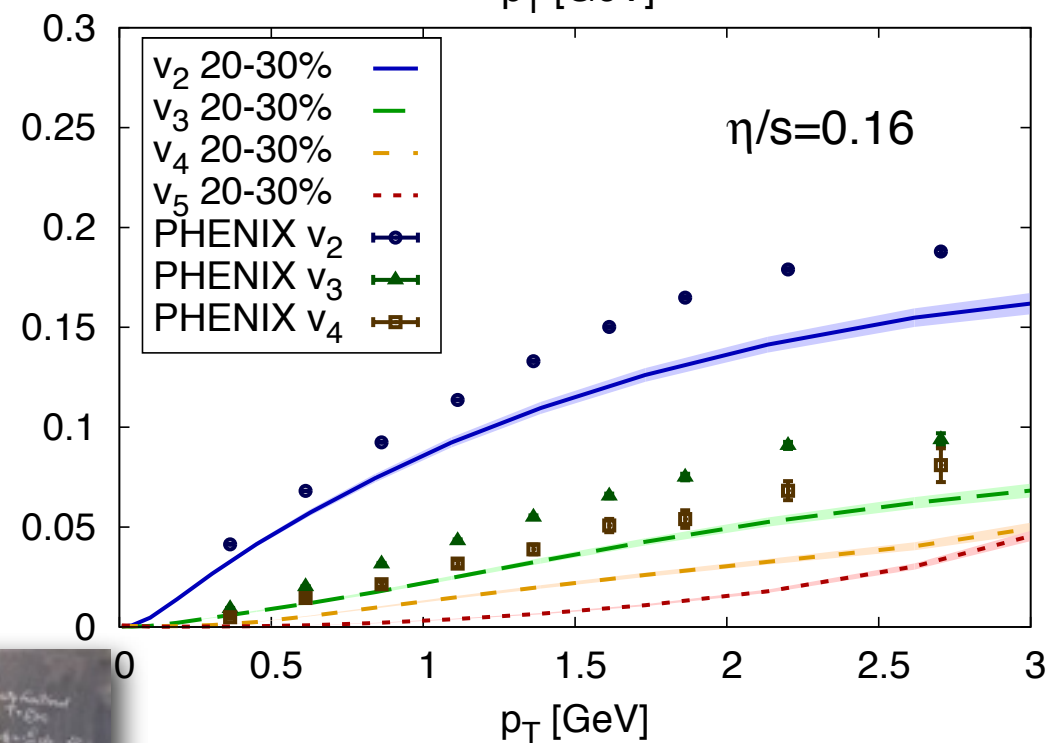
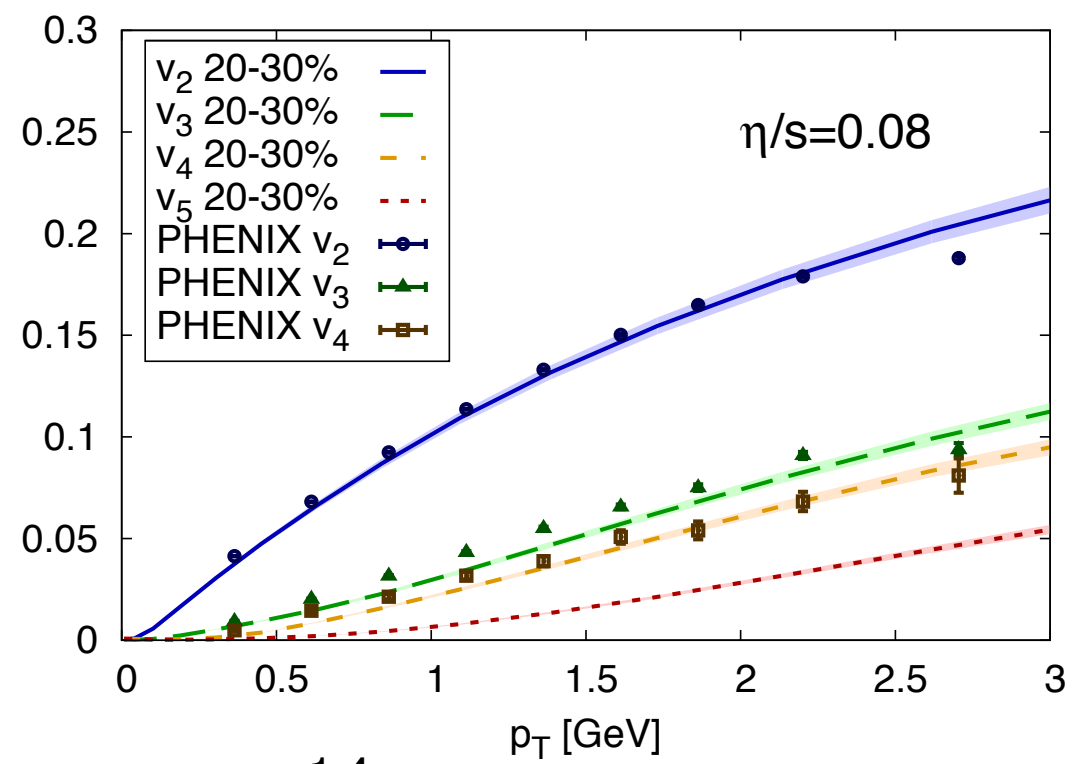
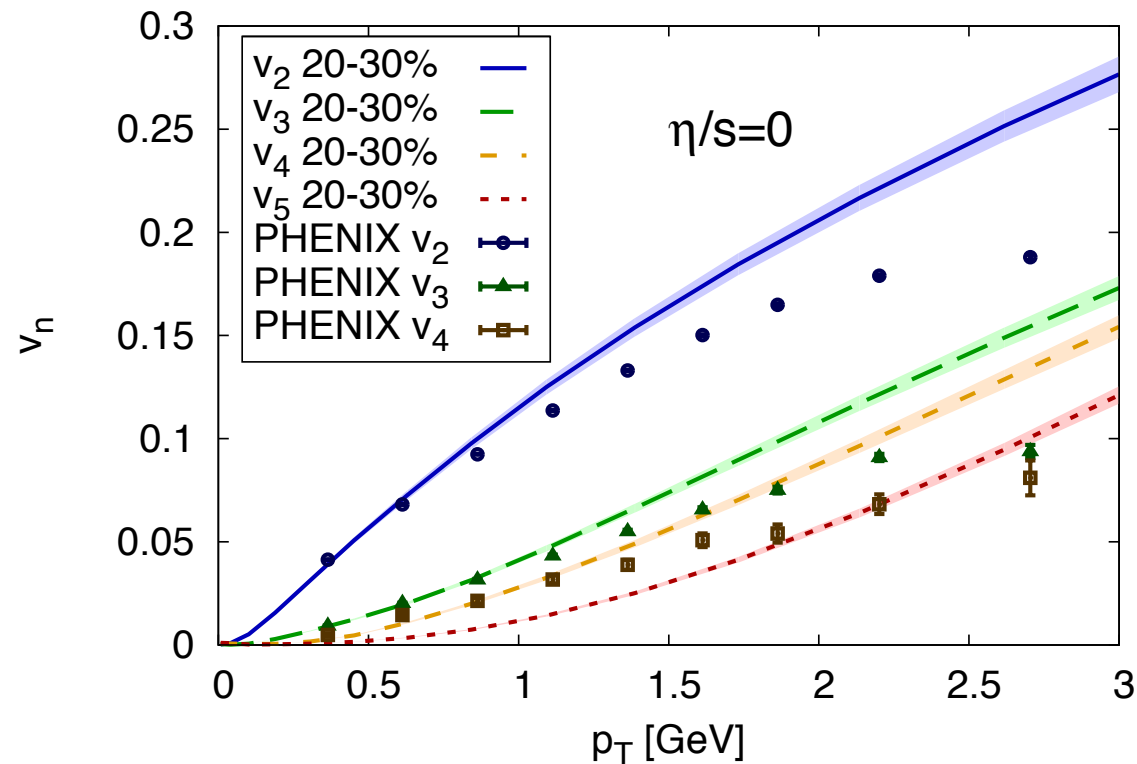
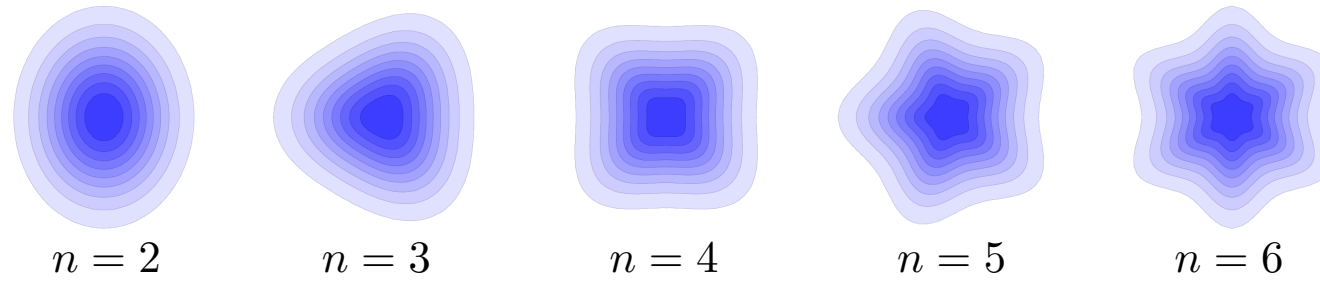
INITIAL STATE FLUCTUATIONS: MC GLAUBER

INITIALIZATION

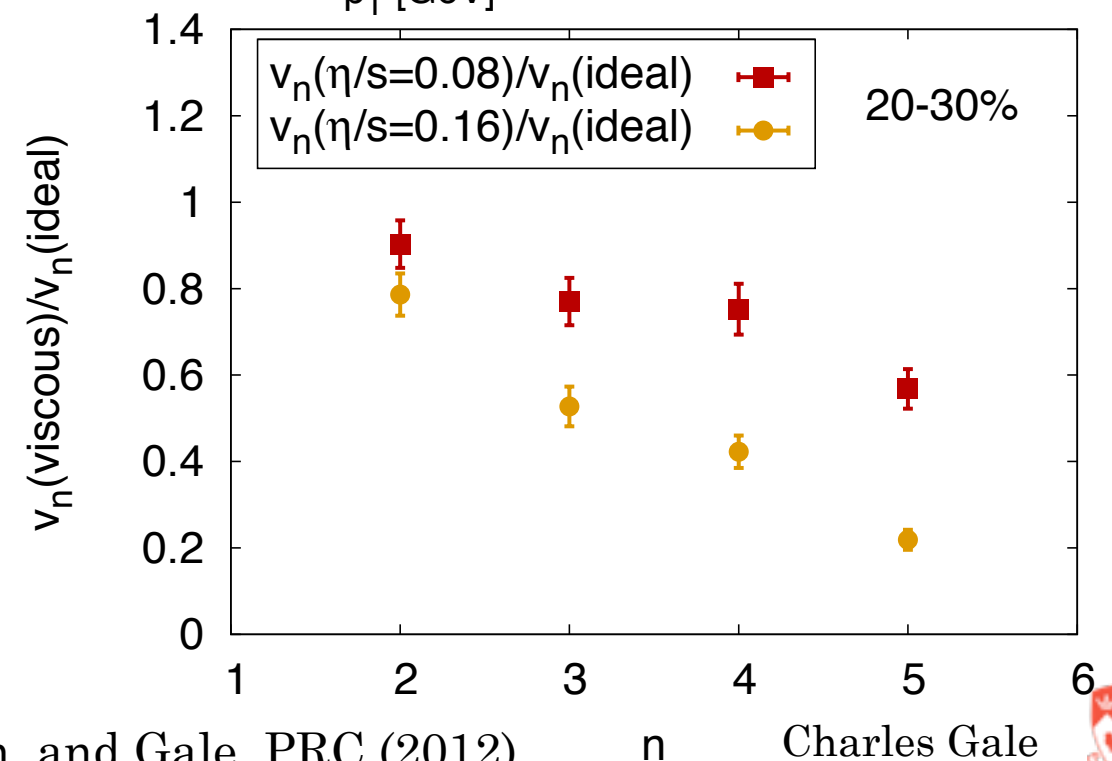
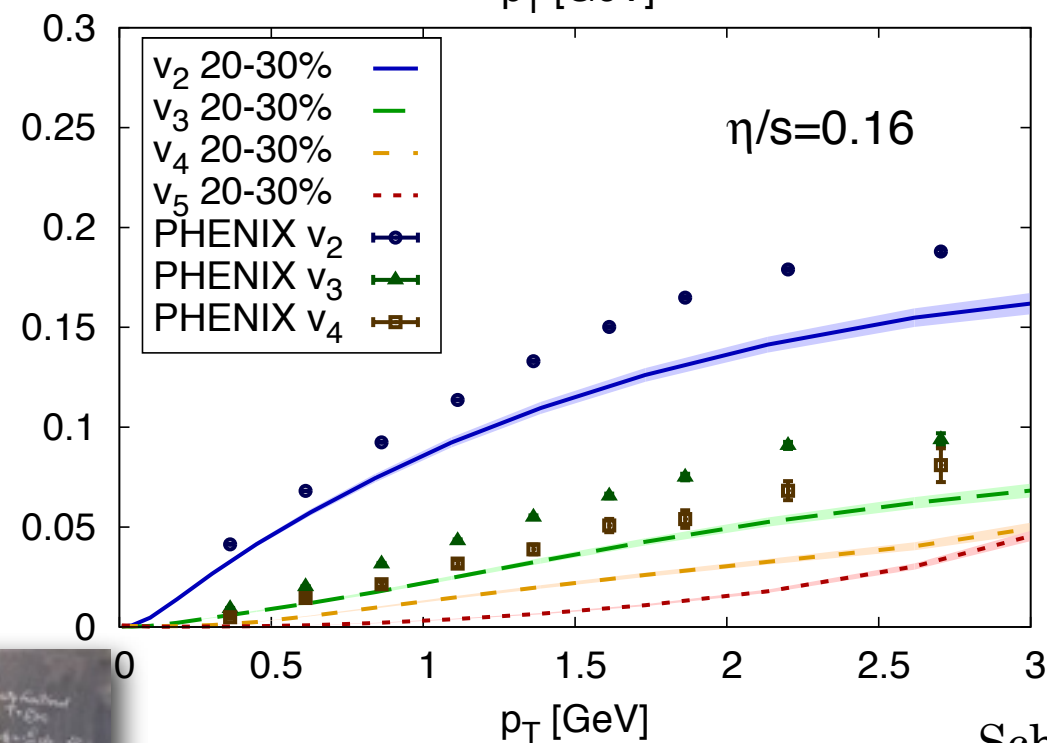
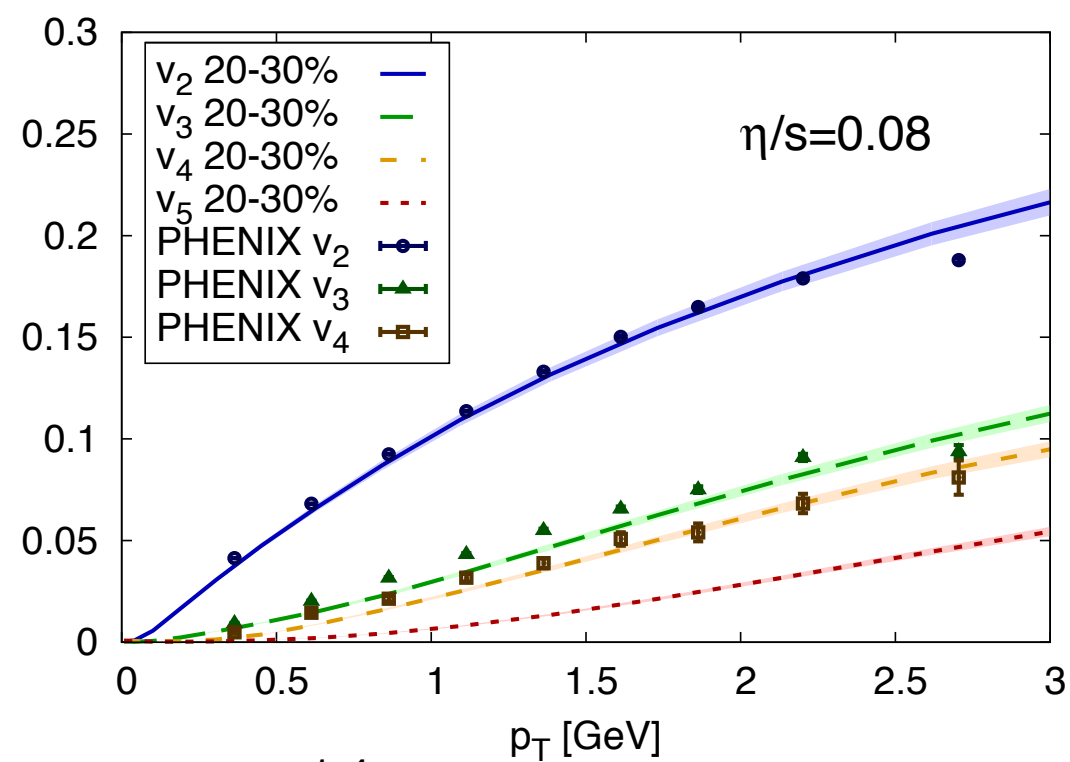
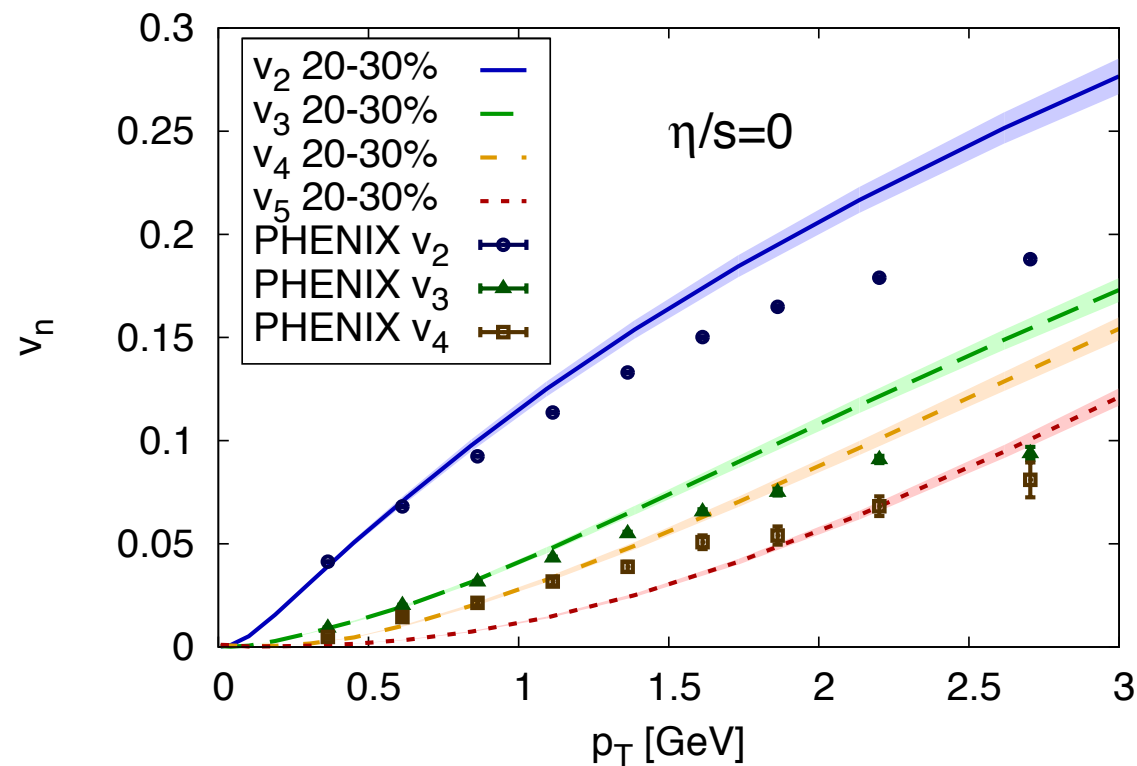
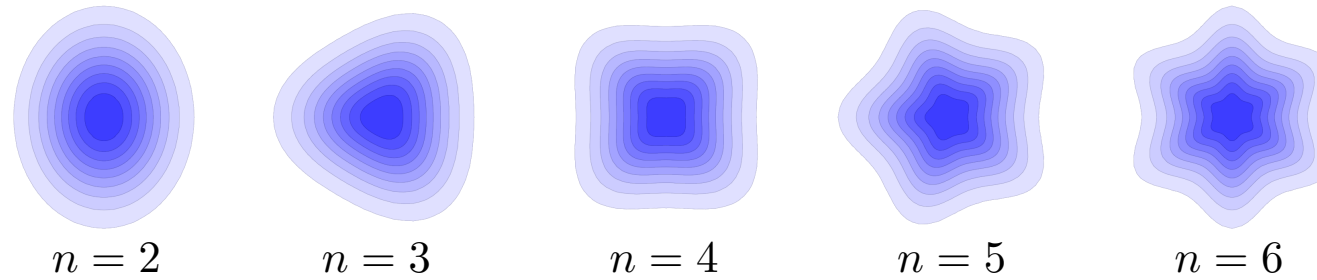
- Sample the nucleon locations from the nuclear density profile (with or without the shell effect deformations)
- Identify the colliding partners ($d \leq \sqrt{\sigma_{NN} / \pi}$)
- Having identified the wounded nucleons, ascribe an energy distribution at each site, with a Gaussian width σ_0 .



MOVING INTO THE "CHARACTERIZATION" PHASE...



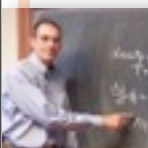
MOVING INTO THE "CHARACTERIZATION" PHASE...



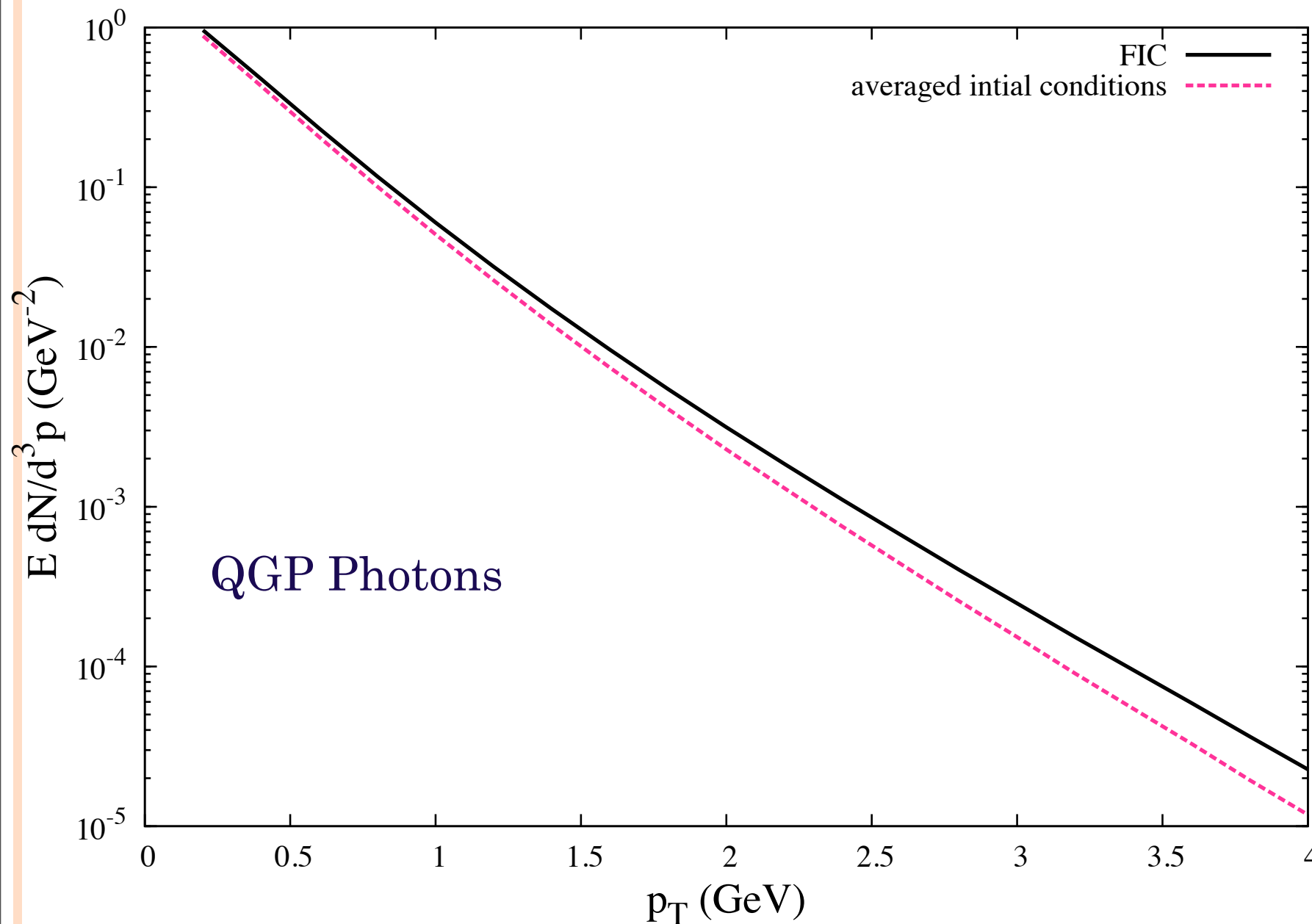
Schenke, Jeon, and Gale, PRC (2012)

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THE EFFECT OF FIC ON THE THERMAL PHOTON SPECTRUM



- FIC produces higher initial T (hot spots), and higher initial gradients
- FIC conditions are demanded by hadronic data (v_{odd})
- These lead to a harder spectrum, *as for hadrons*

Dion et al., PRC (2011)
Chatterjee et al., PRC (2011)

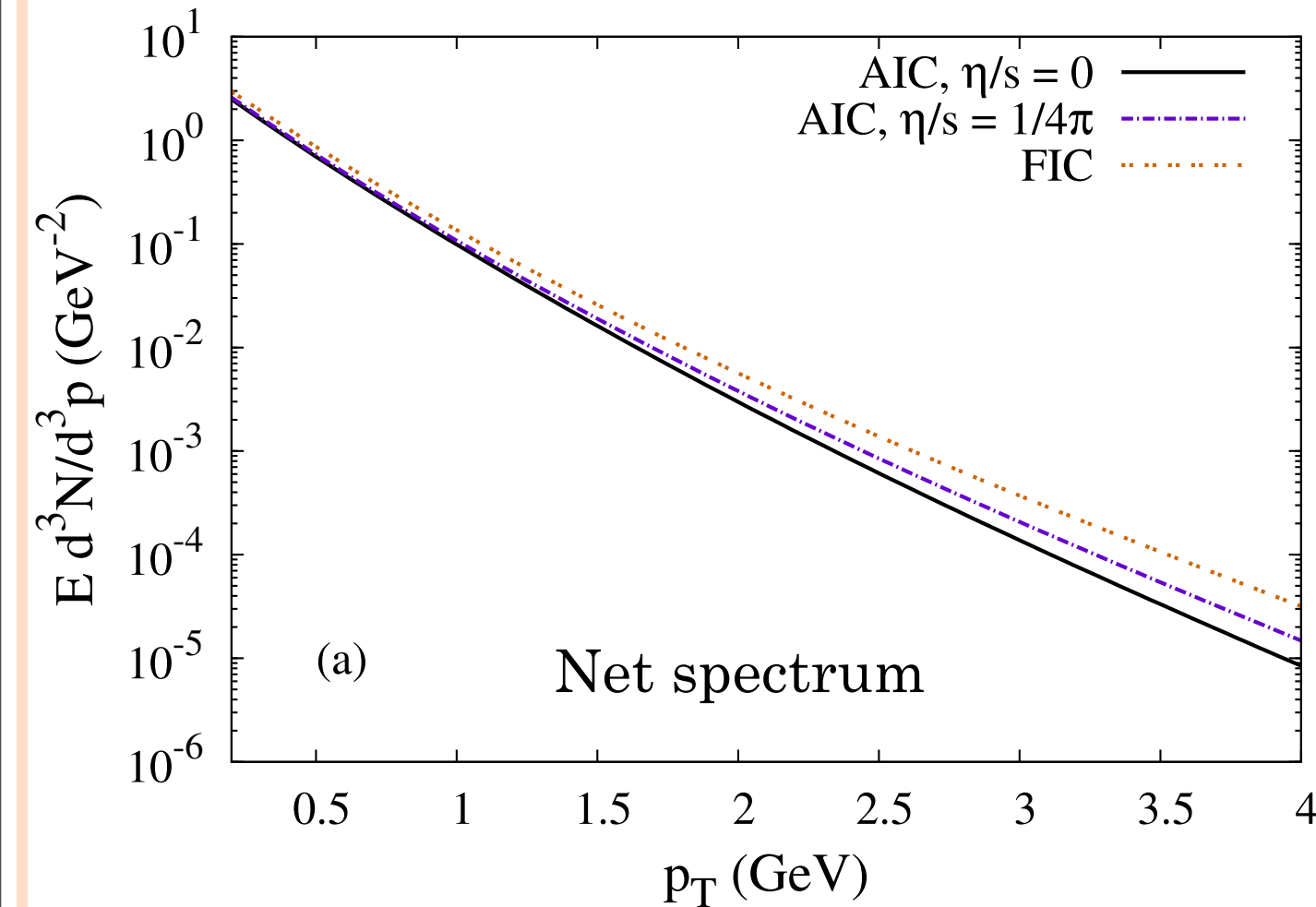


ALL TOGETHER: FIC + VISCOSITY

- Combined with viscous corrections, FIC yield an enhancement by ≈ 5 @ 4 GeV, and ≈ 2 @ 2 GeV
- Temperature estimated by slopes can vary considerably
- A combination of hot spots and blue shift hardens spectra
- Once pQCD photons are included: only modest changes

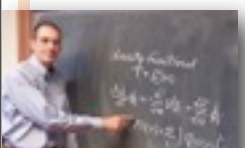


ALL TOGETHER: FIC + VISCOSITY

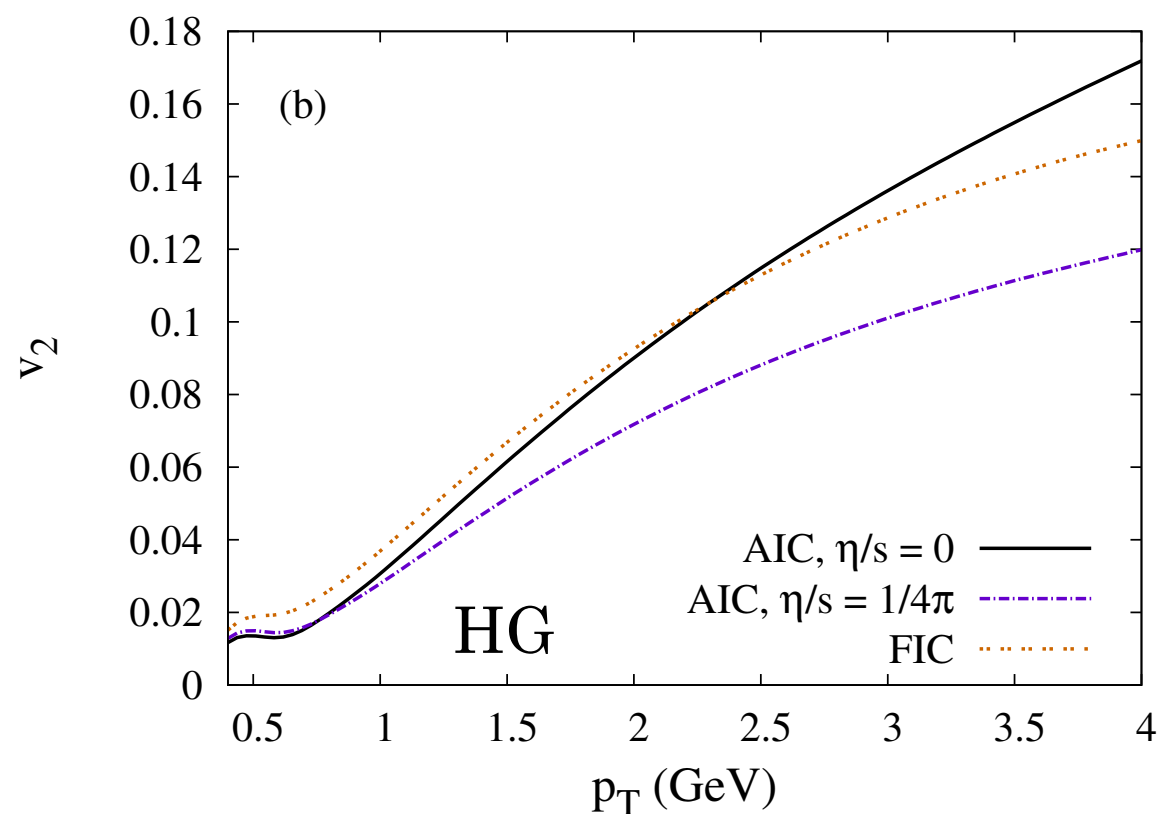
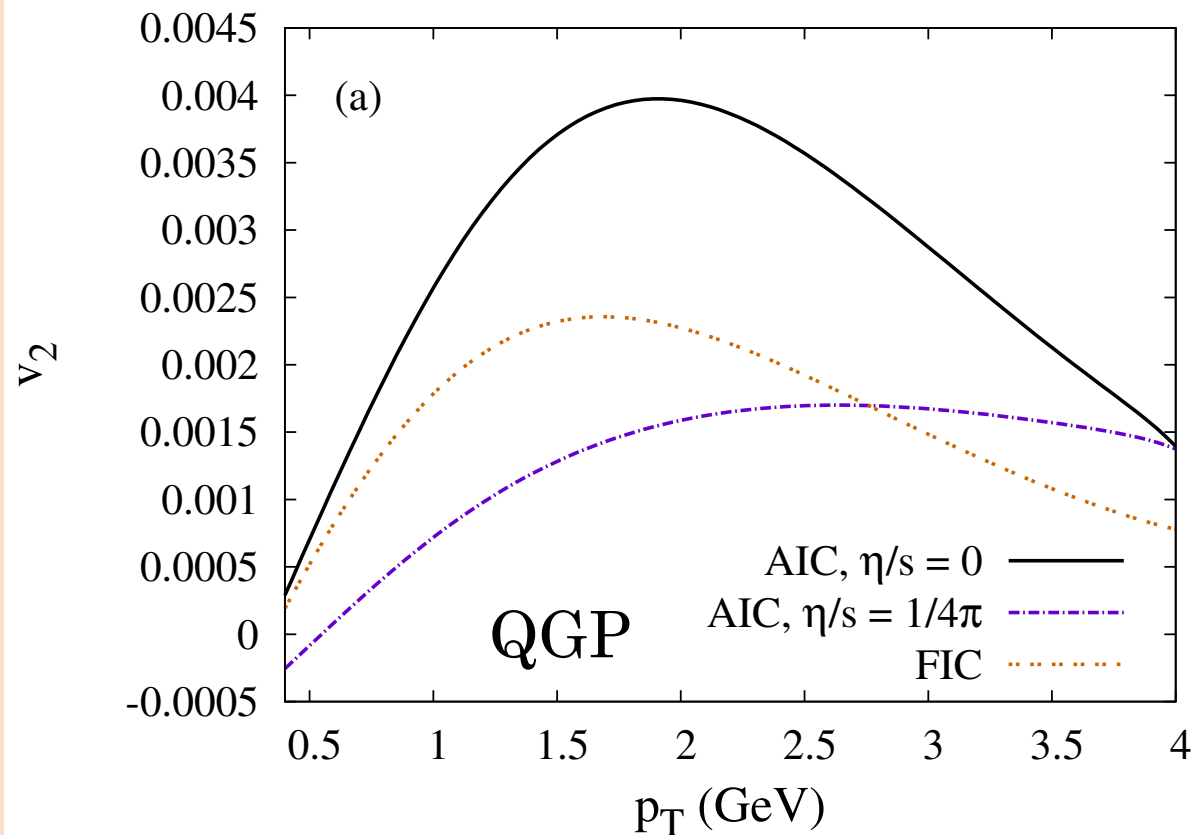


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PROGRESS



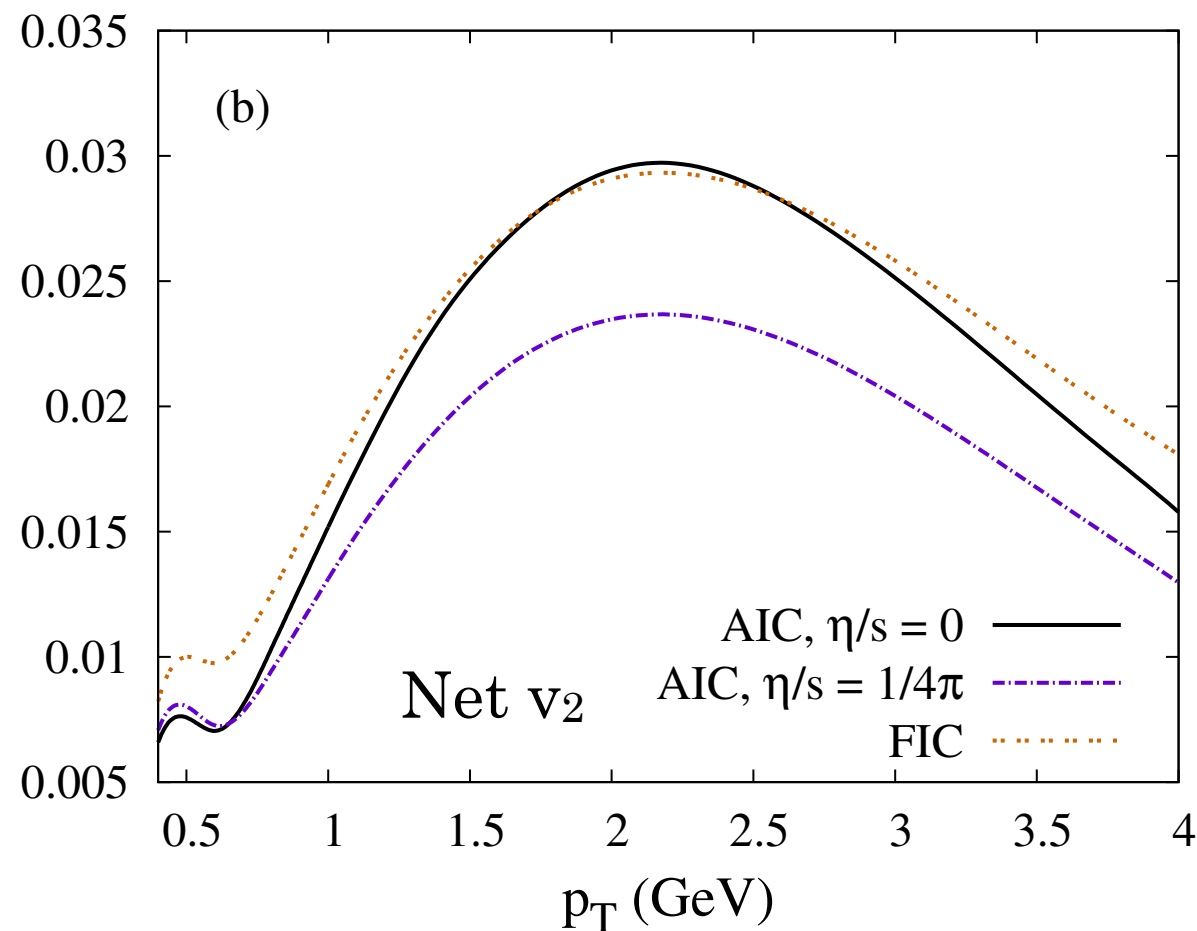
FICs AND THERMAL PHOTON v_2



- FICs enhance v_2 in this centrality class (0-20%), as for hadrons
- For hadrons measured in events belonging to large centrality, FICs will *decrease* v_2
- HG elliptic flow is much larger than QGP elliptic flow, but remember net v_2 is a weighted average

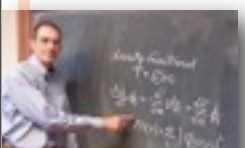


FICs AND THERMAL PHOTON V_2

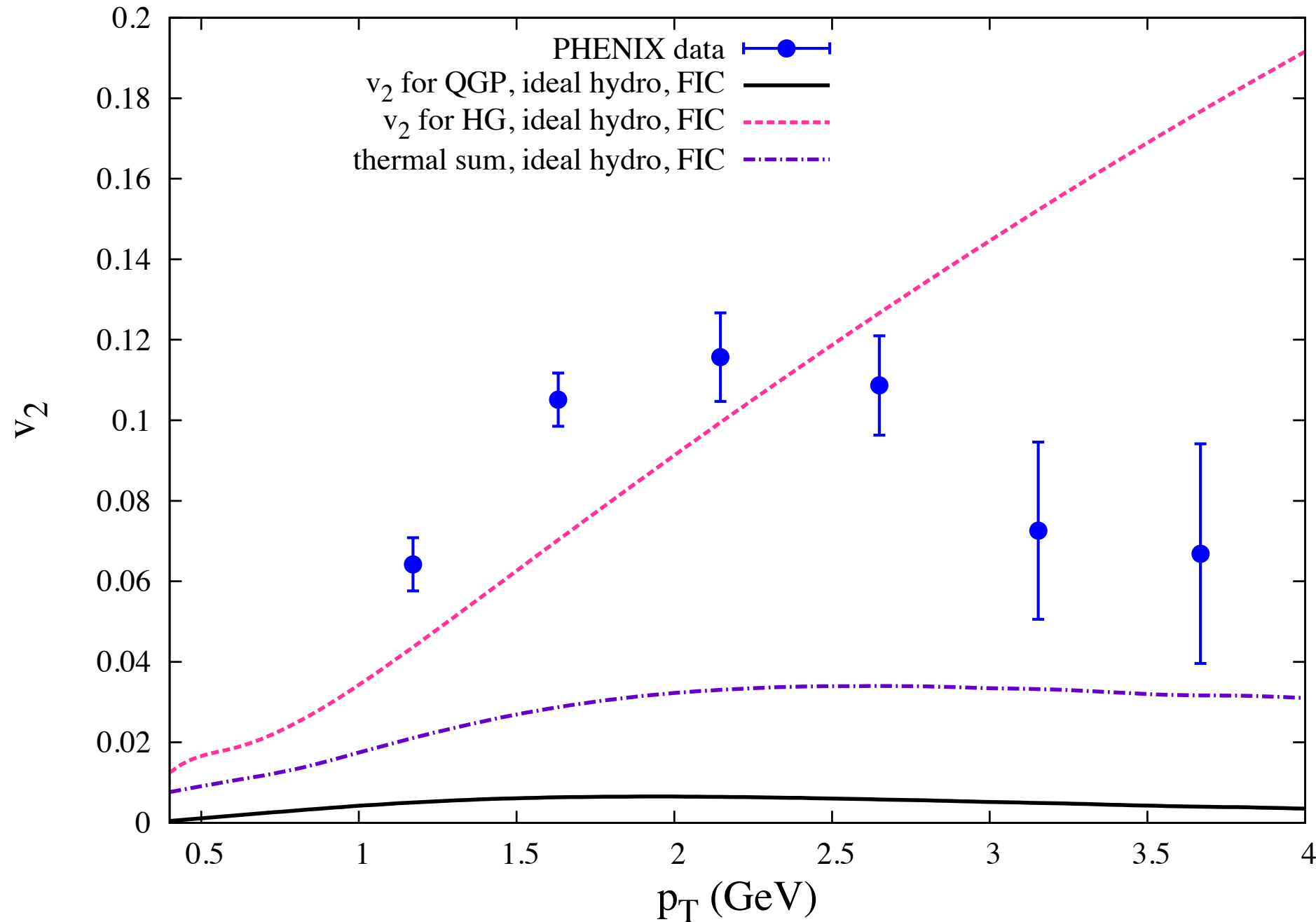


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- HG elliptic flow is much larger than QGP elliptic flow, but remember net v_2 is a weighted average
- Net v_2 is comparable in size to that with ideal medium.

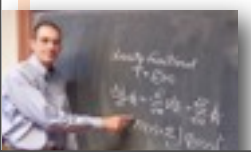
PROGRESS



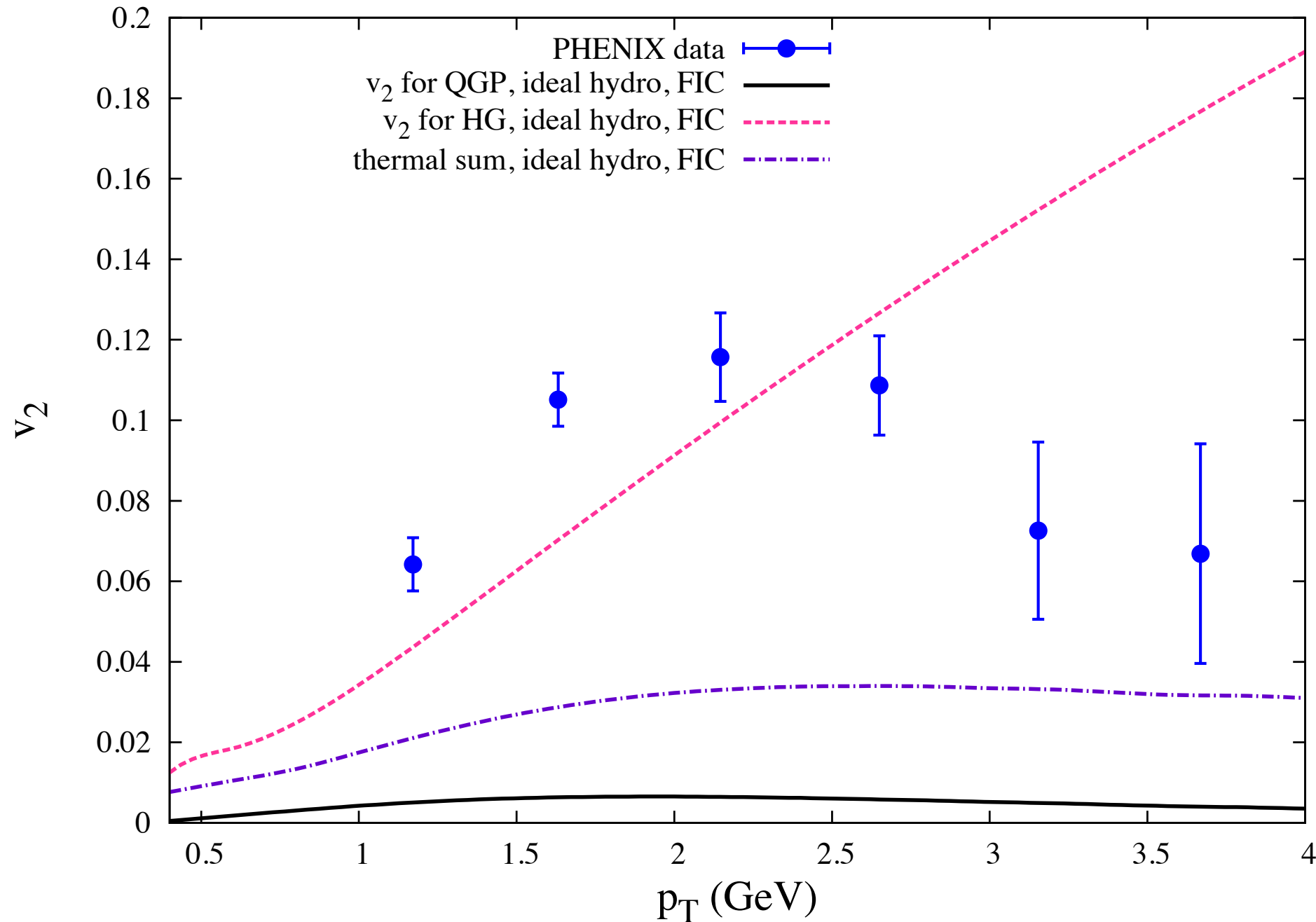
PHOTON v_2 DATA?



- New data is higher than calculation, even with e-b-e initial state fluctuations, and ideal hydro
- Size comparable with HG v_2



PHOTON v_2 DATA?

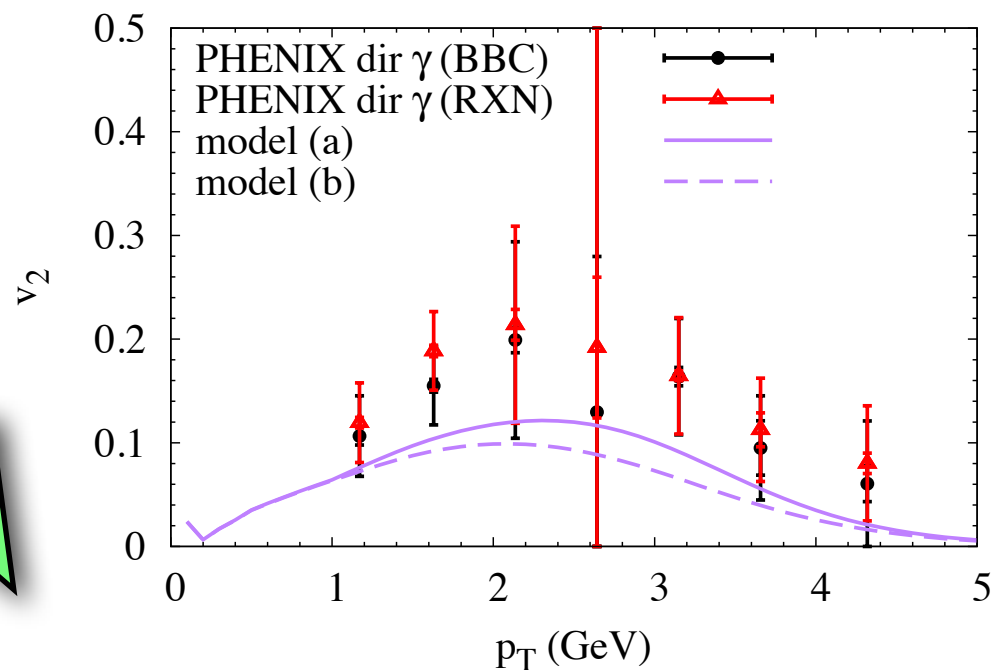
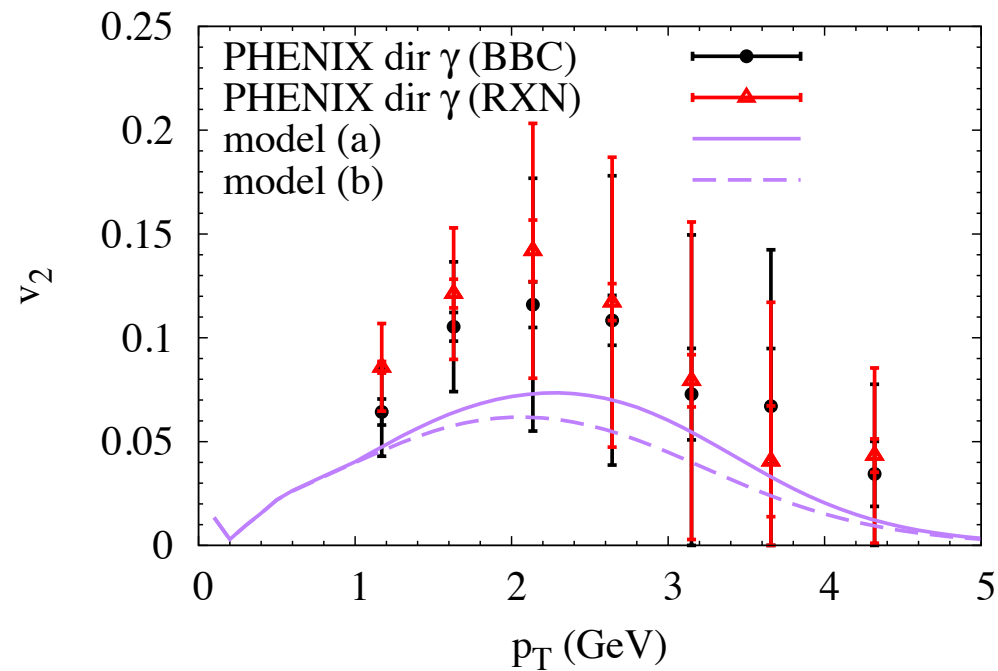
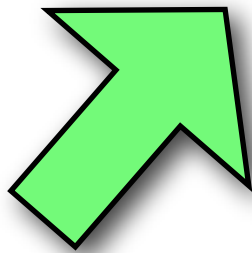
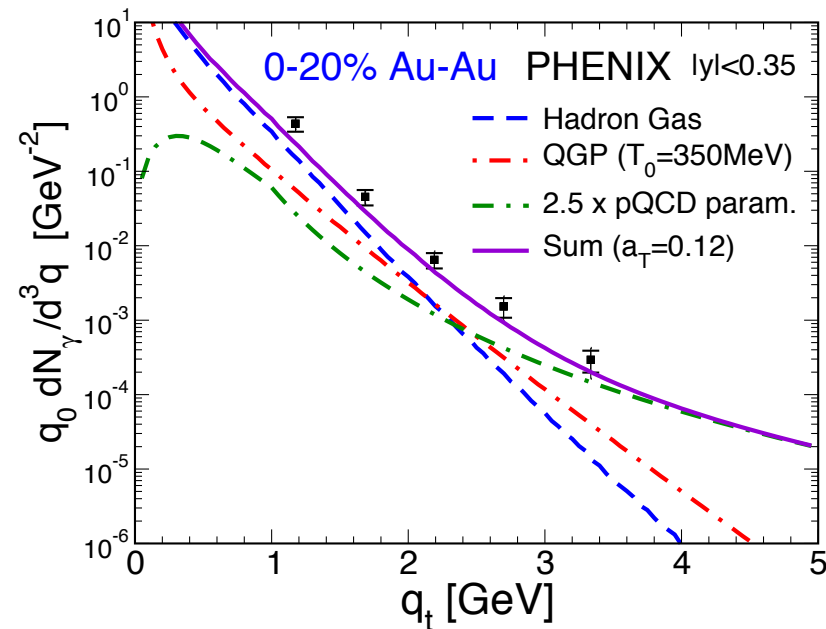
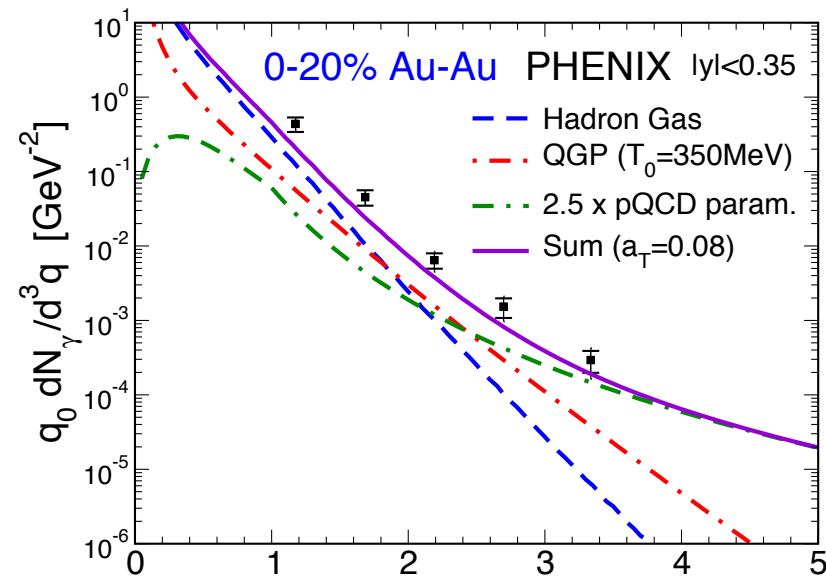
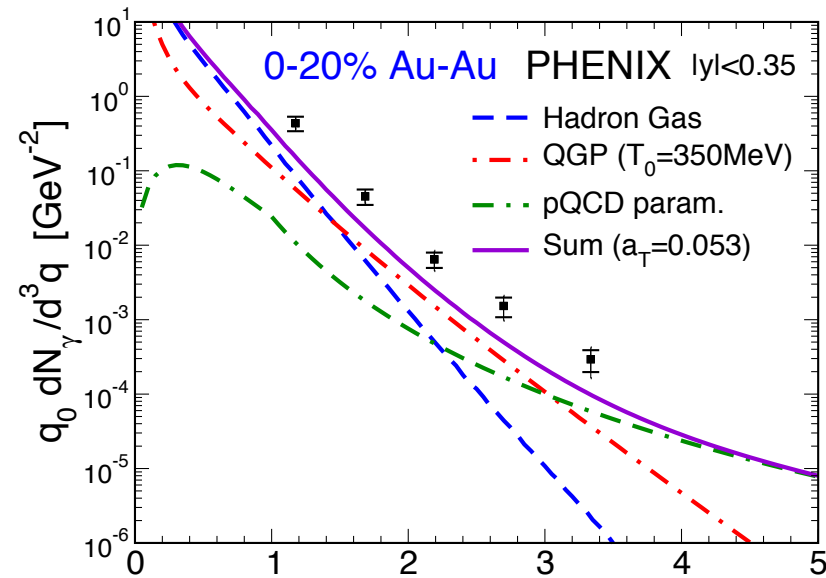


- New data is higher than calculation, even with e-b-e initial state fluctuations, and ideal hydro
- Size comparable with HG v_2

PUZZLE!



CAN THE DATA TEACH US ABOUT DYNAMICS?



van Hees, Rapp, Gale PRC (2011)



SOME FACTS AND SOME LEADS

- FICs are here to stay. The meaning of “initial temperature” is altered.
- Need to explore hydro initialization and parameters. This requires consistency with the hadronic data.
- Making the QGP signal larger will *decrease* the v_2 . Including the $T=0$ photons, will *decrease* v_2 .
- Non-zero initial shear tensor? Primordial flow?

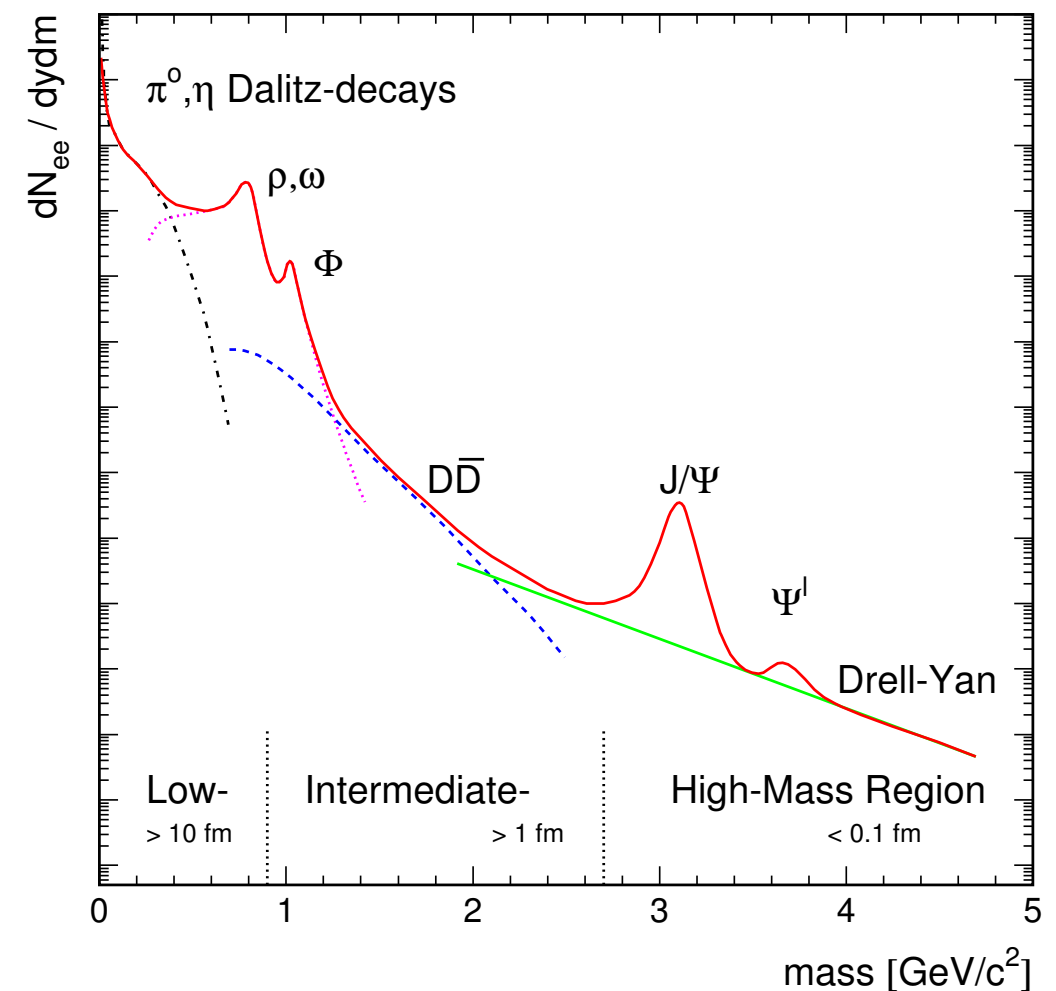


WHAT ABOUT DILEPTONS?

THERMAL DILEPTON SPECTRUM, AND ELLIPTIC FLOW

$$v_2(M, p_T, b) = \frac{\int d\phi \cos(2\phi) \frac{d^4 N}{dM^2 dy p_T dp_T d\phi}}{\int d\phi \frac{d^4 N}{dM^2 dy p_T dp_T d\phi}}$$

Chatterjee, Srivastava, Heinz, Gale, PRC (2007)

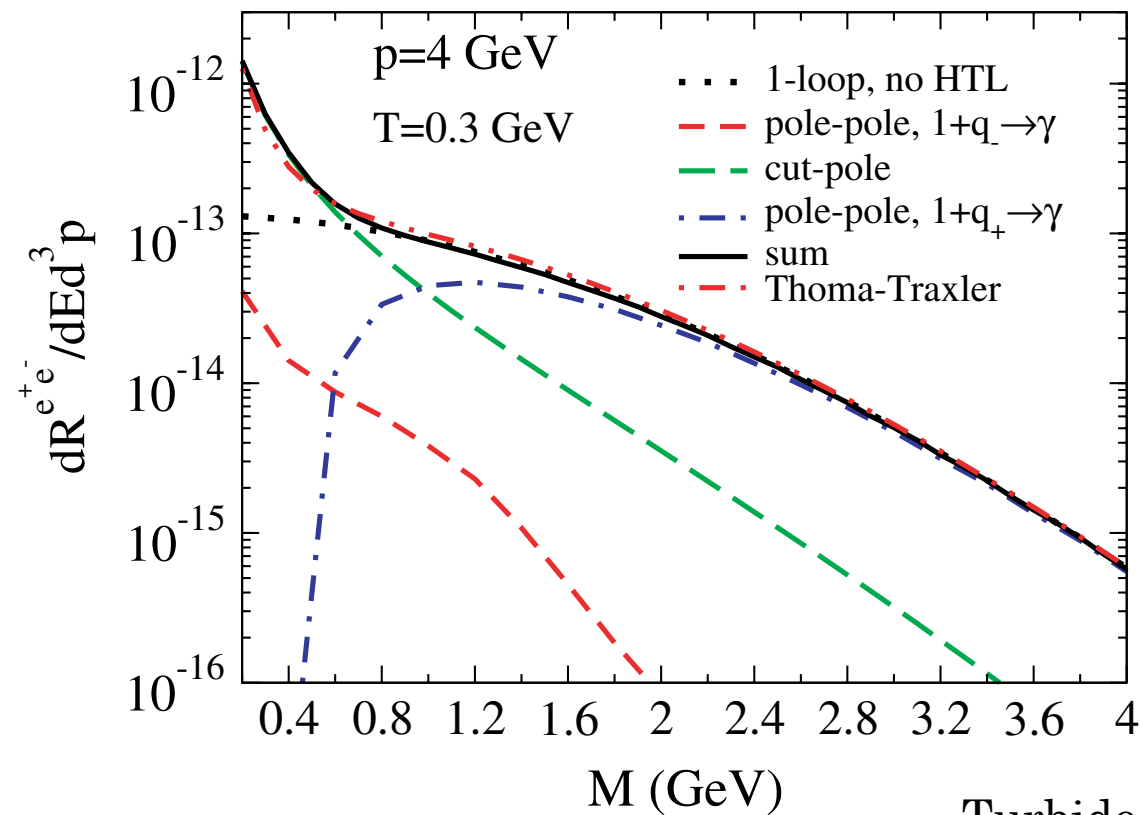


- Additional degree of freedom: M and p_T may be varied independently



THERMAL DILEPTON SOURCES, QGP

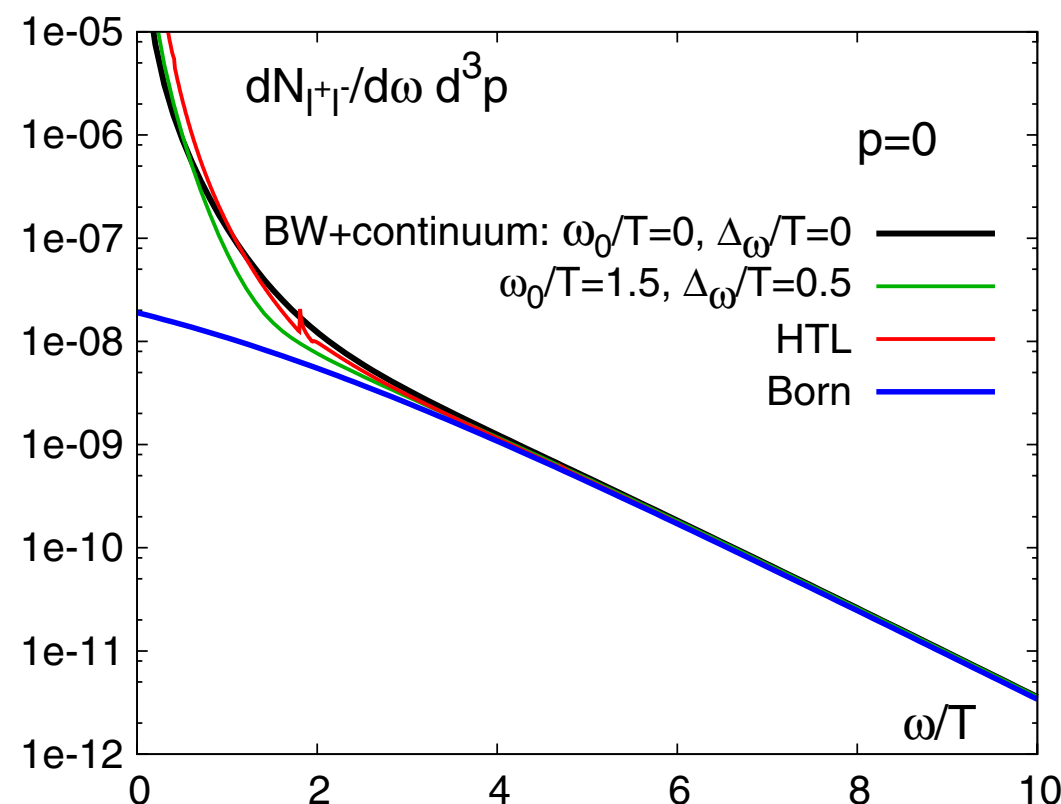
HTL at finite momentum:



Turbide, Gale, Srivastava, Fries PRC (2006)

Non-perturbative estimate:

∴ No single calculation covers the entire dilepton kinematical phase space

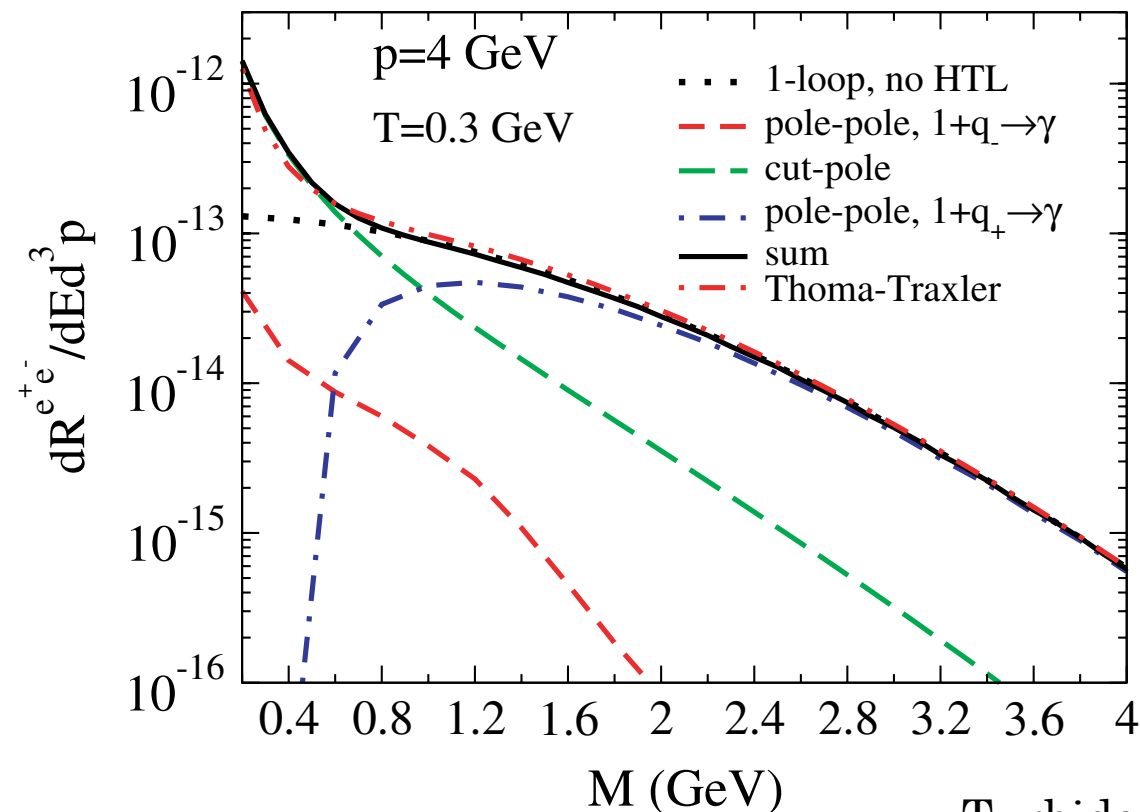


Ding et al., PRD (2011)



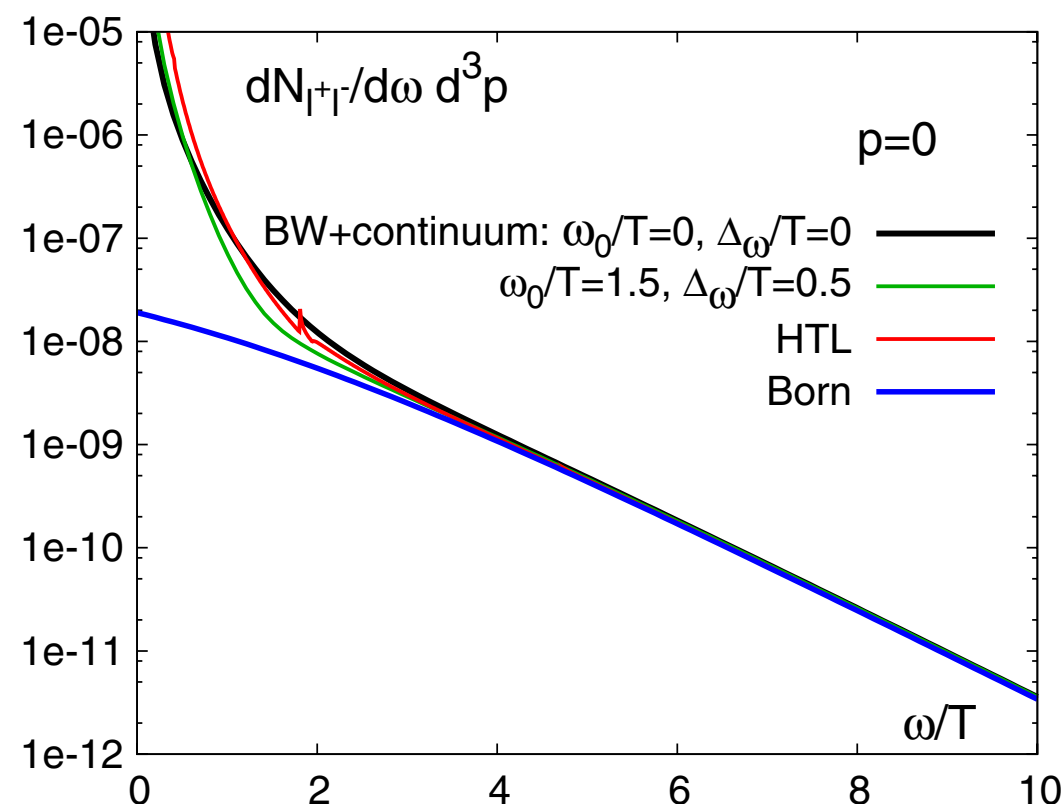
THERMAL DILEPTON SOURCES, QGP

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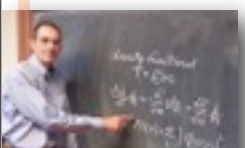
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PROGRESS

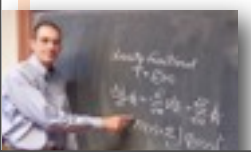


THERMAL DILEPTON SOURCES, HG

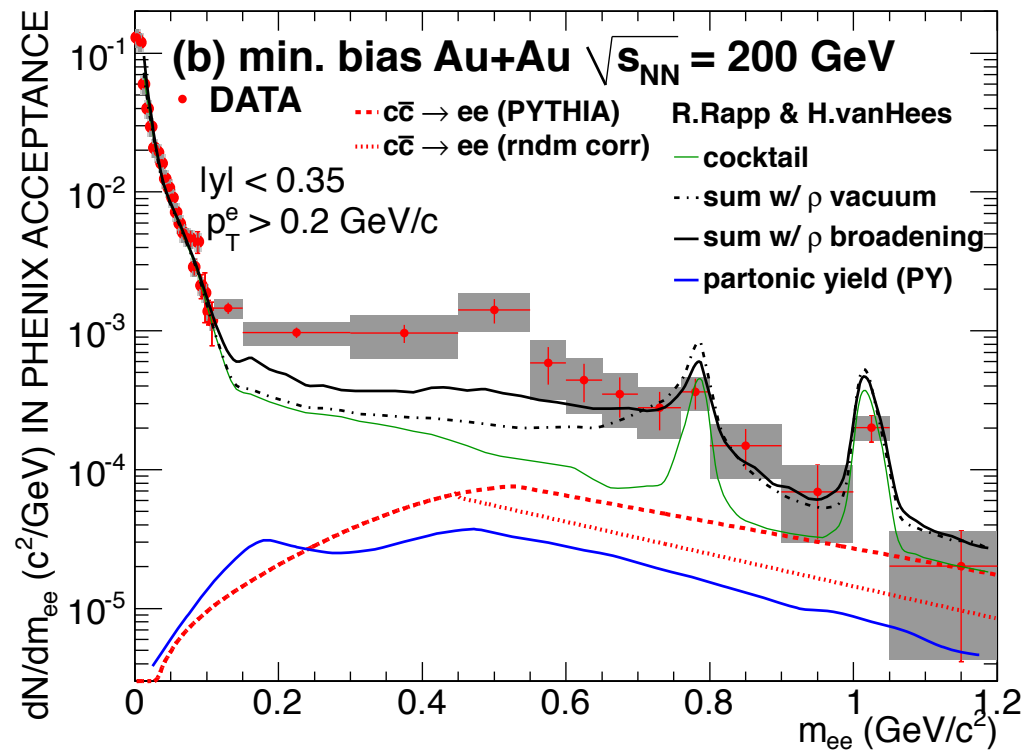
- HG contribution: calculate the in-medium vector spectral density
- Many-Body approach with hadronic effective Lagrangians
 - Rapp and Wambach, ANP (2000)
- Empirical evaluation of the vector mesons forward-scattering amplitudes

$$\Pi_{ab}(E, p) = -4\pi \int \frac{d^3k}{(2\pi)^3} n_b(\omega) \frac{\sqrt{s}}{\omega} f_{ab}^{\text{c.m.}}(s)$$

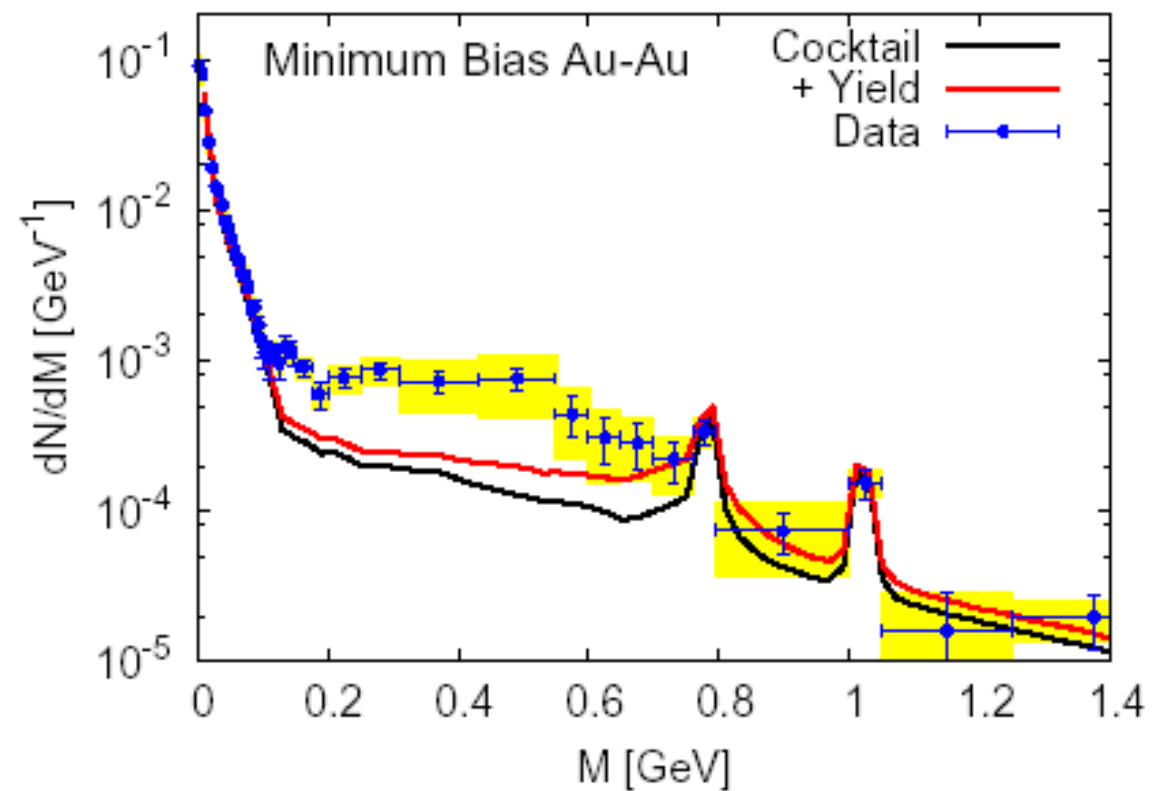
- E. Shuryak, NPA (1991)
- Eletsky, Ioffe, Kapusta (1999)
- Vujanovic, Gale (2009)
- Chiral Reduction formulae
 - Yamagishi, Zahed (1996)



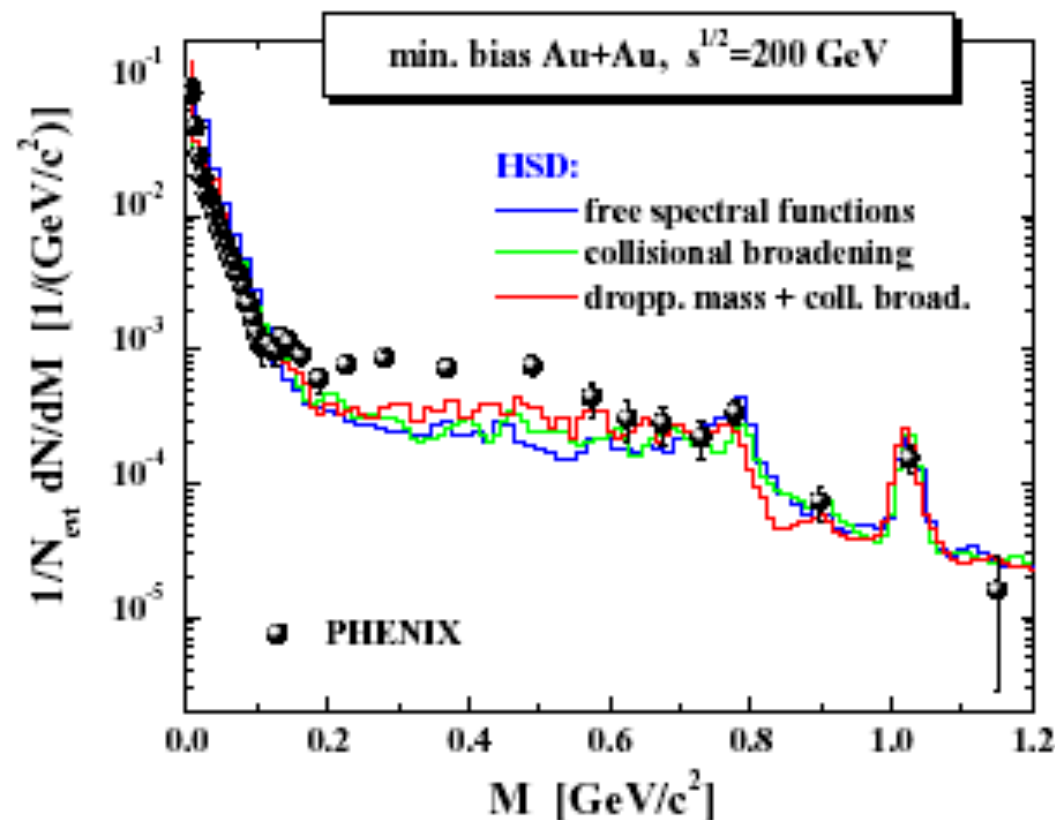
DILEPTONS, THE STORY AS OF A FEW MONTHS AGO



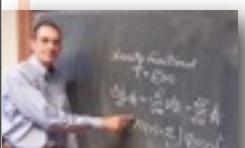
van Hees, Rapp (2010)



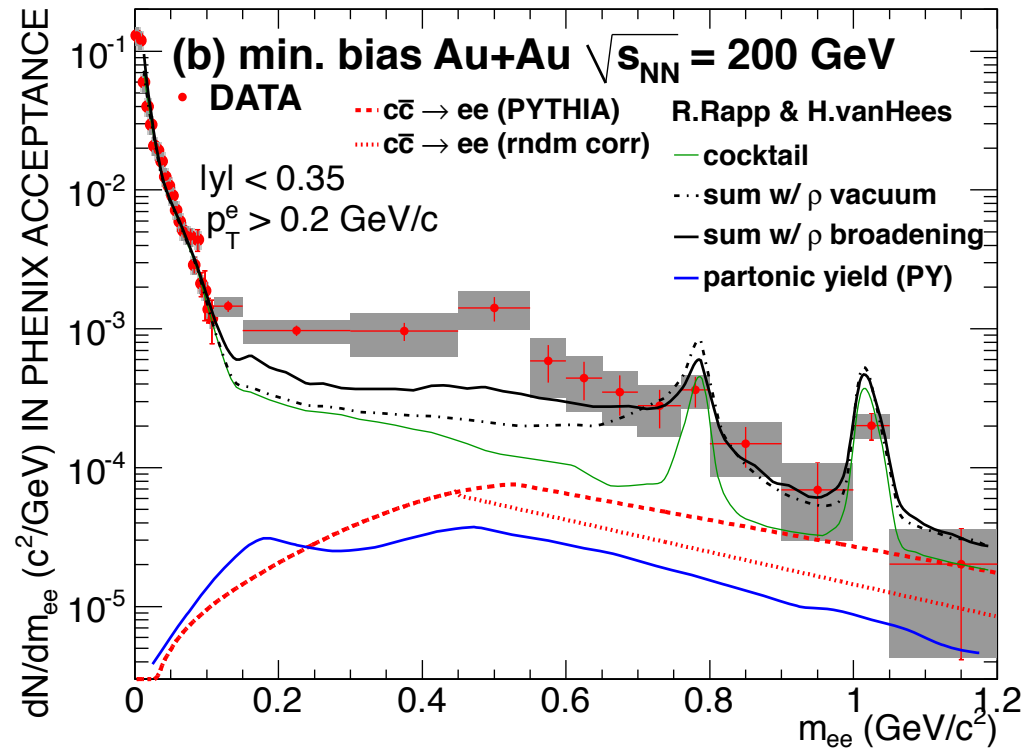
Dusling, Zahed (2009)



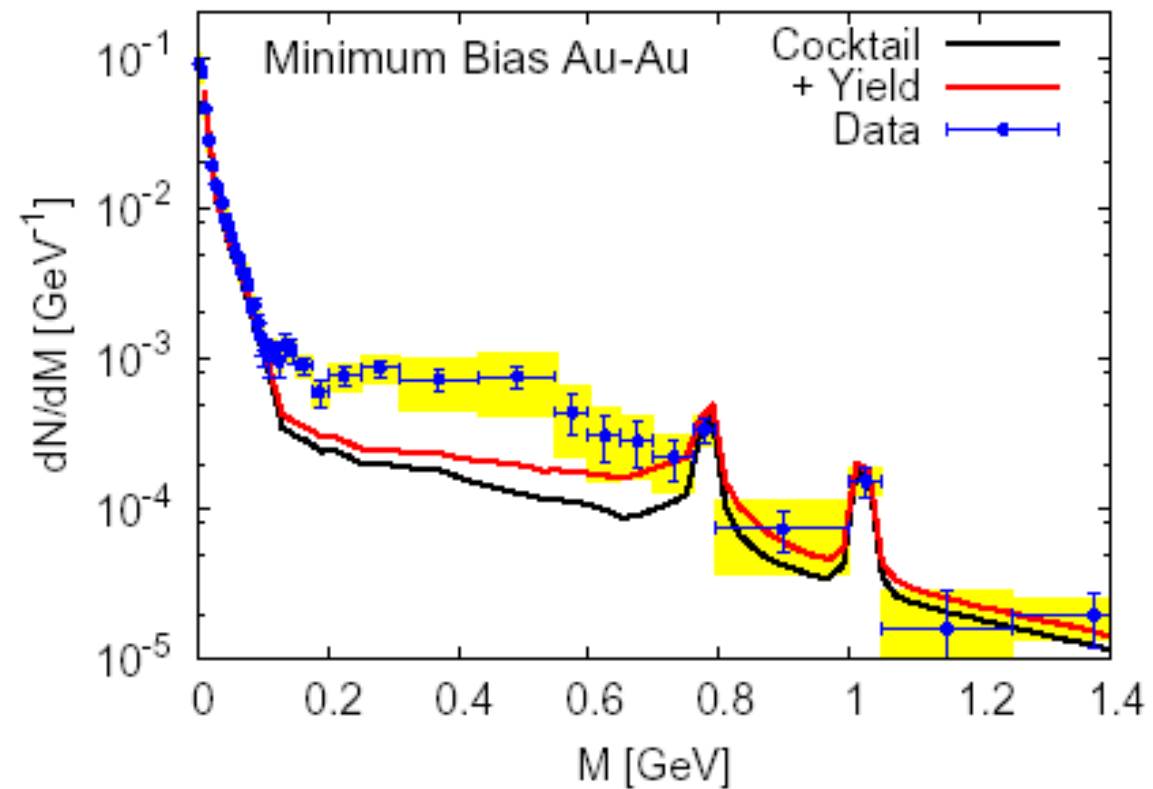
Bratkovskaya, Cassing, Linnyk (2012)



DILEPTONS, THE STORY AS OF A FEW MONTHS AGO

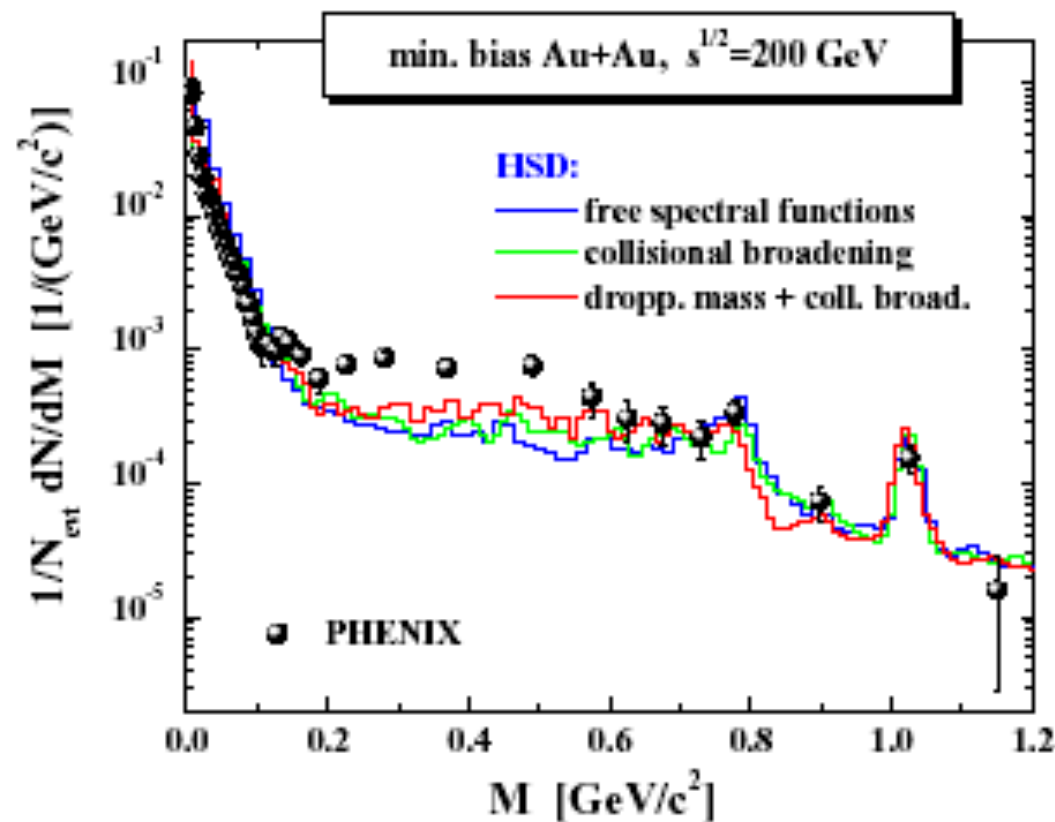


van Hees, Rapp (2010)

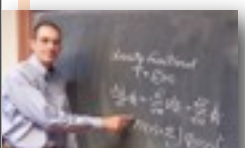


Dusling, Zahed (2009)

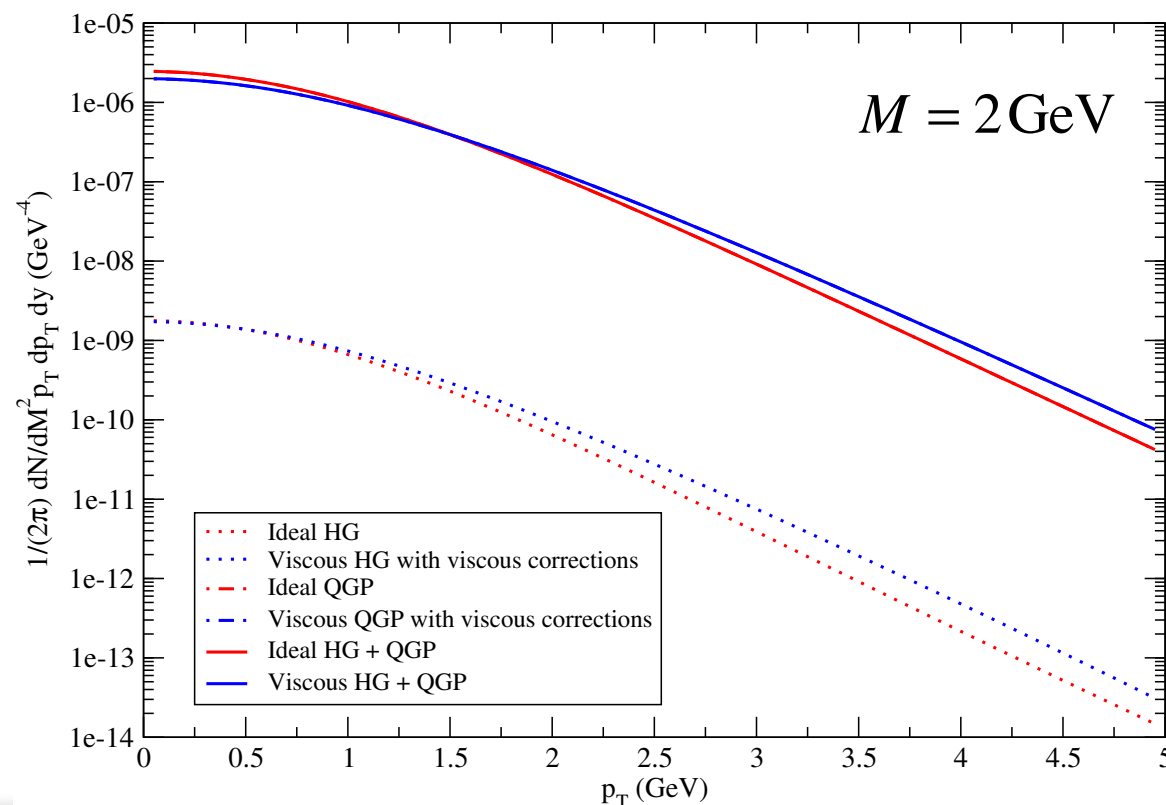
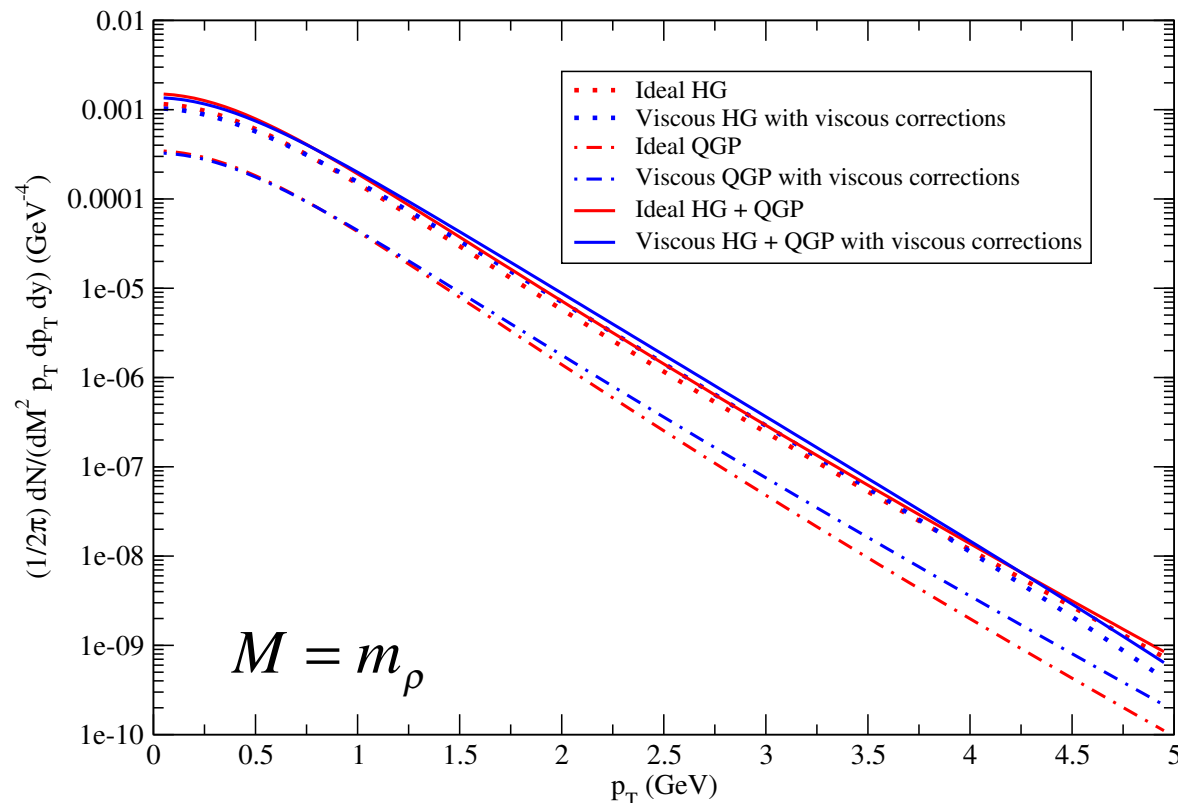
PUZZLE!



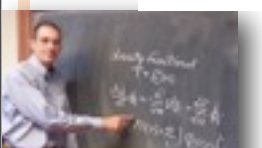
Bratkovskaya, Cassing,
 Linnyk (2012)



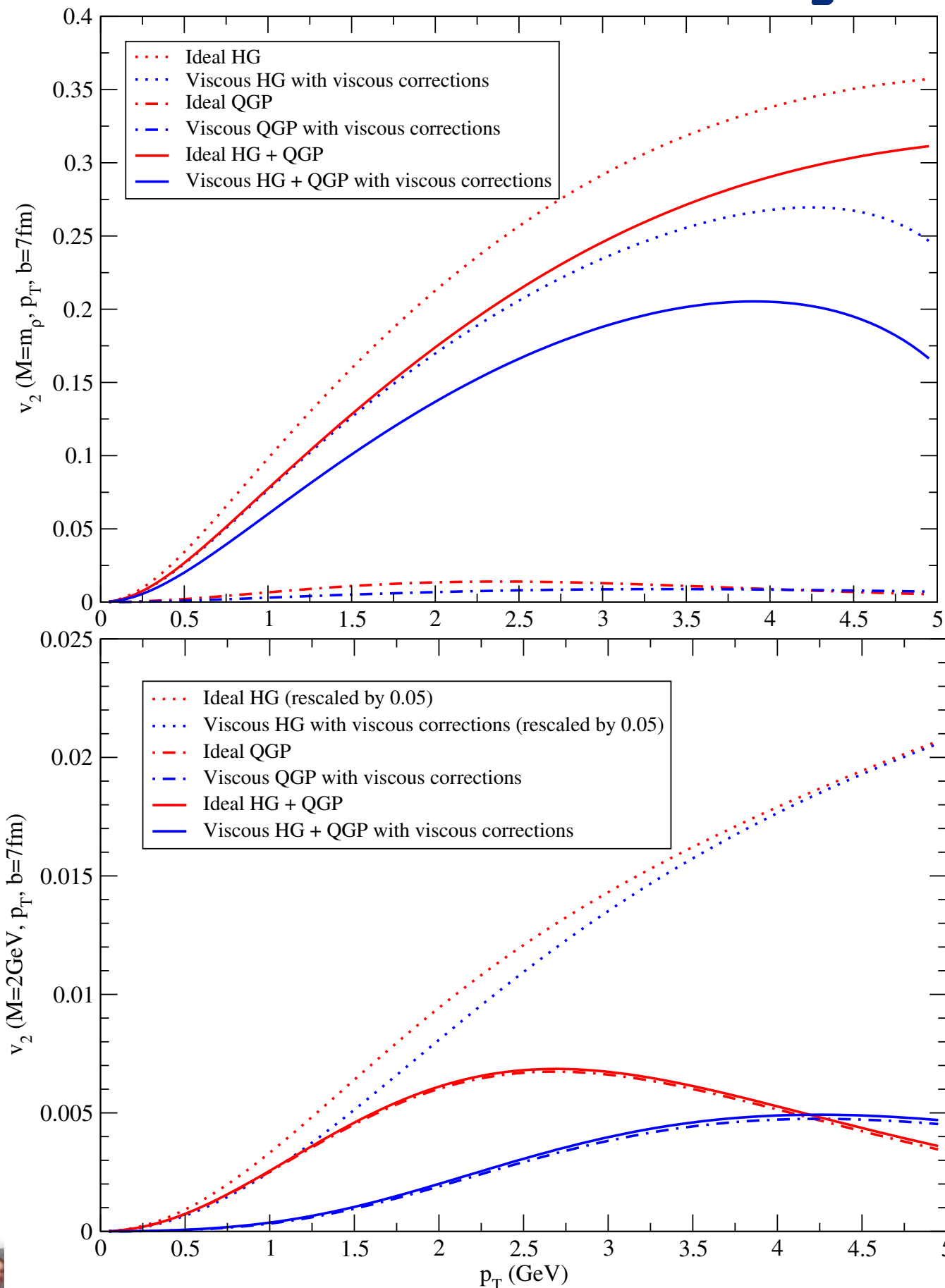
THERMAL DILEPTON SPECTRA: VISCOSITY?



- Transition from HG-dominated to QGP-dominated
- Charm not included here
- Effects of viscous corrections are modest
Dusling & Lin, NPA (2008)
- Same hydro as for photon calculations



THERMAL DILEPTON V_2 WITH VISCOUS EFFECTS



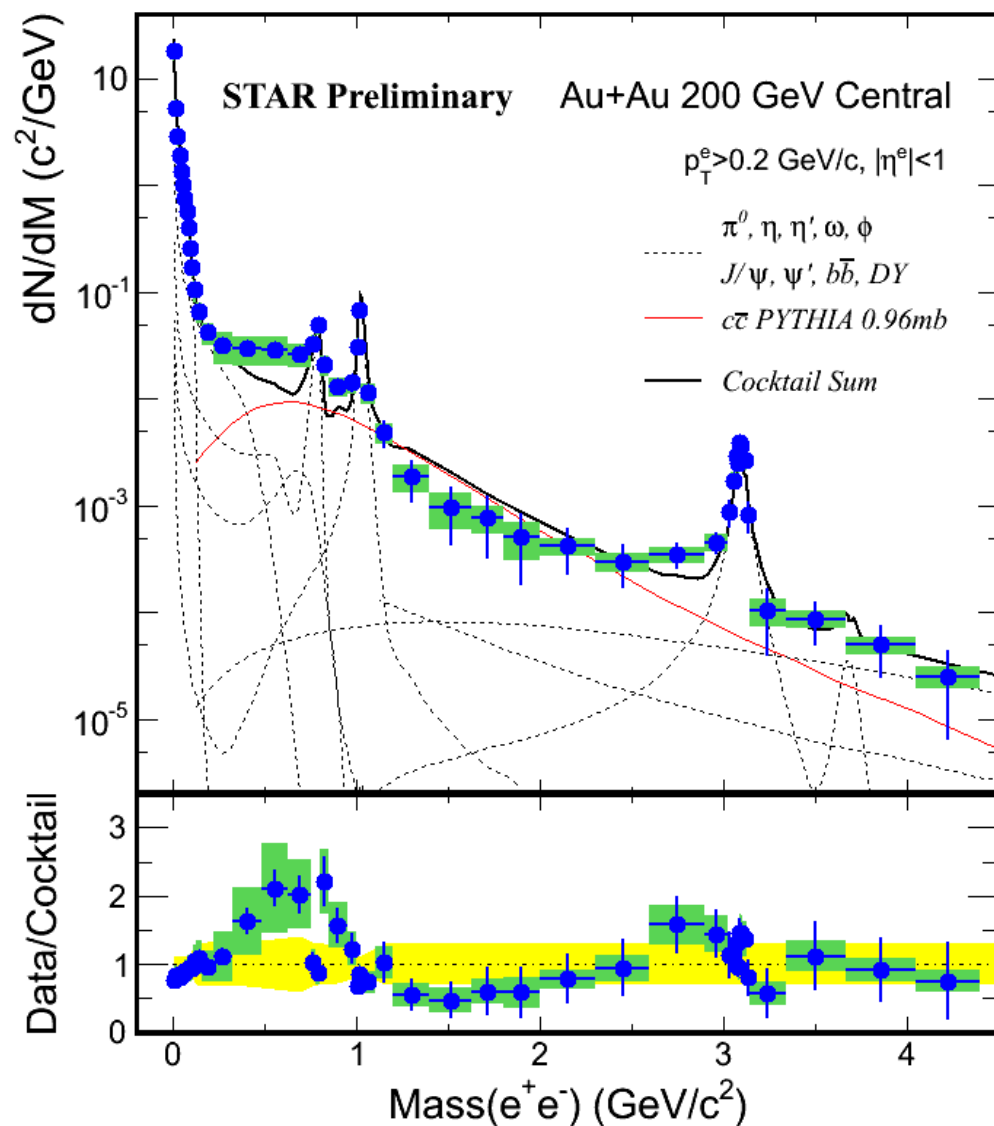
- Low M : HG-dominated
- High- M : QGP dominated

Chatterjee, Srivastava,
Heinz, Gale PRC (2007)

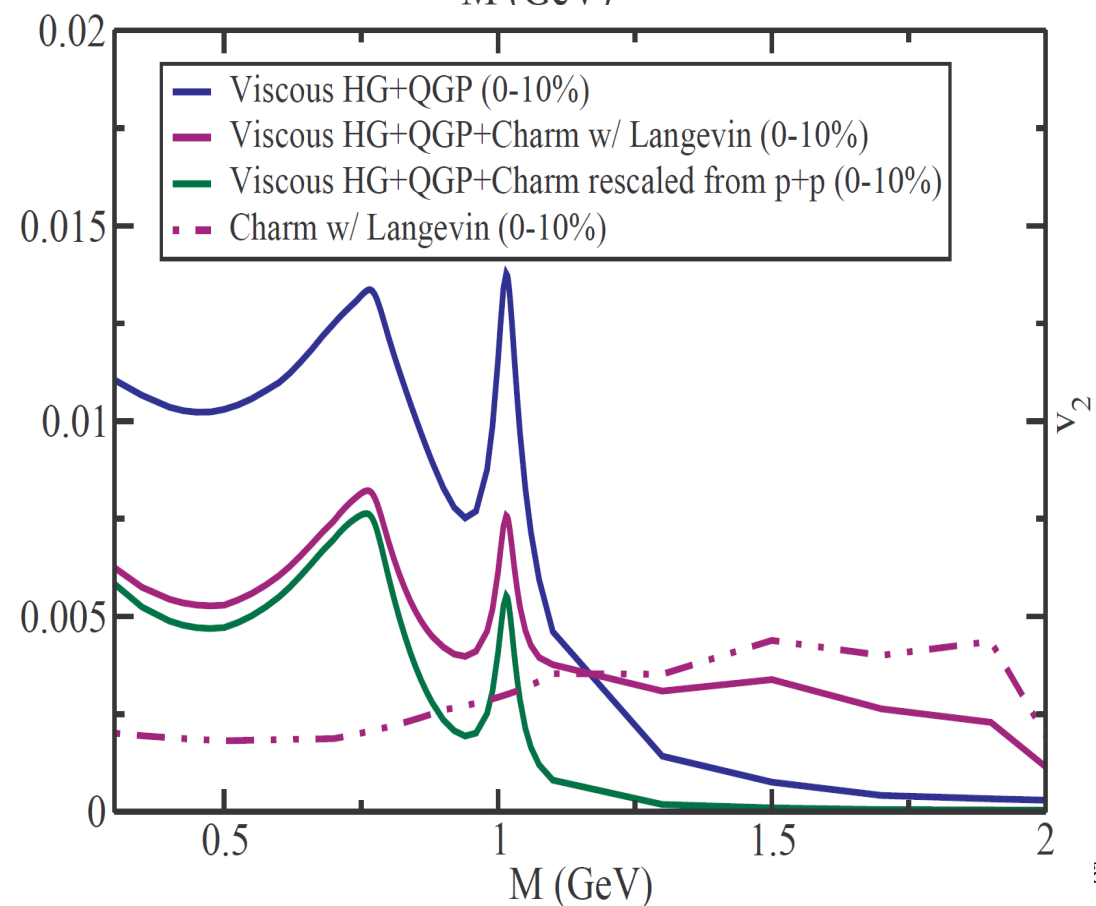
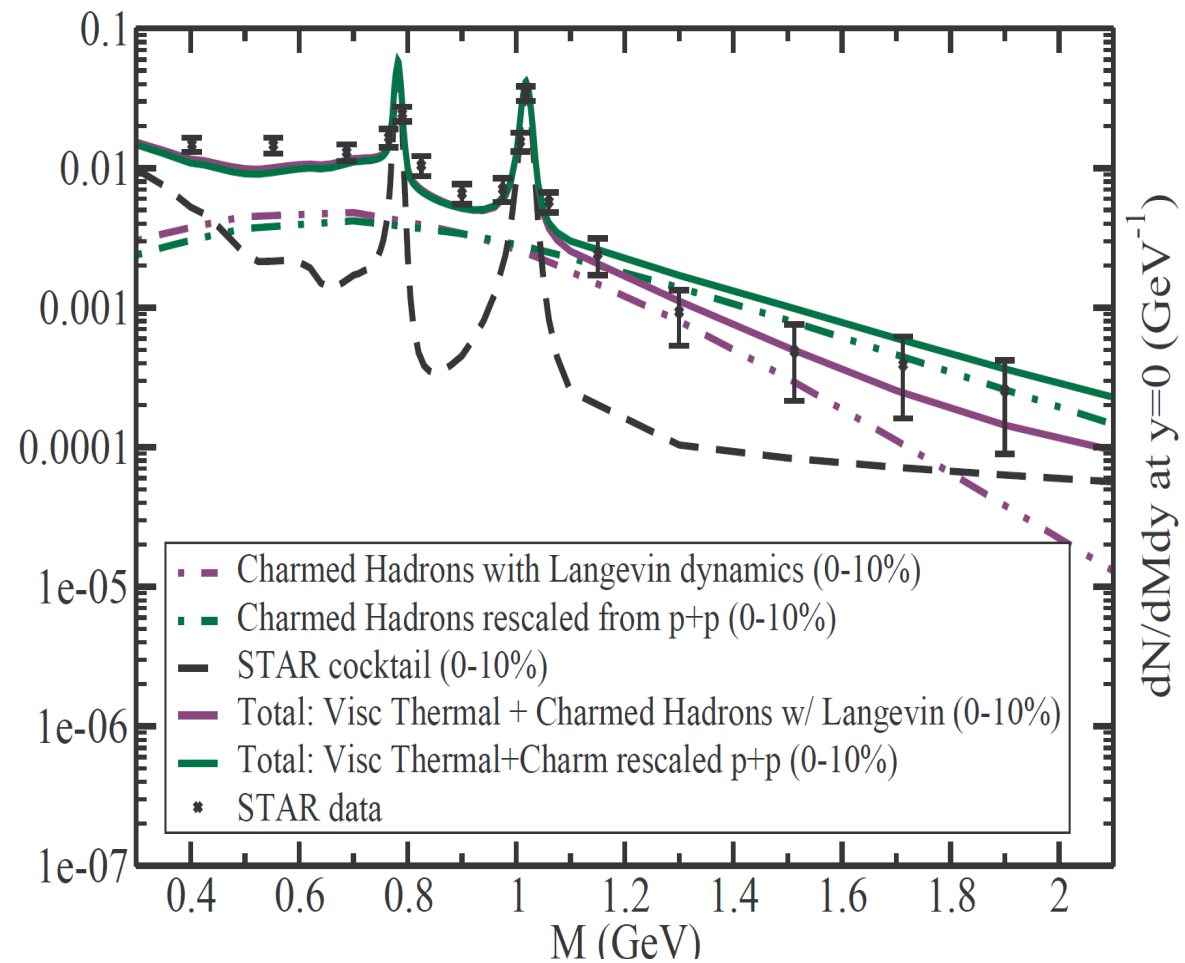
- No open charm here
- v_2 as a function of M will contain some info on the transition regime
- Viscous effects are moderate
- FICs? Coming soon...



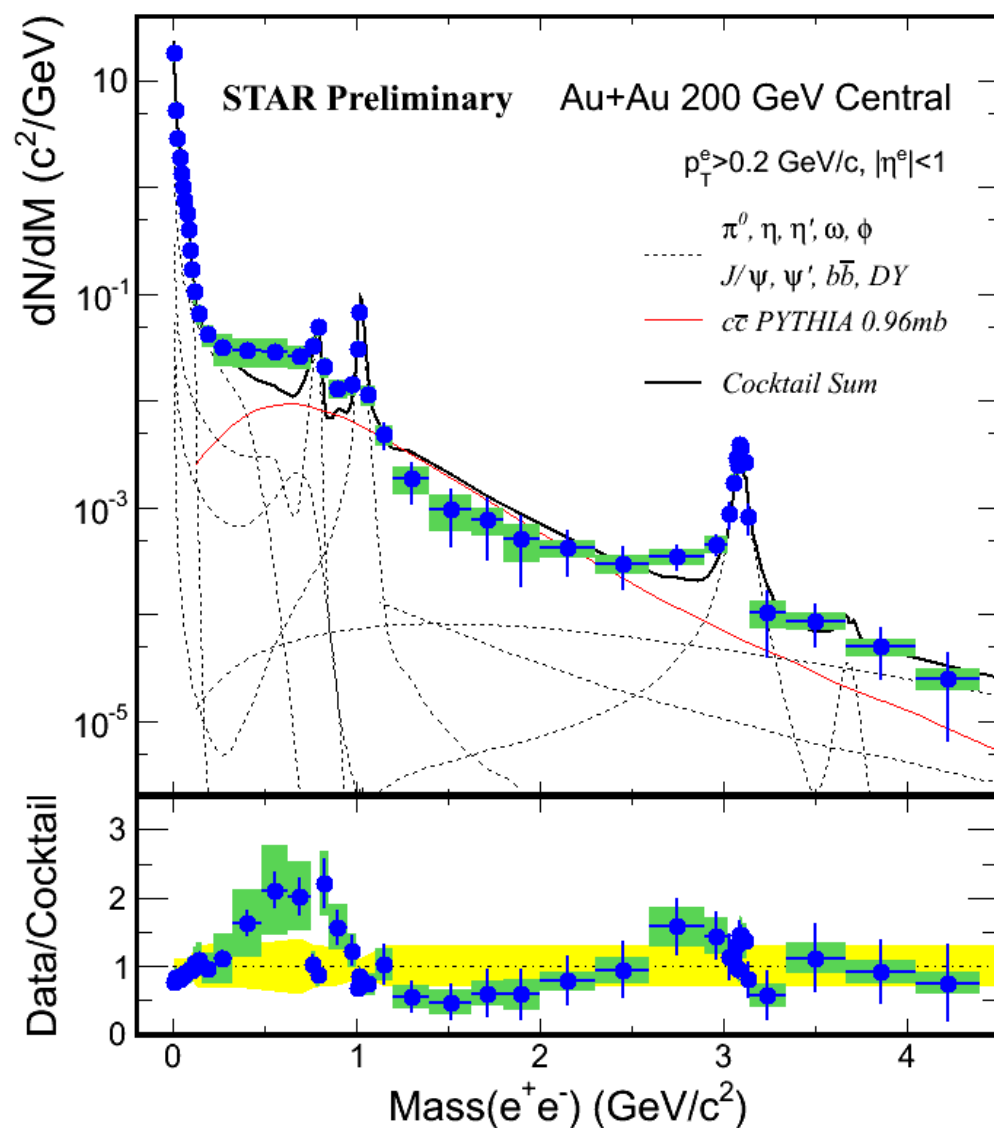
DILEPTONS, SOME RECENT RESULTS



- Uses MUSIC, and rates compatible with NA60 data
- IMR: sensitive to charm energy loss
 - Clint Young et al., arXiv:1111.0647
- IMR: Thermal effects?
 - Li and Gale, PRL (1998)
- FICs to come

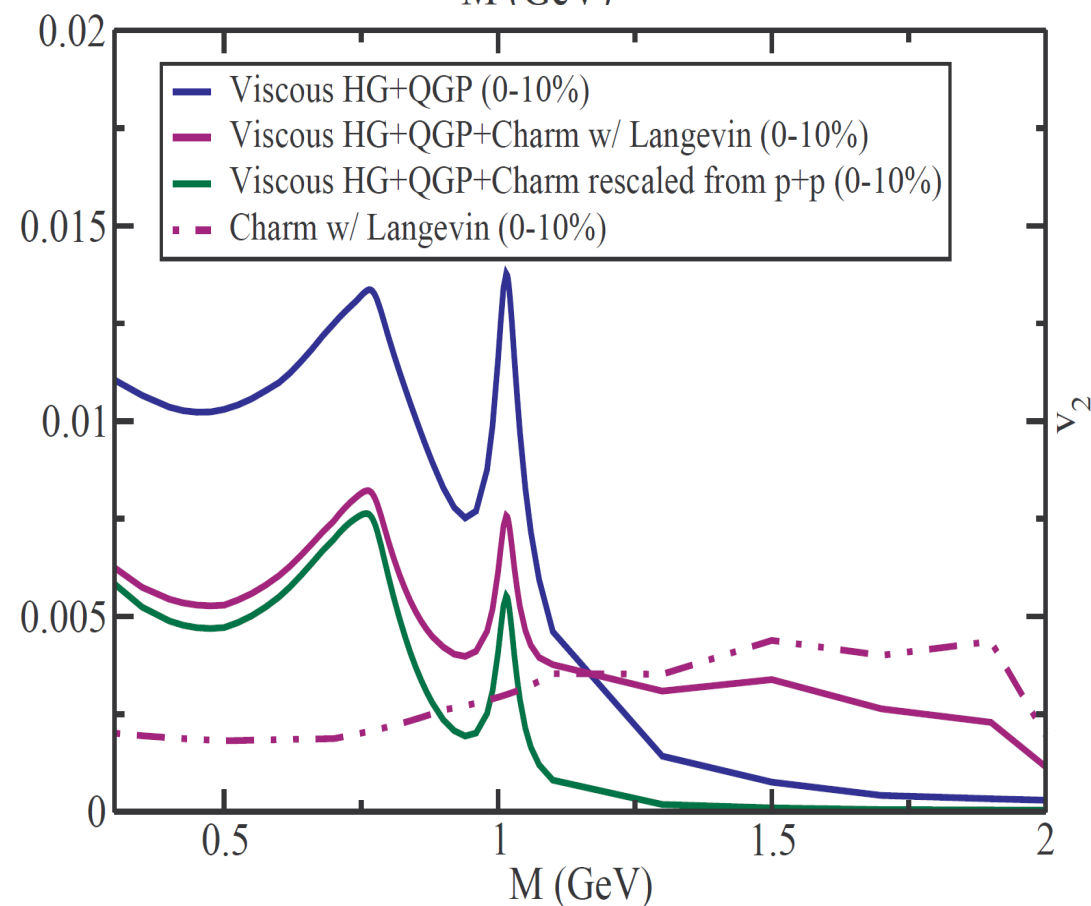
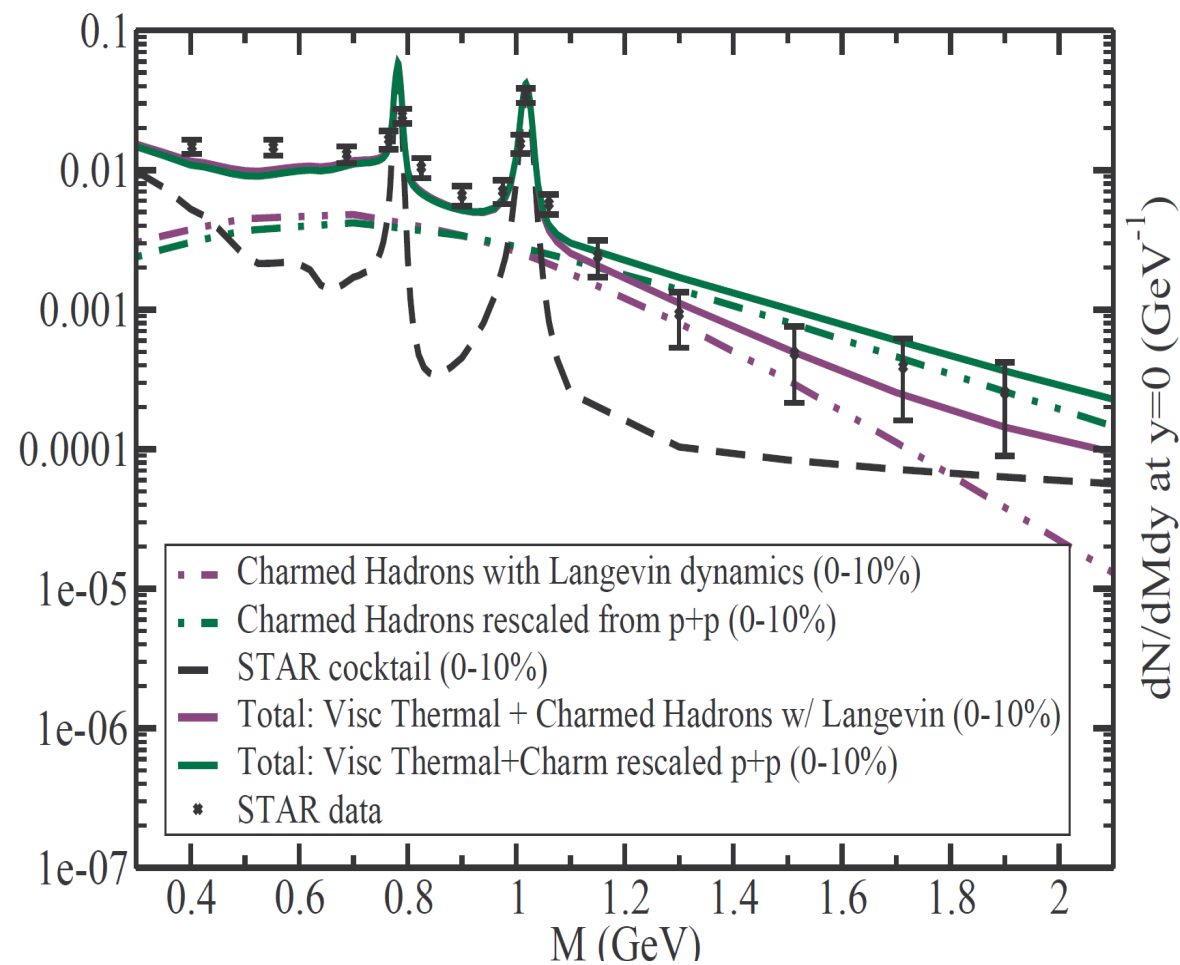


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- FICs to come

PROGRESS



CONCLUSIONS

- The status of EM rates and their integration in dynamical models is still in flux
- Photon v_2 is sensitive to the EOS, and to various hydro parameters such as viscosity, and initial conditions (time and FICs). Current v_2 data is a puzzle. New physics?
- Dilepton v_2 with good statistics is needed to complete the EM emission systematics. STAR & PHENIX?
- FICs and viscosity(ies) make a difference in photon and dilepton characterization: one must be consistent with hadronic data
- Known unknowns: pre-equilibrium radiation, thermal vs. charm components in the dilepton IMR



Heavy-ion collision theory with momentum-dependent interactions

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(Received 29 October 1986)

We examine the influence of momentum-dependent interactions on the momentum flow in 400 MeV/nucleon heavy ion collisions. Choosing the strength of the momentum dependence to produce an effective mass $m^* = 0.7m$ at the Fermi surface, we find that the characteristics of a stiff equation of state can be obtained with a much softer compressibility.



**A GUIDE TO MICROSCOPIC MODELS FOR
INTERMEDIATE ENERGY HEAVY ION COLLISIONS**

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Relativistic Hydrodynamic Theory of Heavy-Ion Collisions*

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By use of finite-difference methods we solve the classical relativistic equations of motion for the head-on collision of two heavy nuclei. For ^{16}O projectiles incident onto various targets at laboratory bombarding energies per nucleon ≤ 2.1 GeV, curved shock waves develop. The target and projectile are deformed and compressed into crescents of revolution. This is followed by rarefaction waves and an overall expansion of the matter into a moderately wide distribution of angles.



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Happy Birthday George!

